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## ARTICLE OPEN



# Orbital selective switching of ferromagnetism in an oxide quasi two-dimensional electron gas

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Multi-orbital physics in quasi-two-dimensional electron gases (q2DEGs) triggers intriguing phenomena not observed in bulk materials, such as unconventional superconductivity and magnetism. Here, we investigate the mechanism of orbital selective switching of the spin-polarization in the oxide q2DEG formed at the (001) interface between the LaAlO<sub>3</sub>, EuTiO<sub>3</sub> and SrTiO<sub>3</sub> band insulators. By using density functional theory calculations, transport, magnetic and x-ray spectroscopy measurements, we find that the filling of titanium-bands with 3d<sub>xz</sub>/3d<sub>yz</sub> orbital character in the EuTiO<sub>3</sub> layer and at the interface with SrTiO<sub>3</sub> induces an antiferromagnetic to ferromagnetic switching of the exchange interaction between Eu-4f<sup>7</sup> magnetic moments. The results explain the observation of the carrier density-dependent ferromagnetic correlations and anomalous Hall effect in this q2DEG, and demonstrate how combined theoretical and experimental approaches can lead to a deeper understanding of emerging electronic phases and serve as a guide for the materials design of advanced electronic applications.

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## INTRODUCTION

Since the discovery of a quasi-two-dimensional electron gas (q2DEG) at the interface between the LaAlO<sub>3</sub> (LAO) and SrTiO<sub>3</sub> (STO) band insulators<sup>1</sup>, studies on oxide surfaces and interfaces uncovered an intriguing and rich physics, boosting the expectations for all-oxide electronics. The LAO/STO q2DEG shows remarkable properties, including electric field effect induced insulator-to-metal transition at room temperature<sup>2</sup>, gate-tunable Rashba-like spin-orbit coupling<sup>3,4</sup>, superconductivity<sup>5–7</sup>, and magnetism<sup>8</sup>. From these studies, a fundamental role of the multi-orbital nature of the carriers in the electronic properties of the oxide q2DEGs clearly emerged.

The possibility to induce ferromagnetic (FM) correlations at the interface between non-magnetic oxides, combined with the large spin to charge conversion efficiency of oxide q2DEGs<sup>9</sup>, paves the way to applications in spintronics.

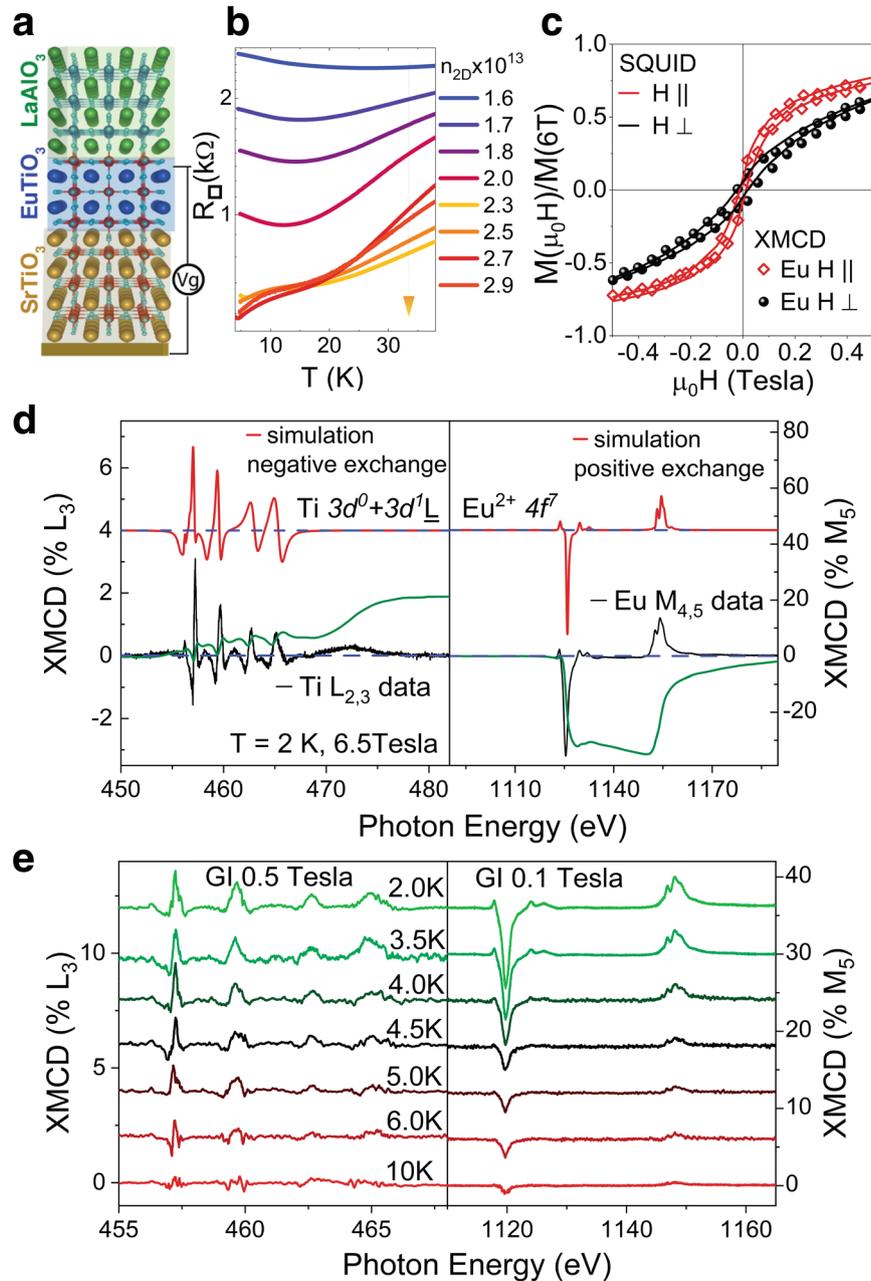
Recently, it has been shown that a feasible method to induce a spin-polarization in oxide q2DEGs is the introduction of a thin magnetic layer between LAO and STO, like EuTiO<sub>3</sub> (ETO) (Fig. 1a)<sup>10,11</sup> and LaMnO<sub>3</sub><sup>12</sup>. FM correlations were reported also in heterostructures where LAO is replaced by a ferro(ferri)-magnetic insulator, as in GdTlO<sub>3</sub>/STO<sup>13,14</sup>, EuO/STO<sup>15</sup> and LaAl<sub>1-x</sub>Mn<sub>x</sub>O<sub>3</sub>/STO<sup>16</sup>.

In LAO/ETO/STO, FM correlations are believed to be induced by the ordering of localized Eu<sup>2+</sup> 4f<sup>7</sup> magnetic moments and their coupling with Ti-3d states forming the conduction band of the q2DEG. However, the microscopic mechanism leading to these phenomena is not straightforward, as rare earth magnetic ions have usually a small hybridization with transition metal 3d orbitals. Unlike the LAO/STO q2DEG<sup>1–9</sup>, there is no detailed study of the

electronic band structure of LAO/ETO/STO heterostructures, and it remains unclear whether a spin-polarized q2DEG is present in both ETO and interfacial STO layers, which is relevant for the intriguing phase diagram showing a transition from a FM to a superconducting state<sup>11</sup>. Additionally, FM correlations were observed only above the Lifshitz transition, where carriers with 3d<sub>xz</sub>/3d<sub>yz</sub> orbital character start to contribute to the transport<sup>11</sup>.

In this work, all these questions are settled by combining extensive experimental and theoretical investigations. Firstly, we show that low-temperature electrical transport, x-ray magnetic circular dichroism (XMCD) and superconducting quantum interference device (SQUID) experiments give evidence of a tunable ferromagnetism originating from the correlation between Ti- and Eu- magnetic moments. Secondly, we provide a direct picture of the q2DEG band structure by using resonant soft-x-ray photoemission spectroscopy (RESPE). These experimental observations are finally compared and blended with density functional theory calculations with on-site Hubbard terms (DFT + *U*). The latter show that the q2DEG forms at the LAO/ETO interface and extends few unit cells into the STO. FM correlations are induced by the delicate balance between different antiferromagnetic (AFM) and FM contributions to the exchange interaction between Eu-4f<sup>7</sup> magnetic moments and Ti-3d electrons at the ETO/STO interface, which involves also the first STO layers, in agreement with the experiments. Moreover, the theoretical calculations show that the ferromagnetism in the q2DEG and the filling of 3d<sub>xz</sub>/3d<sub>yz</sub> bands take place simultaneously, explaining the electric field-induced anomalous Hall effect above the Lifshitz point<sup>11</sup>, and the orbital selective switching of the spin-polarization in an oxide q2DEG.

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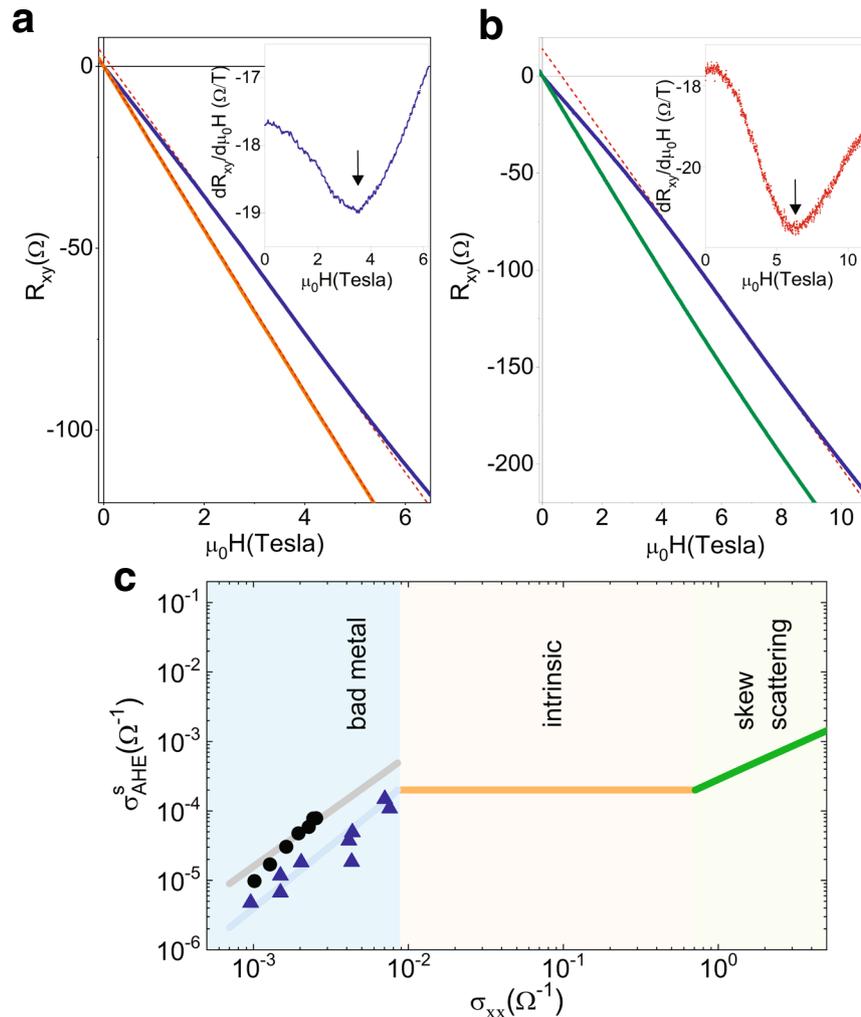
**Fig. 1** Transport and magnetic properties of LAO/ETO/STO heterostructures. **a** Sketch of the LAO/ETO/STO heterostructure. In this work, the carrier concentration was tuned using electric field effect in back-gate configuration, as sketched in the figure. **b** Sheet resistance vs. temperature as a function of the gate voltage  $V_g$ . The arrow indicates increasing values of  $V_g$  from  $-30$  to  $30$  V and 2D carrier density. **c** Eu-XMCD (scatter data) and SQUID magnetization (continuous lines) as function of the magnetic field parallel (red) and perpendicular (black) to the interface. The data are normalized to the saturation value. The Eu-XMCD in parallel direction is obtained by combing data acquired in grazing (60 degrees from the surface normal) and normal incidence conditions. **d** XMCD data at the Ti-L (left) and Eu-M (right) edges at 6.5 Tesla and 2 K. Red-lines are atomic multiplet simulations, green lines are integrals of the XMCD spectra (see main text). **e** Temperature dependence of the Ti-L<sub>2,3</sub> (0.5 Tesla) and Eu-M<sub>4,5</sub> (0.1 Tesla) edges XMCD spectra in grazing incidence conditions. Data are vertically shifted for clarity.

## RESULTS

### Experimental evidences of intrinsic ferromagnetic correlations

Bulk ETO is an AFM insulator with a  $T_N$  of 5.5 K, however doping or lattice strain favor a FM order in epitaxial films<sup>17–21</sup>. It has a quasi-perfect structural matching with STO and a similar conduction band formed by the overlap of  $t_{2g}$  Ti-3d orbitals. By embedding few unit cells (uc) of ETO between LAO and STO using epitaxial growth by pulsed laser deposition (PLD), we realize a q2DEG characterized by electrical transport properties similar to the LAO/

STO system. However, at low temperatures and at a carrier density  $>2.10^{13}$  cm<sup>-2</sup> tuned by electric field effect (Fig. 1a), the sheet resistance exhibits a downturn below  $T=6-8$  K (Fig. 1b), which correlates with the FM transition of doped ETO<sup>11</sup>. The heterostructure itself is characterized by a FM transition ( $T_c = 6-8$  K) and a FM ground state, as shown in Fig. 1c, where we compare the maximum intensity of the Eu-XMCD signal (at  $M_5$ -edge, scatter data), acquired at the X-Treme beamline of the Swiss Light Source (SLS)<sup>22</sup>, to macroscopic SQUID magnetometry data (full lines) at



**Fig. 2 Analysis of the anomalous Hall effect.** Transverse (Hall) resistance  $R_{xy}$  as a function of the magnetic field of: (a) sample A at two gate voltages ( $V_g = +60$  V (blue line),  $V_g = -25$  V (orange line)) and 2 K; (b) sample B at  $V_g = +30$  V (2 K (blue line) and 20 K (green line)). Corresponding 2 K  $dR_{xy}/d\mu_0 H$  data are shown in the insets of each figure, with the arrow indicating the minimum. The minimum of the  $dR_{xy}/d\mu_0 H$  is the slope of the red dashed lines fitting the data in the intermediate 3-4 Tesla range. The intercepts of these lines are the AHE components at saturation. c Anomalous Hall conductivity vs. longitudinal 2D conductivity (see main text Eq. (1)). Black dots and blue triangles refer to the two LAO/ETO/STO samples. Full lines refer to the exponent  $\alpha$  in the relation  $\sigma_{AHE} \propto (\sigma_{xx})^\alpha$ :  $\alpha = 1.8$  (light blue, and gray, intrinsic AHE suppressed by disorder),  $\alpha = 0$  (orange, intrinsic AHE),  $\alpha = 1$  (green, skew scattering).

2 K temperature as a function of the magnetic field in parallel (red) and perpendicular (black) directions. The two sets of data, normalized to the saturation value for direct comparison, show a small, although clear, hysteresis around zero magnetic field and a preferential orientation of the magnetization parallel to the interface, confirming a remanent zero-field magnetization and FM correlations. The magnetization saturates above 2–4 Tesla (depending on the sample and on the magnetic field direction) and a spin moment of the order of  $4 \mu_B/\text{Eu}$  can be estimated from the XMCD using the sum-rules, indicating that  $\text{Eu}^{2+}$ -ions order ferromagnetically above these values.

To further elucidate the magnetic properties of the system, in Fig. 1d we show XAS and XMCD spectra at the Ti- $L_{2,3}$  and Eu- $M_{4,5}$  edges at 2 K and 6.5 Tesla and in Fig. 1e the XMCD temperature dependence in grazing incidence conditions (60 degrees from the surface normal). The data show that Eu- and Ti- XMCD spectra follow each other, thus are correlated. In particular the Ti- (orbital and spin) and the Eu- magnetic moments have the same temperature dependence, with a  $T_c$  of the order of 6–8 K and the same magnetic field dependence as shown in ref. 11.

In Fig. 1d, the XMCD spectra are compared to atomic multiplet simulations (red lines) for  $\text{Eu}^{2+}$  and  $\text{Ti}^{4+}$  ions in  $C_4$  symmetry. In the calculations for the Ti-3d XMCD we used the crystal field splitting reproducing the x-ray linear dichroism data, and included a charge transfer (CT) term to account for the hybridization between Ti-3d and O-2p states in the  $\text{TiO}_6$  cluster, i.e., a  $3d^0 + 3d^1 \underline{L}$  configuration, where  $\underline{L}$  indicates a hole in the O-2p band (see Supplementary Note 1, Supplementary Fig. 1 and ref. 23). The calculations reproduce most of the features shown by the experiment. However, while a positive exchange is needed to simulate the  $\text{Eu}^{2+}$  magnetization, parallel to the magnetic field, only a negative one correctly reproduces the Ti-XMCD peaks at  $L_3$  and  $L_2$ . On the other hand, we also find a finite and positive Ti-XMCD integral that, according to the sum rules<sup>24</sup>, corresponds to a sizeable Ti-3d orbital moment,  $m_{\text{orb}}$ , of the order of  $-0.05 \mu_B/\text{Ti}$ . Following the Hund's rules, the sign of the orbital moment corresponds to an opposite, finite, Ti-spin moment,  $m_{\text{spin}}$ , parallel to the Eu-spin moment. The order of magnitude of  $m_{\text{spin}}$  is  $0.07 \mu_B/\text{Ti}$ , thus the total Ti-magnetic moment is of the order of  $0.02\text{--}0.03 \mu_B/\text{Ti}$ . However, this value has to be considered only an estimate for the order of magnitude of the moment, since the spin

sum rules are heavily affected from the mixing of the  $L_3$  and  $L_2$  edges and inadequate for a precise quantification<sup>24</sup>.

However, the Ti-XMCD integral and the Ti- magnetic moment are expected to be null for a  $Ti^{4+}$  state, even in the presence of a CT between Ti-3d and neighbor O-2p states. Thus, the only way to explain the results is the presence of two contributions to the spectra: (i) a  $3d^0 + 3d^1L$  contribution, with a negative exchange interaction between the ligand Ti-3d electrons and the  $Eu^{2+}$  magnetic moments, which gives rise to the observed sign of the XMCD; (ii) a second contribution coming from Ti- $3d^1$  electronic states, with their associated magnetic moments, parallel to the magnetic field and to the  $Eu^{2+}$  spin-moments, which explain the non-zero integral of the XMCD. These results suggest different typologies of exchange interactions between Eu and Ti ions, i.e., an AFM one between  $Eu^{2+}$  and  $Ti-3d^0 + 3d^1L$  states, and a FM exchange between the  $Eu^{2+}$  and  $Ti-3d^1$  spin-moments.

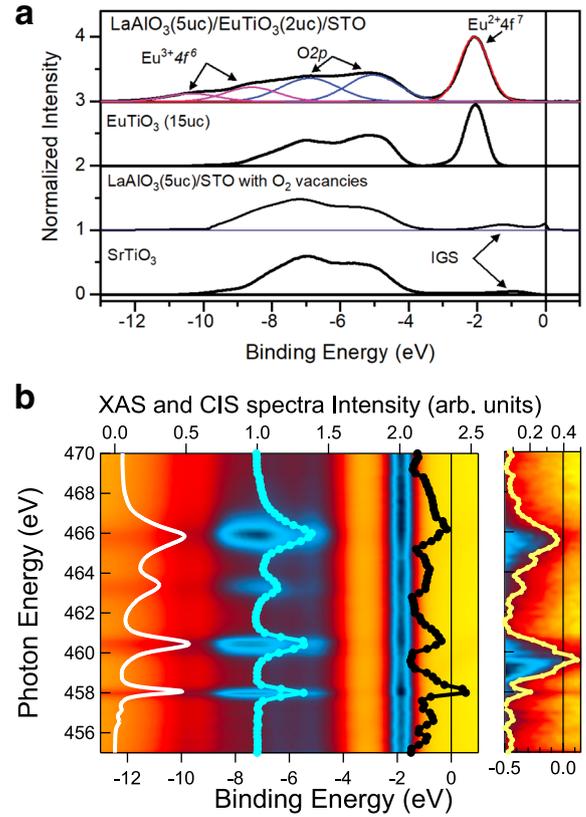
The FM correlations in the q2DEG, and the magnetic field dependence of the magnetization, reveals in the Hall effect data<sup>11</sup>. In Fig. 2a, b we show data on two representative samples, sample A (Fig. 2a,  $V_g = -25$  V and  $V_g = +60$  V at 2 K), and sample B (Fig. 2b,  $V_g = +30$  at 2 K and at 20 K, well-above the FM Tc). A linear Hall effect is present for low gate voltage (curve at  $V_g = -25$  V in Fig. 2a), indicating single band transport. For positive gate voltages and low temperature, instead, both samples exhibit a non-linear Hall effect. The 2 K non-linear Hall effect is characterized by a high field (>5–6 Tesla) curvature change due to the two-band transport above the Lifshitz point ( $n_{2D} = 1.9 \cdot 10^{13} \text{ cm}^{-2}$ ), associated to the contribution of both light ( $3d_{xy}$ ) and heavy ( $3d_{xz}/3d_{yz}$ ) bands, and a low field (<2–4 Tesla) curvature change due to the anomalous Hall effect (AHE). Following the method used in refs. 11,18, the 2D anomalous resistivity at saturation,  $\rho_{AHE}$  is the intercept to zero field of a line fitting the data in the intermediate field region (red dashed lines in Fig. 2a, b). This region can be precisely identified as the field-range at which the  $dR_{xy}/d\mu_0H$  (insets of Fig. 2a, b) reaches a minimum, and the slope of the fitting line is the corresponding value of  $dR_{xy}/d\mu_0H$  (more details about the fitting procedure are shown in Supplementary Note 2 and Supplementary Figs. 3 and 4). Above the FM Tc, the AHE vanishes, and the Hall effect data can be nicely reproduced by a simple two-band model (Fig. 2b).

Further insights on the nature of the AHE is obtained by comparing the 2D anomalous conductivity,  $\sigma_{AHE}$ , to the longitudinal one,  $\sigma_{xx}$  estimated according to Eqs. (1) and (2):

$$\sigma_{xx} = \frac{\rho_{xx}^{2D}(0)}{(\rho_{AHE})^2 + (\rho_{xx}^{2D}(0))^2} \quad (1)$$

$$\sigma_{AHE} = \frac{\rho_{AHE}}{(\rho_{AHE})^2 + (\rho_{xx}^{2D}(0))^2} \quad (2)$$

where  $\rho_{xx}^{2D}(0)$  is the zero field 2D longitudinal resistivity. The origin of the AHE can be attributed to different mechanisms, which can be phenomenologically distinguished using the relation between  $\sigma_{AHE}$  and  $\sigma_{xx}$ . According to the literature<sup>25</sup>, most of the materials showing an AHE can be classified in the following three regimes: (i) a high conductivity range, where  $\sigma_{AHE} \propto \sigma_{xx}$  and the observed AHE is attributed to scattering from unwanted magnetic impurities (skew scattering); (ii) an intermediate conductivity range, where  $\sigma_{AHE}$  is independent from  $\sigma_{xx}$  and the AHE is related only to the Berry curvature of the involved bands (intrinsic AHE); (iii) a bad metal range where  $\sigma_{AHE} \propto (\sigma_{xx})^\alpha$ , with  $\alpha$  in the interval 1.6–1.8, where the intrinsic AHE is suppressed by the disorder. In the latter two ranges the AHE is due to the spin-polarization of the carriers and to the spin-orbit coupling breaking time and inversion symmetry. Data collected on LAO/ETO/STO samples, with carrier densities tuned by a back-gate voltage, are shown in Fig. 2c. We found a correlation  $\sigma_{AHE} \propto (\sigma_{xx})^\alpha$ , with  $\alpha = 1.8$ , which excludes a skew scattering mechanism and a purely intrinsic AHE. We



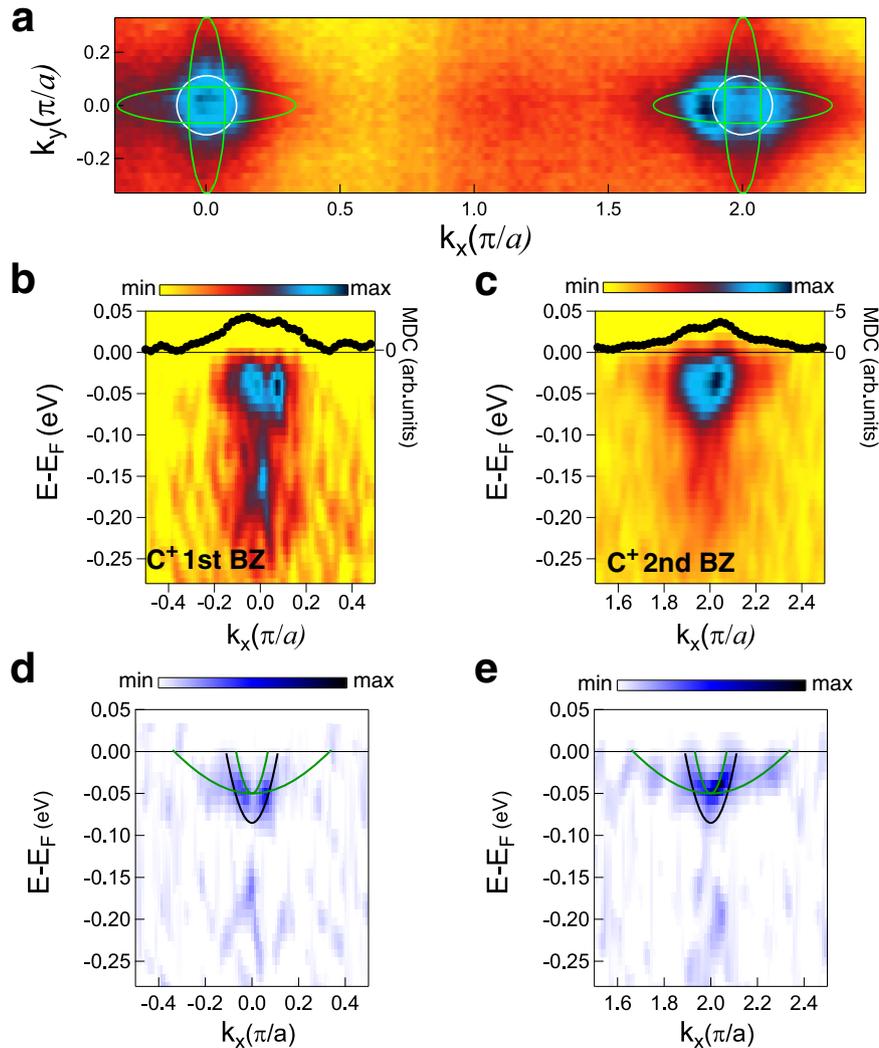
**Fig. 3 RESPEs and ARPES VB data.** **a** Angle integrated ARPES VB data of  $SrTiO_3$ , and  $EuTiO_3$  q2DEG surface state, compared to angle and photon-energy integrated RESPEs VB of (001) LAO(5uc)/STO (containing oxygen vacancies) and LAO(5uc)/ETO(2uc)/STO heterostructures. Multiple gaussian fit of the LAO/ETO/STO VB profile is used to determine the  $Eu^{3+}/Eu^{2+}$  ratio. The IGS in gap states is observed only in STO and LAO/STO q2DEGs due to the presence of oxygen vacancies. **b** (left panel) RESPEs color-map of the valence band region and CIS spectra obtained integrating the RESPEs data around  $BE = -2$  eV ( $Eu^{2+}$ , black scatter data) and between  $-7$  and  $-4$  eV (O-2p, cyan scatter data). The white line is the total electron yield XAS spectrum; (right panel) the same RESPEs map in the Fermi level region and corresponding CIS spectrum (yellow scatter data).

conclude that in LAO/ETO/STO an intrinsic AHE suppressed by disorder is at play, analogously to what found in La-doped ETO films<sup>19</sup>.

### Experimental investigation of the band structure via RESPEs

In order to characterize the electronic structure of the LAO/ETO/STO system, we used the RESPEs technique at the soft-x-ray end station of the ADvanced REsonant Spectroscopies (ADRESS, X03MA) beamline of the Swiss Light Source (SLS)<sup>26,27</sup>. The experimental data were acquired at base temperature (12 K) thus above the FM transition of the system.

In Fig. 3a we report angle and photon-energy integrated (across the Ti- $L_{2,3}$  edge) RESPEs VB data on (001) LAO(5uc)/ETO(2uc)/STO and on a reference PLD grown (001) LAO(5uc)/STO sample containing oxygen vacancies, i.e., cooled down from the high deposition temperature (750 °C) in a reduced oxygen pressure of the order of  $10^{-5}$  mbar, without any high- $O_2$  pressure annealing process. For comparison we show also data acquired by high resolution ARPES at the SIS (Surface Interface Spectroscopy) beamline of SLS ( $h\nu = 85$  eV), on reference VB of an STO crystal and of an ETO film hosting q2DEGs<sup>28</sup>. Since the escape depth of the photoelectrons in the 450–470 eV range is of the order of 2 nm, in LAO(5)/ETO(2)/STO the data have contributions from both



**Fig. 4 Resonant angle-resolved photoemission spectra.** **a** Average between circular plus and circular minus Fermi surface map of the (001) LAO/ETO/STO heterostructure. Dashed lines are Fermi surface contours of light (white) and heavy (green) bands. **b, c**  $E$  vs.  $k_x$  band dispersion maps acquired with C+ polarization in the 1<sup>st</sup> and 2<sup>nd</sup> BZ. Black scatter data on top of each band dispersion map are the Fermi level MDCs curves (integrated around  $E_F$  in a 10 meV range). All the data are measured in resonance conditions (465.5 eV). **d, e** two-dimensional curvature maps of the C+ dispersions in **b** and **c**, with tight-binding fitting of the bands (green lines  $3d_{xz}$  and  $3d_{yz}$  bands, black line  $3d_{xy}$  band).

the ETO layers and from the interfacial STO unit cells. As a result, the VB spectra of ETO heterostructures show prominent differences with the STO and LAO/STO VB. The main one is the presence, in ETO and in LAO/ETO/STO, of a peak at  $-1.95$  eV binding energy (BE), due to  $\text{Eu}^{2+} - 4f^7$  states. Another difference, is the absence in ETO q2DEGs of the  $-1.0$  eV in-gap state (IGS), instead observed in STO and oxygen-deficient LAO/STO<sup>27</sup>. Finally, in LAO/ETO/STO we observe the evidence of a further broad feature at BE lower than  $-8$  eV, which reflects the contribution of Eu in  $\text{Eu}^{3+}$  oxidation state: through a Gaussian deconvolution of the observed spectral features (reported in Fig. 3a), we estimated a fraction of  $\text{Eu}^{3+}$  of about 25%. This result is in agreement with atomically resolved high-resolution transmission electron microscopy and electron energy loss spectroscopy<sup>11</sup>, which shows that a similar fraction of  $\text{Eu}^{3+}$  is present in the system, mostly located into the LAO layer at the interface with the ETO film. The presence of non-magnetic  $\text{Eu}^{3+}$  also explains the reduced value of the saturation magnetization per Eu ions estimated from the XMCD data ( $4 \mu_B/\text{Eu}$  against  $7 \mu_B/\text{Eu}$ ).

In Fig. 3b we report angle integrated VB-RESPES data of LAO (5uc)/ETO(2uc)/STO at photon energies  $h\nu$  across the Ti-2p absorption edge resonance. The data are overlapped to TEY XAS (white line), and to Constant Initial State (CIS) spectra obtained by integrating the valence band over different binding energy ranges. CIS spectra allows the identification of the resonance of different VB features with the XAS intensity spectra. We can see that the same XAS resonances characteristic of  $\text{Ti}^{4+}$  ions appear as higher intensity signal in the RESPES map inside the contribution from the O-2p band. Similar results have been reported for the LAO/STO q2DEG<sup>27</sup>, and are a consequence of the hybridization between O-2p and Ti-3d states. On the other hand, the CIS spectrum around the Fermi level, related to the q2DEG conduction band, has a different shape, reminiscent of the characteristic absorption from  $\text{Ti}^{3+}$  ions in an octahedral environment, similar to bulk  $\text{LaTiO}_3$ <sup>29</sup> (see Supplementary Fig. 2). Surprisingly, also the peak at  $-2$  eV resonates with the Ti-absorption spectrum, a result which would point to a hybridization of Eu-4f with O-2p and Ti-3d states. This is rather unexpected, since 4f rare earth orbitals have usually little overlap with neighbor ions in a crystal.

In order to characterize the band structure of the LAO/ETO/STO q2DEG, in Fig. 4a we show  $k_x$ - $k_y$  in-plane cut of the Fermi Surface (FS), obtained by averaging circular minus (C-) and circular plus (C+) polarization spectra at an incoming photon energy of 465.5 eV, resonant with the  $\text{Ti}^{3+} L_2$  absorption peak. Overlapped to the map we also show Fermi surface contours associated to non-interacting electronic bands originated from atomic  $\text{Ti-3d-}t_{2g}$  states: namely, a ring-shaped feature and two ellipsoidal structures oriented lengthwise along the  $k_x$  and  $k_y$  directions. The ring-shaped Fermi contour corresponds to light effective mass electrons having mainly  $3d_{xy}$  orbital character, while the ellipsoidal ones are related to heavy effective masses electrons with mainly  $3d_{xz}/3d_{yz}$  orbital characters. The qualitative features of the measured FS in (001) LAO/ETO/STO resemble those observed on different oxide systems characterized by the presence of a q2DEG, like the (001) STO surface<sup>30,31</sup>, the (001) LAO/STO interface<sup>32-34</sup>, and the recently investigated (001) ETO surface<sup>28</sup>. However, the FS, in particular at the 2<sup>nd</sup> Brillouin zone (BZ) where heavy bands contributions are better resolved due to matrix element effects, has a shape which departs from the simple tight-binding model of three non-interacting  $t_{2g}$  bands.

In Fig. 4b, c we show band dispersion cuts at the 1<sup>st</sup> and 2<sup>nd</sup> BZ through the  $\Gamma$ -point along the  $k_x$  direction (corresponding to  $\Gamma$ -X in

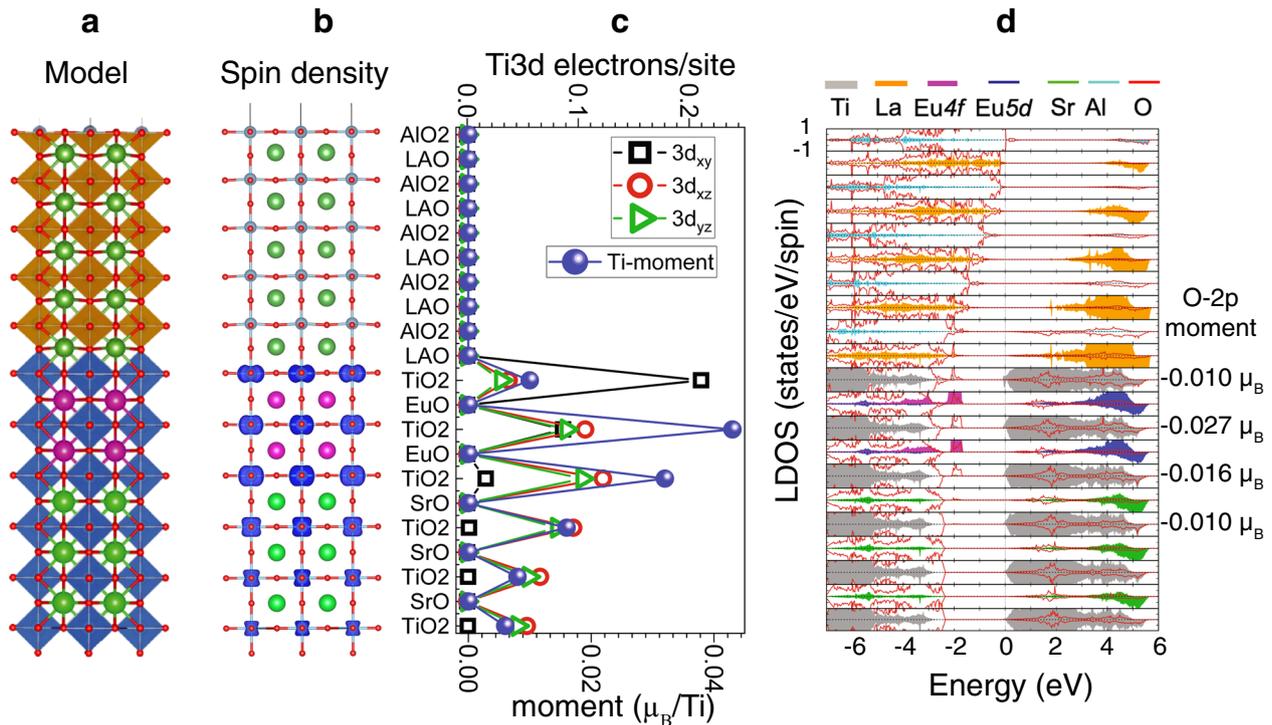
reciprocal space) with C+ polarization of the incoming photons (C+ pol,  $h\nu = 465.5$  eV). Momentum Dispersion Cuts (MDCs) at the Fermi level are shown on the top of each panel. In Fig. 4d, e we show the corresponding 2D-curvature maps to highlight the different bands<sup>35</sup>. Other data acquired with different polarizations of the incoming photons are shown in Supplementary Fig. 5.

The dispersive profile of the different bands were obtained by combining the 2D-curvature maps and the fit of the maxima of MDCs at several energies below the Fermi level, in analogy with previous studies on LAO/STO q2DEGs<sup>33,34</sup> (see Supplementary Note 3). From simple tight-binding fitting assuming three independent bands (superimposed to the 2D curvature maps of Fig. 4d, e), we estimated the effective masses  $m^*$  at the Fermi level of each band (see Table 1). We obtained  $m^* = 0.4 m_e$  for the  $3d_{xy}$  band, while for the heavy bands we estimated  $m^* = 0.25 m_e$  and  $\sim 10 m_e$  in the “light” and “heavy” directions, respectively ( $m_e$  being the free electron mass). The  $3d_{xy}$  band effective mass and the  $3d_{xz}/3d_{yz}$  bands effective masses along the “light” directions are typically smaller than the ones estimated for the LAO/STO system<sup>33</sup>. The splitting between the heavy- and light- band bottoms is of the order of  $\sim 35$  meV at the  $\Gamma$ -point, considerably smaller than that usually reported for the STO-surface 2DEGs<sup>30</sup>, but consistent with our recent report about the differences between the ETO and STO surface states<sup>28</sup>, and in quantitative agreement with x-ray linear dichroism data on (001) LAO/ETO/STO<sup>11</sup>.

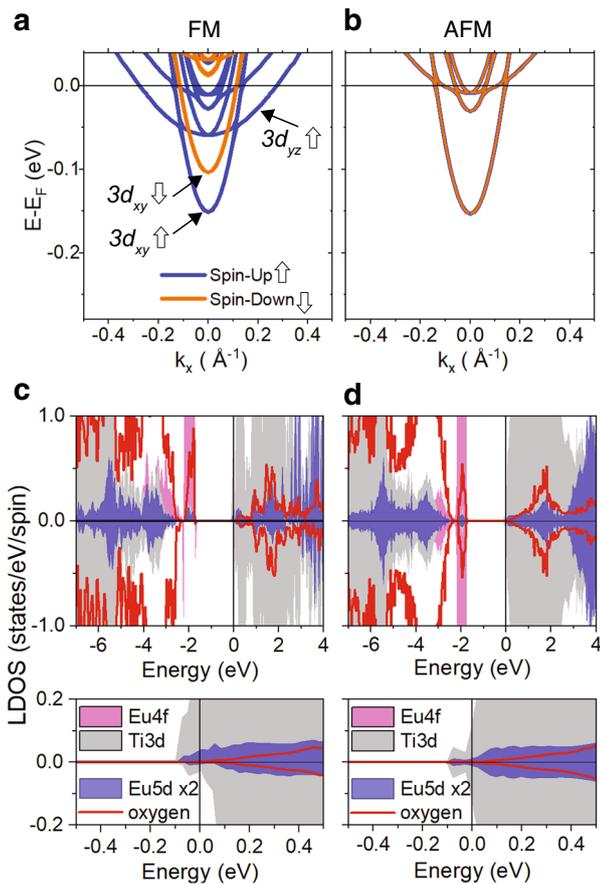
Bands	$k_F$ ( $\text{\AA}^{-1}$ )	$E(0)$ (meV)	$m^*/m_e$
$3d_{xy}$	$0.09 \pm 0.01$	$-85 \pm 5$	$0.4 \pm 0.1$
$3d_{yz}$	$0.27 \pm 0.02$	$-50 \pm 5$	$10 \pm 2$
$3d_{xz}$	$0.055 \pm 0.005$	$-50 \pm 5$	$0.25 \pm 0.05$

#### DFT + $U$ calculations

In order to understand the mechanism at the base of the formation of the q2DEG in LAO/ETO/STO and of the spin-polarization of its carriers, we performed DFT +  $U$  calculations<sup>36-42</sup> on a non-defective  $c(2 \times 2)$  LAO/ETO/STO(001) heterostructure composed of



**Fig. 5** DFT +  $U$  calculations. **a** Side view of the simulated  $c(2 \times 2)$  LAO/ETO/STO(001) heterostructure; **b** spin density, integrated between  $-0.3$  eV and  $E_F$ , with isovalue of  $0.0004 \text{ e}/\text{\AA}^3$ , showing the spatial distribution of the 2DEG and the orbital character at the Ti sites; **c** orbital and layer resolved Ti-3d occupation (upper scale; black squares,  $3d_{xy}$ ; red circles,  $3d_{xz}$ ; green triangles,  $3d_{yz}$ ) and Ti-3d magnetic moment (lower scale; blue circles) obtained by integrating the density of states between  $-0.3$  eV and  $E_F$ . **d** Layer-, spin- and element- resolved density of states. The color code for the partial LDOS contribution of different ions is indicated on top of the figure. On the right side of the figure, we also indicate the values of the calculated oxygen-2p magnetic moment in the corresponding interfacial  $\text{TiO}_2$  layers.



**Fig. 6 Spin-resolved band structure and LDOS in the FM and AFM state from the DFT +  $U$  calculations.** Spin-resolved band structure calculations of the Ti-3d bands around the  $\Gamma$ -point in a small energy range in the (a) FM (left panel) and (b) AFM (right panel) solutions of the model. Note that in the FM solution, the spin-up  $3d_{yz}$  band shift to lower energy and crosses the Fermi level, while the  $3d_{xy}$  spin-down bands shifts up, giving rise to an overall spin-polarization of the q2DEG. Element and spin-resolved density of states from the DFT +  $U$  calculations for  $\text{EuTiO}_3$  layers (average) with (c) FM and (d) AFM coupling (same vertical scales in a, b). In the FM solution we do observe a substantial spin-polarization at  $-2$  eV of both Eu-4f (magenta) and Eu-5d (blue) states, and an overlap between Ti-3d states (gray) and Eu-5d states near and above the Fermi level, absent in the AFM case.

2 uc of ETO and 5 uc of LAO on both sides of a symmetric  $\text{TiO}_2$ -terminated 6.5 uc STO slab (Fig. 5a). Along the  $c$ -axis the LAO surface is separated by a vacuum region of 20 Å from its periodic images. Calculations with larger STO stacks of 10.5 uc and 14.5 uc are shown in Supplementary Fig. 7 and give qualitatively similar results. On-site effective Hubbard parameter  $U = 4$  eV, 7.5 eV and 8 eV were applied on the Ti-3d, Eu-4f and La-4f states, respectively. Similar to findings for the ETO (001) surface state<sup>28</sup>, the choice of  $U$  for the Eu-4f states, and partially for the Ti-3d states, is dictated by the necessity to reproduce the position of the  $\text{Eu}^{2+}$  peak in the valence band at about  $-2$  eV. We performed calculations for both AFM (G-type) and FM arrangement of the Eu-ions. The FM solution is preferred by  $\sim 20$  meV per simulation cell. This result is in full agreement with the experimental evidence of a FM ordering of  $\text{Eu}^{2+}$  magnetic moments shown in Figs. 1, 2 and in ref. 11.

In Fig. 5 we show also the spatial distribution of the spin density (Fig. 5b), the layer- and orbitally- resolved Ti-3d occupation and magnetic moments (Fig. 5c), as well as the layer, atom- and spin-resolved density of states across the interface (Fig. 5d). According to the DFT +  $U$  results, the q2DEG is formed as a result of the polar

discontinuity at the LAO/ETO/STO interface and surface, evidenced by the overlap between the energy positions of the O-2p band of the  $\text{AlO}_2$  surface and of the interfacial  $\text{TiO}_2$  3d-conduction bands in Fig. 5d, analogous to (001) LAO/STO bilayers. The spin density plotted in Fig. 5b and the layer-resolved orbital occupation in Fig. 5c show that the q2DEG is characterized by a substantial occupation of Ti-3d states within the ETO layer, which goes to zero within a few unit cells of STO.

The bands are spin-polarized, with the highest polarization within the ETO layers, which exhibit also the largest electron occupation. In particular, Fig. 5b, c shows that the  $3d_{xy}$  bands are the first to be occupied at the interface and are mostly localized into the ETO film, whereas the occupied states within STO have a  $3d_{xz}/3d_{yz}$  orbital character, reflecting the wider distribution of these carriers, as also observed in (001) STO<sup>30</sup>. The Ti-3d magnetic moment, obtained by integrating the spin-resolved density of states in the range between  $-0.3$  eV and  $E_F$  (i.e., the contribution from spin-polarized conduction bands) is maximum in the second ETO layer, and then goes to zero within the first three-unit cells of STO (Fig. 5c). Consequently, both ETO and the topmost STO layers host electrons which are spin-polarized. The associated Ti-3d magnetic moment is parallel to the large spin-moment of  $\text{Eu}^{2+}$  4f states ( $6.97 \mu_B$ ), and its order of magnitude is in overall agreement with the analysis of the XMCD data (see Supplementary Note 1). The layer resolved map in Fig. 5d shows also a finite spin-polarized electron density of Ti-3d, Eu-5d and O-2p states at the position of the  $\text{Eu}^{2+}$  peak, demonstrating a hybridization between Eu-4f, Eu-5d, O-2p and Ti-3d states, in agreement with RESPAES data. Moreover, O-2p states in the  $\text{TiO}_2$  layers of ETO also acquire a finite magnetic moment, opposite to the Ti-3d<sup>1</sup> and Eu-4f<sup>7</sup> magnetic moments. This can explain the two components of the XMCD spectra, one related to the  $3d^0 + 3d^1L$  states, with a negative exchange interaction with the  $\text{Eu}^{2+}$  4f states, and the other from Ti-3d<sup>1</sup> spin-polarized electrons, parallel to the overall magnetization direction.

In order to compare the experimentally determined band structure shown in Fig. 4 with DFT +  $U$  calculations, and to highlight the differences between the conduction bands in the FM (the ground state) and AFM solutions of the model, we show in Fig. 6 the calculated, spin-resolved, Ti-3d band structure around the Fermi level in the two cases. In the case of the AFM state, we have adopted a G-type ordering of the Eu-spin moments. The calculated band structures are qualitatively consistent with the main characteristics of the experimentally determined band properties shown in Fig. 4. In both the FM (Fig. 6a, left panel) and AFM (Fig. 6b, right panel) solutions, the lowest lying Ti-3d band crossing  $E_F$  is a dispersive parabolic band with  $d_{xy}$  orbital character, which switches at larger momentum into a much flatter band indicating an avoided crossing with the heavy bands with  $3d_{xz}/3d_{yz}$  orbital character. The higher-lying bands around  $\Gamma$  are replicas from different layers due to the confinement of the q2DEG. In the FM solution of the model, spin-up (blue) and spin-down (orange) bands, both of the  $d_{xy}$ , and  $d_{xz}/d_{yz}$  main orbital character, are split in the whole energy range and cross the Fermi level, showing a prevalence of majority spin-polarized electrons in the system, in full agreement with all the experimental evidences of FM correlations in the q2DEG. In the AFM case the spin-degenerate  $3d_{xy}$  bands arrange in the same energy region as the spin-up band of the FM phase, while  $3d_{xz}/3d_{yz}$  bands are lifted above the Fermi level. Thus, the conduction band minimum, associated with the lowest  $3d_{xy}$  band is the same in both AFM and FM phases, but all the other bands are shifted in the FM solution. The DFT +  $U$  calculations in particular show a down-shift of the spin-up  $3d_{xz}/3d_{yz}$  bands and more importantly a filling of these bands taking place only in the FM ground state. Thus, ferromagnetism appears simultaneously with the orbital selective filling of  $3d_{xz}/3d_{yz}$  electrons at the Fermi level, confirming their

crucial role in establishing FM correlations in the q2DEG as earlier suggested in ref. <sup>11</sup>.

## DISCUSSION

The experimental and theoretical results on LAO/ETO/STO (001) heterostructure show an interaction between  $\text{Eu}^{2+}$ -4f and Ti-3d states. This is rather surprising considering the large value of the Hubbard parameter on the 4f orbitals, and their strong electron localization, rendering them not far from a configuration with electrons frozen in the core.

In bulk  $\text{EuTiO}_3$ , the establishment of a FM ground state, instead of an AFM one, is believed to be due to a delicate balance between the different exchange interactions among  $\text{Eu}^{2+}$  magnetic moments in the system: i) Direct FM exchange between  $\text{Eu}^{2+}$ , which is very weak due to the almost null overlap between Eu-4f orbitals; ii) an AFM super-exchange interaction mediated by O-2p states; iii) an AFM super-exchange interaction mediated by Ti-3d states; iv) a carriers mediated Ruderman-Kittel-Kasuya-Yoshida (RKKY)-like indirect FM-exchange, proposed for Layered FM ETO films<sup>17</sup>, due to an overlap between (filled)  $t_{2g}$  Ti-3d and Eu 4f-orbitals; and v) an indirect exchange interaction via Eu-5d  $t_{2g}$  states, which play an important role in Eu-chalcogenides (e.g., EuO). It has been proposed that the filling of the Ti-3d bands reverts the indirect exchange via the Eu-5d  $t_{2g}$  states from AFM to FM through the coupling between  $t_{2g}$  Eu-5d and Ti-3d electrons<sup>43,44</sup>. Our experimental XMCD data confirm that the super-exchange between  $\text{Eu}^{2+}$  magnetic moments through O-2p and Ti-3d states is negative, thus the only two mechanisms which could effectively give rise to ferromagnetism are FM-couplings mediated by itinerant carriers (without any role of Eu-5d states) or mediated by Eu-5d hybridized with Ti-3d states<sup>44</sup>.

In order to verify if the latter can effectively play a role in our heterostructures, in Fig. 6c we show the spin-resolved LDOS of the ETO layers in the FM ground state. The calculations are compared to the LDOS in the AFM configuration (Fig. 6d). We find that Eu-5d states, while in general characterized by a much weaker spectral weight than the O-2p and Ti-3d contributions, have a strong overlap with Ti-3d  $t_{2g}$  bands in the FM ground state close and above the Fermi level. Furthermore, there is also a substantial overlap with the Eu-4f state around  $-2$  eV, where these states also show a clear spin-polarization. On the other hand, in the AFM solution the Eu-5d states are at higher energy, have a lower spectral weight compared to the FM-case, and their overlap with Ti-3d states is substantially reduced. These results suggest that the FM-ordering of  $\text{Eu}^{2+}$  is mostly due to a FM interaction mediated by Eu-5d/Ti-3d hybridized states as suggested earlier<sup>44</sup>, although it is not possible to fully exclude a role of a carrier-mediated RKKY-FM interaction<sup>17</sup>.

Interestingly the calculations also demonstrate that the filling of Ti-3d<sub>xz</sub>/3d<sub>yz</sub> bands and FM-order appear concomitantly, as shown in Fig. 6a. Therefore, Ti-3d<sub>xz</sub>/3d<sub>yz</sub> electrons play an essential role in the establishing a FM-coupling among the  $\text{Eu}^{2+}$  magnetic moments in the confined heterostructure, while the 3d<sub>xy</sub> electrons seems less effective in mediating a dominant FM interaction in this system.

To summarize, in this work we combined different experimental methods and theoretical calculations to analyze and clarify the origin and properties of the spin-polarized q2DEG at the LAO/ETO/STO oxide interface. The DFT + *U* calculations show that a FM and spin-polarized q2DEG is formed at the defect-free interfaces in the (001) LAO/ETO/STO heterostructure. At the same time, theory explains some crucial properties of this q2DEG revealed by x-ray spectroscopy and electrical transport results. In particular, it is shown that the filling of 3d<sub>xz</sub>/3d<sub>yz</sub> bands, inferred from the experimental results, is simultaneous to the transition to a FM state, thus explaining why the q2DEG is spin-polarized only above the Lifshitz transition. Moreover, spin-polarized 3d<sub>xz</sub>/3d<sub>yz</sub> electrons

created at the LAO/ETO interface leak also into the first layers of STO, explaining the contribution from STO to the Ti-3d magnetic moment found in previous investigations<sup>11</sup>.

Our results show that the LAO/ETO/STO system provides a platform for the study of emerging quantum phenomena where superconductivity, magnetism and spin-orbit coupling are fully entangled, and is a clear example on how functional properties can be created in oxide 2D-systems by atomic interface engineering.

More generally, the combination of electrical transport and spectroscopy measurements with band structure calculations reported in this work is a powerful tool to obtain a deeper understanding of complex heterostructures characterized by unexpected functional properties. Such a combined approach is becoming essential to gain detailed understanding of heterostructures showing exceptional properties.

## METHODS

### Samples preparation

LAO(*n*)/ETO(2)/STO(001) heterostructures were fabricated by PLD assisted by Reflection High Energy Electron Diffraction (RHEED) from sintered  $\text{Eu}_2\text{Ti}_2\text{O}_7$  and crystalline LAO targets onto  $\text{TiO}_2$ -terminated (001) STO substrates. The samples were deposited at a temperature of 700 °C in a background  $\text{O}_2$  pressure of  $8 \times 10^{-5}$  mbar and cooled down in the same conditions to room temperature with a rate of 5 °C/min. We used an excimer laser (Lambda Physics, 248 nm wavelength) and 1.3 J/cm<sup>2</sup> fluence, 1 Hz of repetition, resulting in a rate of 0.05 uc/pulse (20 pulses for each unit cells). A q2DEG is formed when  $n > 4$  uc.

### XAS, XLD and XMCD spectroscopy measurements

We have used polarization dependent x-ray absorption spectroscopy (XAS) across the Eu  $M_{4,5}$  and the Ti- $L_{2,3}$  edge to probe directly the magnetic and orbital properties of Eu and Ti at the interface. The experiments were performed at the beamline X-Treme of the Swiss Light Source<sup>24</sup>. XAS performed with circularly or linearly polarized photons can detect the magnetic moments and the 3d-orbital energy splitting, respectively. The two techniques, usually known as x-ray magnetic circular dichroism (XMCD) and x-ray linear dichroism (XLD) are so sensitive that they can be used on single interfaces. The Eu  $M_{4,5}$  edge and the Ti  $L_{2,3}$  XMCD spectra were obtained as difference between the average of 8 and 16 (respectively) XAS spectra acquired with magnetic field parallel and antiparallel to the photon-helicity vector orientations. The 16 and 32 XAS data needed for each XMCD were collected in a sequence alternating reversal of field and polarization at each spectrum. This procedure ensures the best cancellation of spurious effects. The magnetic field dependent magnetization loops, as those shown in Fig. 1, were obtained by measuring, at each field, the difference between the TEY intensity at the  $M_5$ -Eu edge peak, normalized by the intensity below the absorption edge, obtained with two different helicities (combination with polarization and field direction).

### Squid magnetometry measurements

SQUID data were collected by using a Quantum Design MPMS3 magnetometer at the Université de Strasbourg, IPCMS. Magnetization measurements as a function of magnetic field were acquired on both LAO/ETO/STO heterostructure and bare STO substrates that suffered the same heating process mimicking the LAO/ETO growth to exclude the presence of ferromagnetic impurities stemming from the substrate itself. Data were corrected for the diamagnetism of the substrate substrating the linear contribution acquired at high magnetic fields.

### RESPES measurements

RESPES Measurements were carried out at a pressure of  $5 \times 10^{-10}$  mbar and at base temperature of  $\approx 12$  K (above the FM  $T_c$ ) on LAO(5)/ETO(2uc)/STO samples. This technique, employing soft-energy x-rays resonant with the absorption of the relevant ions in the system, enables access to buried systems, and it is, therefore, suitable for the investigation of the q2DEG at the LAO/ETO/STO interface. To reduce any influence of contaminants and to preserve as much as possible the surface of the ex-situ grown samples,

they were transferred, just after the deposition, into a sample-carrier vessel filled by inert Ar-gas. However, an exposure (limited in time) to ambient atmosphere was unavoidable for the sample mounting into the experimental station. During the measurements we moved the beam position around the region investigated in order to avoid variations of the surface and interfacial oxygen by photon irradiation, to which LAO/STO samples of certain preparation protocols have demonstrated some sensitivity at low temperature<sup>27</sup>. As a matter of fact, by raster scanning the beam on the sample, we did not detect any buildup of the  $\text{Ti}^{3+}$  spectral weight, neither in the core level x-ray photoemission (XPS) and XAS spectra, nor in the valence band (VB).

### DFT + $U$ calculations

DFT +  $U$  calculations were performed with the Vienna ab initio simulation package (VASP) (39, 40)<sup>36,37</sup> with the projector augmented wave (PAW) basis<sup>38,39</sup>. The generalized gradient approximation of Perdew, Burke and Ernzerhof<sup>40</sup> and an on-site effective Hubbard parameter<sup>41</sup>,  $U = 4$  eV, 7.5 eV and 8 eV were applied on the Ti-3d, Eu-4f and La-4f states, respectively. Analogous to the ETO (001) surface state<sup>28</sup>, the choice of  $U$  for the Eu-4f states, and partially for the Ti-3d states, is dictated by the need to reproduce the position of the experimental  $\text{Eu}^{2+}$  peak in the valence band at  $\sim 2$  eV. In general, lower values of  $U$  for Eu-4f only shift the position of the Eu-4f band to energies closer to the Fermi level. The results shown here for  $U = 4$  eV for the Ti-3d states give the best agreement with the experimental data, and in particular concerning the position of Eu-4f band. Calculations with  $U$  ranging from 1 eV to 4.5 eV show that the main conclusions are not dependent on the choice of  $U$  for Ti-3d states (see Supplementary Fig. 6).

The simulated supercell consists of a symmetric  $\text{TiO}_2$ -terminated 6.5 uc thick STO part followed by 2 uc ETO and 5 uc LAO on each side. The LAO surface is separated by 20 Å of vacuum from its periodic images to avoid interaction through the vacuum. The resulting structure contains 206 ions. Simulations were also performed with a thicker STO part, containing four and eight additional STO layers (a total of 246 and 286 ions), respectively, showing qualitatively similar results, as shown in Supplementary Note 1.

We used a  $8 \times 8 \times 1$  Monkhorst-Pack  $k$ -point grid<sup>42</sup> and a cut-off energy of 400 eV for the basis set. The ionic positions were fully relaxed until the forces were less than 0.001 eV/Å.

### DATA AVAILABILITY

All data needed to evaluate the conclusions in the paper are present in the paper and/or in the Supplementary Information. The datasets generated and/or analyzed during the current study are available from the corresponding author on reasonable request.

### CODE AVAILABILITY

The code that supports the findings of this study is available from the corresponding author upon reasonable request.

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## COMPETING INTERESTS

The authors declare no competing interests.

## ADDITIONAL INFORMATION

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