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Intrinsic regulation of muscle contraction

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1. Introduction

In the context of cardiac muscle modelling, the active force developed through the interaction between myosin heads and actin sites is regulated at the microscopic level resulting at the organ level into adaptation capabilities, including the Frank-Starling mechanism. This regulation is of two kinds: (i) intrinsic as a result of the sarcomere stretch and (ii) extrinsic through the neuroendocrine regulation. In this work, we focus on the intrinsic regulation, which come from from a variation of the number of available myosin heads and the number of activated actin sites as a result of the stretch of the sarcomere level [4] but leaves the individual actin-myosin mechanical interaction unchanged [2, 10]. We say that the myosin and actin filaments have a varying activation level.

We propose here a new approach that allows to incorporate these two variations into the classical Huxley'57 muscle contraction model equations [7]. This approach allows to bridge the numerous work that have been conducted, on the modelling of the actinmyosin interaction on the one side [6, 9, 1, 3], and the modelling of the thin filament activation on the other side [11, 13], into a complete modelling framework.

2. Model presentation

2.1 Myosin filament activation

The original Huxley'57 model describes the myosin head by two states: attached and detached and assumes that all myosin heads and all actin sites are activated. The varying availability of the myosin heads is introduced by considering two pools of heads: heads that are available for attachment and heads that are not. The heads that are not available can still attach but with a much slower rate. The fraction available heads is given by n₀ and is assumed to depend only on the sarcomere extension ec. In each pool, the probability of being attached for a group of myosin heads located at distance s from their nearest actin site is denoted by $P_1(s, t, \gamma)$, where the additional discrete variable γ describes the belonging to one of the two pools. The variable γ takes the value 1 in the pool of the available heads and takes the value 0 in the other pool. In each pool, the myosin head can transition between the two possible states with the attachment rate f_{γ} and detachment rates g_{γ} . Defining $|x|_{+}$ by x, if $x \geq 0$

and zero otherwise and $|x|_{-}$ = $-x+|x|_{+},$ the dynamics of the system is given by

$$\begin{split} & \frac{d}{dt} \mathsf{P}_1(s,t,1) = \partial_t \mathsf{P}_1(s,t,1) + \dot{x}_c \partial_s \mathsf{P}_1(s,t,1) \\ & + \frac{|\dot{n}_0|_+}{n_0} \Big[\mathsf{P}_1(s,t,1) - \mathsf{P}_1(s,t,0) \Big] = \\ & f_1(s) \big(1 - \mathsf{P}_1(s,t,1) \big) - g_1(s) \mathsf{P}_1(s,t,1), \\ & \frac{d}{dt} \mathsf{P}_1(s,t,0) = \partial_t \mathsf{P}_1(s,t,0) + \dot{x}_c \partial_s \mathsf{P}_1(s,t,0) \\ & + \frac{|\dot{n}_0|_-}{1 - n_0} \Big[\mathsf{P}_1(s,t,0) - \mathsf{P}_1(s,t,1) \Big] = \\ & f_0(s) \big(1 - \mathsf{P}_1(s,t,0) \big) - g_0(s) \mathsf{P}_1(s,t,0), \end{split}$$

where \dot{x}_c is the relative sliding velocity between the myosin and actin filaments.

2.2 Actin filament activation

We now extend this model to rigorously incorporate the variation of the thin filament activation. We supplement our myosin head model with the actin site activation kinetics, which will be characterised by a new parameter n_a , denoting the ratio of activated actin sites. An actin site exists in four possible states: (i) non activated and non occupied by a myosin head, (i) activated and non occupied, (iii) non activated but occupied a myosin head or (iv) activated and occupied. We establish, in each pool of myosin heads, the governing system of equations with conservation laws. The system can be reduced to

$$\begin{cases} \displaystyle \frac{d}{dt} P_1(s,t,\gamma) = f_{\gamma,a}(s)\overline{n}_a(s,t,\gamma) \\ &+ f_{\gamma,na}(s) \big(1 - P_1(s,t,\gamma) - \overline{n}_a(s,t,\gamma)\big) \\ &- g_{\gamma,a}(s) \big(n_a - \overline{n}_a(s,t,\gamma)\big) \\ &- g_{\gamma,na}(s) \big(\overline{n}_a(s,t,\gamma) - n_a + P_1(s,t,\gamma)\big), \\ \\ \displaystyle \frac{d}{dt} \overline{n}_a(s,t,\gamma) = \frac{|\dot{n}_a|_+}{1 - n_a} \big(1 - P_1(s,t,\gamma) - \overline{n}_a(s,t,\gamma)\big) \\ &+ g_{\gamma,a}(s) \big(n_a - \overline{n}_a(s,t,\gamma)\big) \\ &- \frac{|\dot{n}_a|_-}{n_a} \overline{n}_a(s,t,\gamma) - f_{\gamma,a}(s)\overline{n}_a(s,t,\gamma), \end{cases}$$

where \overline{n}_a denotes the ratio of actin sites that are activated but non-occupied. The active tension is given by the average between the two pools of myosin heads.

3. Validation

We validate our model by comparing its outputs with experimental data. The model is calibrated

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choosing n_0 such that the predicted tension interpolates data points in maximal calcium activation conditions. The input parameter n_a is written as $n_a = n_{a,\infty}(e_c,C)n_{a,t}(e_c,t)$ The function $n_{a,\infty}(e_c,C)$, where $C \in [0,1]$ denotes a normalised level of calcium activation, accounts for the steady-state activation and is chosen as a sigmoid function to mimic the experimental data obtained on skinned cells [5]. The function $n_{a,t}(e_c,t)$ captures the transient nature of the thin filament activation.

We first compare the steady-state model prediction with experimental data obtained in [12] (see Figure 1) and note that our model displays a good match with the data. We then compare the dynamics predicted by the model with data from [8]. We simulate twitch contractions in "high" and "low" contractility conditions (corresponding to a "high" and "low" level of calcium supply, respectively) in isometric conditions, i.e. with $\dot{x}_c = 0$ (see Figure 2). The results match well the experimental data in high contractility conditions. However, in low contractility conditions, the decreases of the peak force with the sarcomere extension is faster in the simulations than in the data. One can note that this decrease originates from the constraints imposed on the function $n_{a,\infty}$ by the steady-state data. However, these data are obtained on a range of sarcomere extensions where the function $n_{a,\infty}$ extrapolates the data.



Figure 1: Steady state thin filament activation. Experimental data (black closed symbols) and model outputs (coloured lines).

4. Conclusion

In this work, we proposed a novel framework to rigorously incorporate the varying myosin heads and actin sites activation level into the Huxley'57 model family. This framework may be coupled to a separately defined thin filament activation model or, as done in this work, used as a standalone model with the timedependent thin filament activation being provided as a model input. The model is calibrated on experimental twitch contractions obtained in controlled sarcomere length conditions and a preliminary validation is performed.

5. References

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Figure 2: Transient twitch contraction at various sarcomere stretch in sarcomere length control conditions. Experimental data (red lines) and model outputs (blue lines). The sarcomere extension e_c varies between 0.027 and 0.189.

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