

Propagation or extinction in bistable equations: the non-monotone role of initial fragmentation

Matthieu Alfaro, François Hamel, Lionel Roques

▶ To cite this version:

Matthieu Alfaro, François Hamel, Lionel Roques. Propagation or extinction in bistable equations: the non-monotone role of initial fragmentation. Discrete and Continuous Dynamical Systems - Series S, 2024, 17 (4), pp.1460-1484. 10.3934/dcdss.2023165. hal-04055502v2

HAL Id: hal-04055502 https://cnrs.hal.science/hal-04055502v2

Submitted on 29 Jun 2024

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers. L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

Propagation or extinction in bistable equations: the non-monotone role of initial fragmentation

Matthieu Alfaro ^a, François Hamel ^b and Lionel Roques ^c *

^a Univ Rouen Normandie, LMRS, CNRS, Rouen, France
 ^b Aix Marseille Univ, CNRS, I2M, Marseille, France
 ^c INRAE, BioSP, 84914, Avignon, France

To Professor Yihong Du, a distinguished scholar and esteemed mathematician

Abstract

In this paper, we investigate the large-time behavior of bounded solutions of the Cauchy problem for a reaction-diffusion equation in \mathbb{R}^N with bistable reaction term. We consider initial conditions that are chiefly indicator functions of bounded Borel sets. We examine how geometric transformations of the supports of these initial conditions affect the propagation or extinction of the solutions at large time. We also consider two fragmentation indices defined in the set of bounded Borel sets and we establish some propagation or extinction results when the initial supports are weakly or highly fragmented. Lastly, we show that the large-time dynamics of the solutions is not monotone with respect to the considered fragmentation indices, even for equimeasurable sets.

Keywords: Reaction-diffusion equations; invasion; extinction; fragmentation.

1 Introduction

This paper is concerned with the large-time behavior of solutions to the Cauchy problem for the bistable reaction-diffusion equation

$$\begin{cases}
\frac{\partial u}{\partial t} = \Delta u + f(u), & t > 0, \ x \in \mathbb{R}^N, \\
u(0, x) = u_0(x), & x \in \mathbb{R}^N,
\end{cases}$$
(1.1)

in any dimension $N \geq 1$, where Δ stands for the Laplacian with respect to the spatial variables $x \in \mathbb{R}^N$. The function $f: [0,1] \to \mathbb{R}$ is assumed to be of class C^1 and of the

^{*}This work has received funding from Excellence Initiative of Aix-Marseille Université - A*MIDEX, a French "Investissements d'Avenir" programme, from the French ANR RESISTE (ANR-18-CE45-0019), ReaCh (ANR-23-CE40-0023-02) and DEEV (ANR-20-CE40-0011-01) projects.

bistable type with positive mass, that is,

$$\begin{cases}
f(0) = f(1) = 0, & f'(0) < 0, & f'(1) < 0, \int_{0}^{1} f(s)ds > 0, \\
\exists \theta \in (0, 1), & f < 0 \text{ in } (0, \theta), & f > 0 \text{ in } (\theta, 1), & f'(\theta) > 0.
\end{cases}$$
(1.2)

The initial conditions $u_0: \mathbb{R} \to [0,1]$ are assumed to be Lebesgue-measurable and compactly supported. The Cauchy problem (1.1) is well posed and, for each such u_0 , the solution $u: [0, +\infty) \times \mathbb{R}^N \to [0, 1]$ exists and is unique, it is classical in $(0, +\infty) \times \mathbb{R}^N$ and $u(t, \cdot) \to u_0$ as $t \to 0^+$ in $L^1(\mathbb{R}^N)$. Furthermore, from the strong parabolic maximum principle, one has 0 < u(t, x) < 1 for all $(t, x) \in (0, +\infty) \times \mathbb{R}^N$ provided that $||u_0||_{L^1(\mathbb{R}^N)} > 0$. In biological or ecological models, the quantity u stands for the normalized concentration of a species, subject to local dispersion on the one hand, and on growth and death processes taking into account a strong Allee effect on the other hand, meaning that the per capita growth rate f(u)/u (hence f(u) itself) is negative at low densities.

Some notations

Throughout the paper, " $\| \|_{\infty}$ ", " $\| \|_{\infty}$ ", " $\| \|_{\infty}$ " denote respectively the maximum norm, the Euclidean norm and the Euclidean inner product in \mathbb{R}^N , $B_r(x) := \{y \in \mathbb{R}^N : |y - x| < r\}$ denotes the open Euclidean ball of center $x \in \mathbb{R}^N$ and radius $r \geq 0$, and $B_r := B_r(0)$. We call

$$\mathcal{B} := \{ B_r(x) : x \in \mathbb{R}^N, \ r > 0 \}.$$

We denote λ the Lebesgue measure in \mathbb{R}^N , and $\omega_N := \lambda(B_1)$. For $x \in \mathbb{R}^N$ and a Borel subset E of \mathbb{R}^N , we define the essential distance between x and E as

$$\operatorname{dist}(x, E) := \operatorname{essinf}\left(y \in E \mapsto |x - y|\right) = \sup\left\{r \ge 0 : \lambda(B_r(x) \cap E) = 0\right\}$$

(hence $\operatorname{dist}(x, E) = +\infty$ if $\lambda(E) = 0$, and the supremum is a maximum if $\lambda(E) > 0$). The function $x \mapsto \operatorname{dist}(x, E)$ is then Lipschitz continuous in \mathbb{R}^N . The d_1 -distance between two bounded Borel subsets E and F of \mathbb{R}^N is given by the Lebesgue measure of the symmetric difference $E\Delta F$, that is,

$$d_1(E,F) := \lambda ((E \cup F) \setminus (E \cap F)) = \lambda(E \cup F) - \lambda(E \cap F), \tag{1.3}$$

which is also the L^1 norm of the difference between the indicator functions of the sets E and F. The essential Hausdorff distance between two bounded Borel subsets E and F of \mathbb{R}^N is given by

$$d_{H}(E, F) := \max \Big(\| \operatorname{dist}(\cdot, F) \|_{L^{\infty}(E)}, \| \operatorname{dist}(\cdot, E) \|_{L^{\infty}(F)} \Big), \tag{1.4}$$

with the conventions $d_H(E,F) = d_H(F,E) = +\infty$ if $\lambda(E) > 0 = \lambda(F)$, and $d_H(E,F) = 0$ if $\lambda(E) = \lambda(F) = 0$. Throughout the paper, for two Borel subsets E and F of \mathbb{R}^N , we also say that E is included in F up to a negligible set, or equivalently F contains E up to a negligible set, if $\lambda(E \setminus F) = 0$. We say that E and F are equal up to a negligible set if $E \subset F$ up to a negligible set and $F \subset E$ up to a negligible set, that is, the indicator functions of the sets E and F are equal almost everywhere in \mathbb{R}^N . In other words, for any two bounded Borel subsets E and F of \mathbb{R}^N , $d_1(E,F) = 0$ if and only if E and E are equal up to a negligible set. Lastly, for a bounded Borel set $E \subset \mathbb{R}^N$ and for E0, we call

$$\mathbb{B}_r(E) := \{ F \text{ bounded Borel subsets of } \mathbb{R}^N : d_1(E, F) < r \}.$$
 (1.5)

Some results of the literature

The main goal of the paper is to understand the role of the fragmentation, in a sense to be made precise, of the initial condition u_0 on the large-time dynamics of the solution u. But before presenting various notions of fragmentation and their properties and effects on the dynamics of (1.1), let us recall some important known results of the literature.

First of all, it turns out that, for each given compactly supported initial condition $u_0: \mathbb{R}^N \to [0,1]$, u can have only three types of asymptotic behaviors as $t \to +\infty$. Namely, it follows from [33] that either $u(t,\cdot) \to 0$ as $t \to +\infty$ uniformly in \mathbb{R}^N (the extinction case), or $u(t,\cdot) \to 1$ as $t \to +\infty$ locally uniformly in \mathbb{R}^N (the invasion case), or $u(t,\cdot) \to \Phi(\cdot + x_0)$ as $t \to +\infty$ uniformly in \mathbb{R}^N for some $x_0 \in \mathbb{R}^N$, where $\Phi: \mathbb{R}^N \to (0,1)$ is the unique stationary solution of (1.1) converging to 0 at infinity and such that $\max_{\mathbb{R}^N} \Phi = \Phi(0)$ (the function Φ is actually radially symmetric and decreasing with respect to the origin). We point out that, if f satisfies (1.2) but with a nonpositive integral over [0,1], then only the extinction case is possible. Under the full assumption (1.2), when $u(t,\cdot) \to 1$ as $t \to +\infty$ locally uniformly in \mathbb{R}^N , then it is also known from [10] (in dimension N = 1) and from [45] (in any dimension $N \geq 1$) that

$$\liminf_{t\to +\infty} \Big(\min_{|x|\leq ct-\frac{N-1}{c}\ln t-A} u(t,x) \Big) \underset{A\to +\infty}{\longrightarrow} 1 \text{ and } \limsup_{t\to +\infty} \Big(\max_{|x|\geq ct-\frac{N-1}{c}\ln t+A} u(t,x) \Big) \underset{A\to +\infty}{\longrightarrow} 0,$$

where c>0 is the unique speed of a traveling front $\varphi(x-ct)$ solving (1.1) in dimension N=1 and connecting 1 to 0, that is, with $\varphi:\mathbb{R}\to(0,1)$ such that $\varphi(-\infty)=1$ and $\varphi(+\infty)=0$. In other words, the levels sets of u with a given level $\rho\in(0,1)$ spread with speed c in all directions as $t\to+\infty$ (this spreading result was originally proved in [3]) and are even located at bounded Hausdorff distance from spheres of radii $ct-((N-1)/c)\ln t$. As a matter of fact, the level sets become asymptotically locally planar as $t\to+\infty$, since $\sup_{t\geq A,\,|x|\geq A}\nabla u(t,x)/|\nabla u(t,x)|+x/|x|\to 0$ as $A\to+\infty$ by [18] (see [16, 17] for a study of further flattening properties), but they nevertheless do not necessarily converge to families of spheres, see [40, 41, 43]. For further spreading properties and estimates of the location of the level sets for various types of functions f, we refer to e.g. [7, 9, 13, 37] in the case of compactly supported initial conditions, and to [5, 11, 14, 20, 22, 32, 36, 44] for initial conditions with unbounded initial support in dimension N=1 and [15, 16, 17, 24, 25, 35, 38, 39] in any dimension.

Moreover, if for the problem (1.1)-(1.2) one considers a family $[0, +\infty) \ni \sigma \mapsto u_0^{\sigma}$ of compactly supported initial conditions, which is continuous and increasing in the $L^1(\mathbb{R}^N)$ sense and which is such that $u_0^0 = 0$, there is a unique threshold $\sigma^* \in (0, +\infty]$ such that the solutions u^{σ} of (1.1) with initial conditions u_0^{σ} satisfy: 1) $u^{\sigma}(t, \cdot) \to 0$ as $t \to +\infty$ uniformly in \mathbb{R}^N if $0 \le \sigma < \sigma^*$, 2) $u^{\sigma}(t, \cdot) \to 1$ as $t \to +\infty$ locally uniformly in \mathbb{R}^N if $\sigma^* < +\infty$ and $\sigma > \sigma^*$, 3) $u^{\sigma^*}(t, \cdot) \to \Phi(\cdot + x_0)$ as $t \to +\infty$ uniformly in \mathbb{R}^N for some $x_0 \in \mathbb{R}^N$ if $\sigma^* < +\infty$, see [33]. In particular, for any given $\alpha \in (\theta, 1]$, since the solutions $u^{\alpha, R}$ of (1.1) with initial conditions

$$u_0^{\alpha,R} = \alpha \, \mathbb{1}_{B_R} := \left\{ \begin{array}{ll} \alpha & \text{if } |x| < R, \\ 0 & \text{if } |x| \ge R, \end{array} \right.$$

converge to 1 as $t \to +\infty$ locally uniformly in \mathbb{R}^N provided R > 0 is large enough (see

e.g. [3, 8]), there is then from [30] or [33] a unique $R_{\alpha} \in (0, +\infty)$ such that

$$\begin{cases} u^{\alpha,R} \to 0 \text{ as } t \to +\infty \text{ uniformly in } \mathbb{R}^N & \text{if } 0 \le R < R_{\alpha}, \\ u^{\alpha,R} \to 1 \text{ as } t \to +\infty \text{ locally uniformly in } \mathbb{R}^N & \text{if } R > R_{\alpha}, \\ u^{\alpha,R} \to \Phi \text{ as } t \to +\infty \text{ uniformly in } \mathbb{R}^N & \text{if } R = R_{\alpha}. \end{cases}$$

$$(1.6)$$

The asymptotic behavior of R_{α} as $\alpha \xrightarrow{>} \theta$ was estimated in [1]. The first sharp threshold results of that type were obtained in dimension N=1, in [46] with initial conditions that are indicator functions of bounded intervals, and then in [6] for more general initial conditions. In both papers [6, 46], more general functions f can be considered (such as ignition nonlinearities for which there is $\theta \in (0,1)$ such that f=0 in $[0,\theta] \cup \{1\}$, f>0 in $(\theta,1)$ and f is non-decreasing in a neighborhood of θ). We point out that reference [6] also contains a general convergence result for the bounded nonnegative solutions of (1.1) with compactly supported initial conditions in dimension N=1 under the sole assumption f(0)=0: namely such solutions converge as $t \to +\infty$ locally in \mathbb{R} to a stationary solution, which is either constant or even and decreasing with respect to a point. Other sharp threshold results were obtained in [26, 33] for bistable-type autonomous or non-autonomous equations in \mathbb{R} or \mathbb{R}^N with initial conditions converging to 0 at infinity, and in [29, 30] for more general functions f and radially symmetric non-increasing initial conditions in $L^2(\mathbb{R})$ and $L^2(\mathbb{R}^N)$. Earlier references [3, 19, 23] also established non-sharp extinction/invasion results with respect to the size or the amplitude of the initial condition u_0 in \mathbb{R}^N for various reaction terms f. We refer to [6, 8, 26, 27, 34] for further results on the convergence to a stationary solution or the convergence to the set of stationary solutions (quasiconvergence) for various equations of the type (1.1). Extinction and invasion results have also been derived for equations of the type (1.1) in general unbounded domains Ω , see [4, 42]. We lastly mention [2, 28, 31] some results on the optimization of $\int_{\Omega} u(T,\cdot)$ (or other integral quantities) with respect to the initial condition under some pointwise and integral constraints, for some reaction-diffusion equations set in bounded domains Ω with Neumann boundary conditions on $\partial\Omega$.

Main goal and outline of the paper

In this paper, we consider the Cauchy problem (1.1) with a bistable function f of the type (1.2), and with initial conditions which are mostly indicator functions of bounded Borel sets E, that is,

$$u_0(x) = \mathbb{1}_E(x) := \begin{cases} 1 & \text{if } x \in E, \\ 0 & \text{if } x \in \mathbb{R}^N \setminus E. \end{cases}$$
 (1.7)

Our goal is to understand the effect of the fragmentation of the initial set E on the large-time dynamics of the solution of (1.1) with initial condition (1.7). Unlike the aforementioned sharp threshold results obtained for monotone families of initial conditions, that would correspond in (1.7) to monotone (with respect to the inclusion) families of sets E, we typically consider here some Borel sets E having the same Lebesgue measure (hence being in general not comparable with respect to the inclusion) and we look for some properties of these sets which guarantee the extinction or the invasion of the solutions u. As we will see, the large-time dynamics will strongly depend on the fragmentation of the sets E. Two fragmentation indices are considered and their properties are discussed in Section 2. The main results on the role of the fragmentation of E on the large-time dynamics of the

solutions of (1.1) and (1.7) are stated in Section 3. We especially show that there is no monotonicity of the large-time dynamics with respect to the fragmentation indices in the class of equimeasurable sets. The main proofs are given in Section 4.

2 Fragmentation indices and their properties

For the solutions u of problem (1.1)-(1.2) with initial conditions of the type (1.7) with $E = B_R$, there is a unique $R_1 \in (0, +\infty)$ such that extinction holds as $t \to +\infty$ if $R < R_1$ while invasion happens if $R > R_1$, as recalled in Section 1, see (1.6). Now, for a set E with Lebesgue measure $\lambda(E)$ larger than $\lambda(B_{R_1})$, can one provide some conditions guaranteeing the extinction or the invasion? The invasion obviously holds if E itself is a ball or if E contains a ball of radius larger than R_1 , from the maximum principle. But the invasion does not hold in general. For instance, in dimension N = 1, it follows from [12, Theorem 2] that, if

$$D_a := (-a - r, -a) \cup (a, a + r)$$

with a given $r \in (0, 2R_1)$, then the solutions u^a of (1.1) with initial conditions $\mathbb{1}_{D_a}$ go to extinction as $t \to +\infty$ for all a > 0 large enough, while $\lambda(D_a) = 2r > 2R_1 = \lambda(B_{R_1})$ if r is chosen such that $r > R_1$.

In order to quantify how a bounded non-negligible Borel set E deviates from the set of balls having the same measure as E, we consider in this paper two different fragmentation indices. We then list their main properties and compare them. For such a set E, we define

$$R_E := \left(\frac{\lambda(E)}{\omega_N}\right)^{1/N} = \left(\frac{\lambda(E)}{\lambda(B_1)}\right)^{1/N}$$

which is nothing but the radius of all equimeasurable balls. We then denote ρ_E the smallest radius of a ball containing E up to a negligible set, which can also be expressed as

$$\rho_E = \min_{x \in \mathbb{R}^N} \|d(x, \cdot)\|_{L^{\infty}(E)}. \tag{2.1}$$

Notice that

$$\rho_E \ge \frac{\operatorname{diam}(E)}{2} \tag{2.2}$$

for any $N \geq 1$, with equality when N = 1, where

$$\operatorname{diam}(E) := \| (x, y) \mapsto |x - y| \|_{L^{\infty}(E \times E)}$$

is the essential diameter of E. There is actually a unique x_E such that $E \subset B_{\rho_E}(x_E)$ up to a negligible set. Lastly,

$$R_E < \rho_E$$

and the equality holds if and only if E is a ball, up to a negligible set.

The fragmentation index δ_1

The first considered fragmentation index of a bounded non-negligible Borel set $E \subset \mathbb{R}^N$ is, up to the multiplicative factor 1/2, the Fraenkel asymmetry. It is based on the L^1 distance between E and the equimeasurable balls:

$$\delta_1(E) := \inf_{B \in \mathcal{B}, \ \lambda(B) = \lambda(E)} \frac{\|\mathbb{1}_B - \mathbb{1}_E\|_{L^1(\mathbb{R}^N)}}{2\lambda(E)}.$$

Since $\|\mathbb{1}_B - \mathbb{1}_E\|_{L^1(\mathbb{R}^N)} = \lambda(B\Delta E) = 2\lambda(E) - 2\lambda(E \cap B_{R_E}(x)) = 2\lambda(E \setminus B_{R_E}(x))$ for every ball $B = B_{R_E}(x)$ (namely, such that $\lambda(B) = \lambda(E)$), it follows from the continuity of the maps $x \mapsto \lambda(E \setminus B_{R_E}(x))$ and $x \mapsto \lambda(E \cap B_{R_E}(x))$ in \mathbb{R}^N that

$$\delta_{1}(E) = \min_{x \in \mathbb{R}^{N}} \frac{\lambda(E \setminus B_{R_{E}}(x))}{\lambda(E)} = \min_{x \in \mathbb{R}^{N}} \frac{\lambda(E) - \lambda(E \cap B_{R_{E}}(x))}{\lambda(E)}$$
$$= 1 - \max_{x \in \mathbb{R}^{N}} \frac{\lambda(E \cap B_{R_{E}}(x))}{\lambda(E)}.$$
 (2.3)

In particular,

$$0 \le \delta_1(E) < 1$$
,

and $\delta(E) = 0$ if and only if E is a ball, up to a negligible set. Furthermore, the constant 1 in the inequality $\delta_1(E) < 1$ is optimal: for instance, the non-negligible Borel sets

$$E_n := \bigcup_{x \in \mathbb{Z}^N \cap (0,n)^N} B_{1/n^2} \left(\frac{x}{n}\right), \quad n \ge 1, \tag{2.4}$$

which are all included into the cube $(0,1)^N$, are highly fragmented for the index δ_1 as $n \to +\infty$, in the sense that

$$\delta_1(E_n) \to 1 \text{ as } n \to +\infty.$$
 (2.5)

Lastly, for any $\nu \in (0,1)$, consider the non-negligible Borel sets

$$F_n := \bigcup_{x \in \mathbb{Z}^N \cap B_n} \left[\frac{x}{n} + \left(-\frac{\nu}{2n}, \frac{\nu}{2n} \right)^N \right] = \bigcup_{y \in (\mathbb{Z}^N/n) \cap B_1} \left[y + \left(-\frac{\nu}{2n}, \frac{\nu}{2n} \right)^N \right], \quad n \ge 1, \qquad (2.6)$$

which are all included into the ball $B_{1+\nu\sqrt{N}/2}$. Actually, counting the number A_n of points $x \in \mathbb{Z}^N \cap B_n$ corresponds to the N-dimensional extension of Gauss circle problem. It is a difficult issue, but some estimates can be obtained when n is large. In particular, $A_n = n^N \omega_N + O(n^{N-1})$ as $n \to +\infty$, see [21, Chapter 4] for sharper estimates. Thus, $\lambda(F_n) = A_n (\nu/n)^N \to \nu^N \omega_N$ as $n \to +\infty$. The same arguments imply that $\lambda(F_n \cap B_\nu) \to \nu^{2N} \omega_N$ as $n \to +\infty$, so that

$$\delta_1(F_n) \to 1 - \nu^N \text{ as } n \to +\infty.$$

The fragmentation index δ_H

The second considered fragmentation index of a bounded non-negligible Borel set $E \subset \mathbb{R}^N$ is based on the essential Hausdorff distance between E and the equimeasurable balls:

$$\delta_H(E) := \inf_{B \in \mathcal{B}, \, \lambda(B) = \lambda(E)} \frac{d_H(E, B)}{\rho_E + R_E}.$$

From the continuity of the map $x \mapsto d_H(E, B_{R_E}(x))$ in \mathbb{R}^N , one can also write

$$\delta_H(E) = \min_{x \in \mathbb{R}^N} \frac{d_H(E, B_{R_E}(x))}{\rho_E + R_E}.$$
(2.7)

Furthermore, for the unique $x_E \in \mathbb{R}^N$ such that $E \subset B_{\rho_E}(x_E)$ up to a negligible set, one has $\|\operatorname{dist}(\cdot, B_{R_E}(x_E))\|_{L^{\infty}(E)} \leq \rho_E - R_E < \rho_E + R_E$. Moreover,

 $\|\operatorname{dist}(\cdot, E)\|_{L^{\infty}(B_{R_E}(x_E))} < \rho_E + R_E \text{ (indeed, otherwise, since the map } \operatorname{dist}(\cdot, E) \text{ is continuous in } \mathbb{R}^N, \text{ there would exist a point } y_E \in \overline{B_{R_E}(x_E)} \text{ such that } \operatorname{dist}(y_E, E) \geq \rho_E + R_E, \text{ hence } E \subset \mathbb{R}^N \setminus B_{\rho_E + R_E}(y_E) \text{ up to a negligible set, and then } E \subset \mathbb{R}^N \setminus B_{\rho_E}(x_E) \text{ up to a negligible set from the triangle inequality, a contradiction with the positivity of } \lambda(E) \text{ and the inclusion } E \subset B_{\rho_E}(x_E) \text{ up to a negligible set)}. \text{ Finally, the inequality } \|\operatorname{dist}(\cdot, E)\|_{L^{\infty}(B_{R_E}(x_E))} < \rho_E + R_E \text{ has been proved, hence } d_H(E, B_{R_E}(x_E)) < \rho_E + R_E \text{ and } d_H(E, B_{R_E}(x_E)) < \rho_E + R_E \text{ and } d_H(E, B_{R_E}(x_E)) < \rho_E + R_E \text{ and } d_H(E, B_{R_E}(x_E)) < \rho_E + R_E \text{ and } d_H(E, B_{R_E}(x_E)) < \rho_E + R_E \text{ and } d_H(E, B_{R_E}(x_E)) < \rho_E + R_E \text{ and } d_H(E, B_{R_E}(x_E)) < \rho_E + R_E \text{ and } d_H(E, B_{R_E}(x_E)) < \rho_E + R_E \text{ and } d_H(E, B_{R_E}(x_E)) < \rho_E + R_E \text{ and } d_H(E, B_{R_E}(x_E)) < \rho_E + R_E \text{ and } d_H(E, B_{R_E}(x_E)) < \rho_E + R_E \text{ and } d_H(E, B_{R_E}(x_E)) < \rho_E + R_E \text{ and } d_H(E, B_{R_E}(x_E)) < \rho_E + R_E \text{ and } d_H(E, B_{R_E}(x_E)) < \rho_E + R_E \text{ and } d_H(E, B_{R_E}(x_E)) < \rho_E + R_E \text{ and } d_H(E, B_{R_E}(x_E)) < \rho_E + R_E \text{ and } d_H(E, B_{R_E}(x_E)) < \rho_E + R_E \text{ and } d_H(E, B_{R_E}(x_E)) < \rho_E + R_E \text{ and } d_H(E, B_{R_E}(x_E)) < \rho_E + R_E \text{ and } d_H(E, B_{R_E}(x_E)) < \rho_E + R_E \text{ and } d_H(E, B_{R_E}(x_E)) < \rho_E + R_E \text{ and } d_H(E, B_{R_E}(x_E)) < \rho_E + R_E \text{ and } d_H(E, B_{R_E}(x_E)) < \rho_E + R_E \text{ and } d_H(E, B_{R_E}(x_E)) < \rho_E + R_E \text{ and } d_H(E, B_{R_E}(x_E)) < \rho_E + R_E \text{ and } d_H(E, B_{R_E}(x_E)) < \rho_E + R_E \text{ and } d_H(E, B_{R_E}(x_E)) < \rho_E + R_E \text{ and } d_H(E, B_{R_E}(x_E)) < \rho_E + R_E \text{ and } d_H(E, B_{R_E}(x_E)) < \rho_E + R_E \text{ and } d_H(E, B_{R_E}(x_E)) < \rho_E + R_E \text{ and } d_H(E, B_{R_E}(x_E)) < \rho_E + R_E \text{ and } d_H(E, B_{R_E}(x_E)) < \rho_E + R_E \text{ and } d_H(E, B_{R_E}(x_E)) < \rho_E + R_E \text{ and } d_H(E, B_{R_E}(x_E)) < \rho_E + R_E \text{ and } d$

$$0 \le \delta_H(E) < 1.$$

As for the index δ_1 , a bounded non-negligible Borel E satisfies $\delta_H(E) = 0$ if and only if E is a ball, up to a negligible set. Furthermore, the constant 1 in the inequality $\delta_H(E) < 1$ is optimal: for instance, the sets E_n defined in (2.4) satisfy $\delta_H(E_n) \to 1$ as $n \to +\infty$. Similarly, taking two points $x \neq y$ in \mathbb{R}^N , the non-negligible Borel sets

$$G_n := B_{1/n}(x) \cup B_{1/n}(y), \quad n \ge 1,$$
 (2.8)

which are all included in the ball $B_{|x-y|/2+1}((x+y)/2)$, are highly fragmented for the index δ_H as $n \to +\infty$, in the sense that

$$\delta_H(G_n) \to 1 \text{ as } n \to +\infty.$$

Observe on the other hand that these sets are not highly fragmented for the index δ_1 , since $\delta_1(G_n) \to 1/2$ as $n \to +\infty$. Lastly, for $\nu \in (0,1)$, the sets F_n defined by (2.6) are such that

$$\delta_H(F_n) \to \frac{1-\nu}{1+\nu}$$
 as $n \to +\infty$.

In dimension N=1, for a bounded non-negligible Borel set E, calling

$$m_E := \operatorname{essinf} E$$
 and $M_E := \operatorname{esssup} E$,

one has $m_E < M_E \in \mathbb{R}$, $x_E = (m_E + M_E)/2$, and $\operatorname{diam}(E) = M_E - m_E = 2\rho_E \ge 2R_E$. Moreover, $\|\operatorname{dist}(\cdot, (x - R_E, x + R_E))\|_{L^{\infty}(E)} \ge \rho_E - R_E$ for every $x \in \mathbb{R}$, with equality if and only if $x = x_E$, while $\|\operatorname{dist}(\cdot, E)\|_{L^{\infty}(x_E - R_E, x_E + R_E)} \le \rho_E - R_E$ (otherwise, there would be $y_E \in [x_E - R_E, x_E + R_E]$ such that $d := \operatorname{dist}(y_E, E) > \rho_E - R_E$ and $E \subset (x_E - \rho_E, x_E + \rho_E) \setminus (y_E - d, y_E + d)$ up to a negligible set, a contradiction with $|y_E - x_E| \le R_E$ and $\lambda(E) = 2R_E$). Finally, $d_H(E, (x - R_E, x + R_E)) \ge \rho_E - R_E$ for each $x \in \mathbb{R}$, with equality if and only if $x = x_E$, hence

$$\delta_H(E) = \frac{\rho_E - R_E}{\rho_E + R_E}$$
 in dimension $N = 1$.

Furthermore, in any dimension $N \ge 1$, (2.1) implies that $||d(x,\cdot)||_{L^{\infty}(E)} \ge \rho_E$ for any $x \in \mathbb{R}^N$, hence $d_H(E, B_{R_E}(x)) \ge ||d(\cdot, B_{R_E}(x))||_{L^{\infty}(E)} \ge \rho_E - R_E$ and

$$\delta_H(E) \ge \frac{\rho_E - R_E}{\rho_E + R_E}$$
 in any dimension $N \ge 1$, (2.9)

by (2.7). Notice that, unlike the case N=1, the inequality (2.9) is in general strict in dimensions $N \geq 2$. For instance, with $N \geq 2$, pick any a > 0 and consider the spherical shell

$$E := B_{(a^N+1)^{1/N}} \setminus B_a.$$

For this set, one has $R_E = 1$, $\rho_E = (a^N + 1)^{1/N}$, and, since $a > (a^N + 1)^{1/N} - 1$ (because $N \ge 2$), it follows that

$$\min_{x \in \mathbb{R}^N} d_H(E, B_1(x)) = d_H(E, B_1) = a > (a^N + 1)^{1/N} - 1 = \rho_E - R_E,$$

hence $\delta_H(E) = a/(\rho_E + R_E) > (\rho_E - R_E)/(\rho_E + R_E)$.

An important consequence of (2.9) is that $\delta_H(E) \to 1$ if $\rho_E/R_E \to +\infty$, as for the sets E_n and G_n defined in (2.4) and (2.8), for which $\rho_{E_n}/R_{E_n} \sim n\sqrt{N}/2$ and $\rho_{G_n}/R_{G_n} \sim n|x-y|/2^{1+1/N}$ as $n \to +\infty$.

Some properties of these indices

Let us now list some further properties satisfied by the fragmentation indices δ_1 and δ_H . First of all, they are clearly invariant by rigid motion. They are also invariant by contraction or dilation, that is,

$$\delta_1(\mu E) = \delta_1(E)$$
 and $\delta_H(\mu E) = \delta_H(E)$

for every bounded non-negligible Borel set E and for every $\mu > 0$.

Furthermore, the comparisons between δ_1 and δ_H , respectively between $1-\delta_1$ and $1-\delta_H$, are summarized in the following proposition. In short, weakly fragmented sets for the index δ_H are also weakly fragmented for δ_1 , and highly fragmented sets for the index δ_1 are also highly fragmented for δ_H , whereas the reverse comparisons are false.

Proposition 2.1. (i) There is a constant $\gamma > 0$ such that

$$0 \le \delta_1(E) \le \gamma \, \delta_H(E) \tag{2.10}$$

for every bounded non-negligible Borel set E. In particular, $\delta_1(E) \to 0$ as $\delta_H(E) \to 0$.

(ii) On the other hand,

$$\delta_H(E) \not\to 0$$
 as $\delta_1(E) \to 0$.

(iii) There is a constant $\eta > 0$ such that

$$0 < 1 - \delta_H(E) \le \eta (1 - \delta_1(E))^{1/N} \tag{2.11}$$

for every bounded non-negligible Borel set E. In particular, $\delta_H(E) \to 1$ as $\delta_1(E) \to 1$.

(iv) On the other hand,

$$1 - \delta_1(E) \not\to 0$$
 as $1 - \delta_H(E) \to 0$.

Part (iv) is immediate since the sets G_n defined in (2.8) are such that $\delta_H(G_n) \to 1$ as $n \to +\infty$ and $\delta_1(G_n) \to 1/2 \neq 1$ as $n \to +\infty$. For (ii), consider a point $x \in \mathbb{R}^N$ with |x| > 1 and the sets $H_n := B_1 \cup B_{1/n}(x)$ for $n \geq 1$. There holds $\delta_1(H_n) \to 0$ and $\delta_H(H_n) \to (|x| - 1)/(|x| + 3) > 0$ as $n \to +\infty$, yielding the desired conclusion. The proofs of parts (i) and (iii) are not as immediate, and are done in Subsection 4.1.

3 Influence of initial fragmentation on the large-time behavior

Let us come back to the Cauchy problem (1.1) with initial conditions $u_0 = \mathbb{1}_E$ as in (1.7), where $E \subset \mathbb{R}^N$ always stands for a bounded non-negligible Borel set and f is a bistable-type function satisfying (1.2). Remember that 0 < u(t,x) < 1 for all $(t,x) \in (0,+\infty) \times \mathbb{R}^N$ for any such set E, and let us denote

$$\left\{ \begin{array}{l} \mathcal{I} := \left\{ E : u(t, \cdot) \to 1 \text{ as } t \to +\infty \text{ locally uniformly in } \mathbb{R}^N \right\}, \\ \mathcal{E} := \left\{ E : u(t, \cdot) \to 0 \text{ as } t \to +\infty \text{ uniformly in } \mathbb{R}^N \right\}, \\ \mathcal{T} := \left\{ E : \exists \, a \in \mathbb{R}^N, \ u(t, \cdot) \to \Phi(\cdot + a) \text{ as } t \to +\infty \text{ uniformly in } \mathbb{R}^N \right\}, \end{array} \right.$$

where $\Phi: \mathbb{R}^N \to (0,1)$ is the unique radially symmetric stationary solution of (1.1) such that $\Phi(x) \to 0$ as $|x| \to +\infty$. From the results of [33] recalled in Section 1, any bounded non-negligible Borel set E belongs to $\mathcal{I} \cup \mathcal{E} \cup \mathcal{T}$. The calligraphic letters \mathcal{I} , \mathcal{E} and \mathcal{T} respectively stand for invasion, extinction and threshold.

Our goal in this section is to determine sufficient conditions for E to belong to \mathcal{I} or \mathcal{E} , and to compare them. Notice at once from the comparison principle that if $E \subset F$ and if $F \in \mathcal{E}$ (resp. if $E \in \mathcal{I}$), then $E \in \mathcal{E}$ (resp. $F \in \mathcal{I}$). Actually, for monotone (for the inclusion) continuously increasing (for the Lebesgue measure) families of sets E, these sets belong to either $\mathcal{I} \cup \mathcal{E}$ up to at most one threshold value, from [33], and this is why we mainly focus on the conditions for which $E \in \mathcal{I}$ or $E \in \mathcal{E}$. We however do not consider here monotone families of sets E and we rather look for some conditions involving the fragmentation indices $\delta_1(E)$ and $\delta_H(E)$, for constant values of the Lebesgue measure $\lambda(E)$.

First of all, it turns out that the sets \mathcal{I} and \mathcal{E} are open for the topology generated by the balls $\mathbb{B}_r(E)$ defined in (1.5), as stated in the following proposition (whose proof is given in Subsection 4.2).

Proposition 3.1. For every $E \in \mathcal{I}$, there is r > 0 such that $\mathbb{B}_r(E) \subset \mathcal{I}$, under the notation (1.5). Similarly, for every $E \in \mathcal{E}$, there is s > 0 such that $\mathbb{B}_s(E) \subset \mathcal{E}$.

Secondly, it is easy to see that

$$\exists \varepsilon > 0, \ (\lambda(E) \le \varepsilon) \Longrightarrow (E \in \mathcal{E}),$$
 (3.1)

and this implication is independent of the fragmentation indices $\delta_1(E)$ and $\delta_H(E)$. Indeed, calling

$$M' := \max_{[0,1]} |f'| \tag{3.2}$$

and $\varepsilon := e^{-M'} (4\pi)^{N/2} \theta/2 > 0$, with $\theta \in (0,1)$ as in (1.2), and assuming that $\lambda(E) \leq \varepsilon$, one has

$$0 < u(1,x) \le \frac{e^{M'}}{(4\pi)^{N/2}} \int_E e^{-|x-y|^2/4} dy \le \frac{e^{M'} \lambda(E)}{(4\pi)^{N/2}} \le \frac{\theta}{2}$$

for all $x \in \mathbb{R}^N$ from the maximum principle, hence $||u(t,\cdot)||_{L^{\infty}(\mathbb{R}^N)} \to 0$ as $t \to +\infty$ since f(0) = 0 and f < 0 in $(0,\theta)$.

Thirdly, highly contracted sets of a given set lead to extinction, and highly dilated sets of a given set lead to invasion, as stated in the following proposition.

Proposition 3.2. For any bounded non-negligible Borel $E \subset \mathbb{R}^N$, there are some real numbers $0 < \underline{\mu}_E \leq \overline{\mu}_E$ such that

$$\begin{cases}
(0 < \mu < \underline{\mu}_E) \implies (\mu E \in \mathcal{E}), \\
(\mu > \overline{\mu}_E) \implies (\mu E \in \mathcal{I}).
\end{cases}$$
(3.3)

Furthermore, the sets $\{\mu > 0 : \mu E \in \mathcal{E}\}$ and $\{\mu > 0 : \mu E \in \mathcal{I}\}$ are open, and

$$\{\mu > 0 : \mu E \in \mathcal{T}\} \neq \emptyset. \tag{3.4}$$

As a matter of fact, (3.1) implies that $\underline{\mu}_E := (\varepsilon/\lambda(E))^{1/N} > 0$ satisfies the first assertion of (3.3). If E has a non-empty interior \mathring{E} , that is, if there are $x_0 \in \mathbb{R}^N$ and r > 0 such that $E \supset B_r(x_0)$, then $\mu E \supset B_{\mu r}(\mu x_0)$ and $\mu E \in \mathcal{I}$ as soon as $\mu r > R_1$, with $R_1 > 0$ as in (1.6), from the comparison principle and the invariance of (1.1) with respect to translations. The end of the proof of the second assertion of (3.3), with the remaining case $\mathring{E} = \emptyset$, as well as the last part of Proposition 3.2, are done in Subsection 4.2. Having in hand (1.6) and Proposition 3.2, it is tempting to conjecture that, for any bounded non-negligible Borel $E \subset \mathbb{R}^N$, the set $\{\mu > 0 : \mu E \in \mathcal{T}\}$ would be a singleton, that is, $\underline{\mu}_E = \overline{\mu}_E$ if $\underline{\mu}_E$ and $\overline{\mu}_E$ respectively denote the largest and the smallest real numbers satisfying (3.3). This property is true if E is further assumed to be star-shaped with respect to a point, say x_0 , as follows from [33], since then the initial condition $\mathbbm{1}_{\mu(E-x_0)}$ are pointwise nondecreasing with respect to $\mu > 0$ and increasing in $L^1(\mathbb{R}^N)$, and since for each $\mu > 0$ the sets $\mu(E-x_0)$ and μE belong to the same set \mathcal{E} , \mathcal{T} or \mathcal{I} . The proof of the uniqueness of the element of $\{\mu > 0 : \mu E \in \mathcal{T}\}$ is however still open in the general case of sets E that are not star-shaped.

All remaining results are concerned with the role of the fragmentation indices $\delta_1(E)$ and $\delta_H(E)$ on the membership of E in \mathcal{I} or \mathcal{E} . The first such result asserts that equimeasurable and highly fragmented sets for the index δ_1 belong to the extinction set \mathcal{E} .

Proposition 3.3. For each given m > 0, there is $\varepsilon_m > 0$ such that any bounded Borel set E satisfying $\lambda(E) = m$ and $\delta_1(E) \geq 1 - \varepsilon_m$ belongs to \mathcal{E} .

Two comments are in order on this result, which is proved in Subsection 4.2. First of all, the conclusion does not hold without the hypothesis that the sets have a given Lebesgue measure m, because of Proposition 3.2 and the invariance of δ_1 with respect to dilations. More explicitly, for instance, the sets Rn^2E_n , with $n \geq 2$, E_n as in (2.4) and $R > R_1$, all belong to \mathcal{I} from (1.6) and the comparison principle, whereas $\delta_1(Rn^2E_n) = \delta_1(E_n) \to 1$ as $n \to +\infty$.

We also point out that a similar statement as Proposition 3.3 would be false if δ_1 were replaced by δ_H . For instance, consider $R > R_1$ and $E'_n := B_R \cup B_1(x_n)$, with $R+1 < |x_n| \to +\infty$ as $n \to +\infty$. Then, for every $n \in \mathbb{N}$, one has $\lambda(E'_n) = \omega_N(R^N + 1)$ and $E'_n \in \mathcal{I}$ from the comparison principle, whereas $\delta_H(E'_n) \to 1$ as $n \to +\infty$. On the other hand, the extinction is nevertheless possible for highly fragmented sets for the index δ_H with fixed Lebesgue measure, even if they are not highly fragmented for the index δ_1 . For instance, for any $r \in (0, 2R_1)$, the sets $D_a := (-a - r, -a) \cup (a, a + r)$, given in the first paragraph of Section 2, with measure 2r, belong to \mathcal{E} for all a > 0 large enough, while $\delta_H(D_a) \to 1$ and $\delta_1(D_a) \to 1/2$ as $a \to +\infty$ (this follows from [12], and a similar property holds immediately in any dimension $N \geq 1$).

After dealing with highly fragmented sets, let us now consider weakly fragmented sets, in the following proposition and an immediate corollary.

Proposition 3.4. For each given $m > \lambda(B_{R_1}) = \omega_N R_1^N$, with R_1 given in (1.6), there is $\eta_m > 0$ such that any bounded Borel set E satisfying $\lambda(E) = m$ and $\delta_1(E) \leq \eta_m$ belongs to \mathcal{I} .

Proposition 3.4 is proved in Subsection 4.2. Together with Proposition 2.1 (i), the following corollary immediately holds.

Corollary 3.5. For each given $m > \lambda(B_{R_1}) = \omega_N R_1^N$, any bounded Borel set E satisfying $\lambda(E) = m$ and $\delta_H(E) \leq \eta_m/\gamma$ belongs to \mathcal{I} , where $\eta_m > 0$ is given in Proposition 3.4 and $\gamma > 0$ in (2.10).

As for Proposition 3.3, the conclusions of Proposition 3.4 and Corollary 3.5 immediately do not hold without the hypothesis that the sets have a given Lebesgue measure m. For instance, the balls B_r with $0 < r < R_1$ belong to \mathcal{E} by (1.6), but nevertheless have fragmentation indices $\delta_1(B_r)$ and $\delta_H(B_r)$ equal to 0.

Finally, after the previous results about highly or weakly fragmented sets with given Lebesgue measure, we investigate the following question: if two bounded Borel sets E_1 and E_2 have the same Lebesgue measure, is it possible to decide about their membership of \mathcal{E} or \mathcal{I} according to the comparison of the values of $\delta_1(E_1)$ and $\delta_1(E_2)$, or $\delta_H(E_1)$ and $\delta_H(E_2)$? In other words, is there a kind of monotonicity of the large-time dynamics of the solutions of (1.1) and (1.7) with respect to the fragmentation indices δ_1 or δ_H of the initial set E?

Actually, the answer to this question is easily seen to be false for the index δ_H . Let us explain why in this paragraph. Corollary 3.5 says that, for a given $m > \lambda(B_{R_1})$, the weakly fragmented bounded Borel sets E for the index δ_H (namely, $\delta_H(E) \leq \eta_m/\gamma$) with Lebesgue measure equal to m belong to the invasion set \mathcal{I} . Now, call $R := (m/\omega_N)^{1/N} > R_1$, pick any $e \in \mathbb{R}^N \setminus \{0\}$ and define the bounded Borel sets

$$O_n := \bigcup_{k=1}^n B_{R/n^{1/N}}(ke).$$

One has $\lambda(O_n) = m$ for all $n \in \mathbb{N}$ large enough, and $\delta_1(O_n) \to 1$ as $n \to +\infty$ (and also $\delta_H(O_n) \to 1$ as $n \to +\infty$ by Proposition 2.1). Using Proposition 3.3, there is $n_0 \ge 2$ large enough such that $\lambda(O_{n_0}) = m$ and $O_{n_0} \in \mathcal{E}$. Pick now any $R' \in (R_1, R)$ and call r' > 0 such that $\omega_N r'^N = m - \omega_N R'^N$. Choose any sequence $(x_p)_{p \in \mathbb{N}}$ in \mathbb{R}^N such that $|x_p| \to +\infty$ as $p \to +\infty$. The bounded Borel sets

$$Q_p := B_{R'} \cup B_{r'}(x_p)$$

satisfy $\lambda(Q_p) = m$ for all p large enough, and they belong to \mathcal{I} since they contain the ball $B_{R'}$ with $R' > R_1$. Furthermore, $\lim_{p \to +\infty} \delta_H(Q_p) = 1 > \delta_H(O_{n_0}) > 0$ (the inequality $\delta_H(O_{n_0}) > 0$ holds since $n_0 \ge 2$ and thus O_{n_0} is not a ball up to a negligible set). Therefore, there is $p_0 \in \mathbb{N}$ large enough such that $\lambda(Q_{p_0}) = m$ and $\delta_H(Q_{p_0}) > \delta_H(O_{n_0}) > 0$, while $Q_{p_0} \in \mathcal{I}$. As a conclusion, the sets B_R , O_{n_0} and O_{p_0} are equimeasurable,

$$0 = \delta_H(B_R) < \delta_H(O_{n_0}) < \delta_H(Q_{p_0}),$$

while

$$B_R \in \mathcal{I}, \ O_{n_0} \in \mathcal{E} \ \text{and} \ Q_{p_0} \in \mathcal{I}.$$

In other words, there is no monotonicity of the large-time dynamics of the solutions of (1.1) and (1.7) with respect to $\delta_H(E)$ in the class of equimeasurable sets.

As far as the fragmentation index δ_1 is concerned, the answer to the same monotonicity question is not that clear. For a given $m > \lambda(B_{R_1})$, we know from Proposition 3.4 that the weakly fragmented bounded Borel sets E for the index δ_1 (namely, $\delta_1(E) \leq \eta_m$) with Lebesgue measure equal to m belong to the invasion set \mathcal{I} , while the highly fragmented ones (namely, $\delta_1(E) \geq 1 - \varepsilon_m$) belong to the extinction set \mathcal{E} , from Proposition 3.3. However, what happens for intermediate values of $\delta_1(E)$ is not as clear as with the fragmentation index $\delta_H(E)$. The last main result of the paper actually shows that there is in general no monotonicity of the large-time dynamics of the solutions of (1.1) and (1.7) with respect to $\delta_1(E)$ in the class of equimeasurable sets.

Theorem 3.6. There are some bounded non-negligible Borel sets E_1 and E_2 such that $\lambda(E_1) = \lambda(E_2)$, $0 < \delta_1(E_2) < \delta_1(E_1)$, while $E_1 \in \mathcal{I}$ and $E_2 \in \mathcal{E}$.

With the notations of Theorem 3.6 and Proposition 3.3, calling $m := \lambda(E_1) = \lambda(E_2) > 0$, one has $0 < \delta_1(E_2) < \delta_1(E_1) < 1 - \varepsilon_m$. Consider now the sets E_n defined in (2.4) (for $n \geq 3$, to avoid the confusion with the above sets E_1 and E_2) and satisfying (2.5). There is $n_0 \geq 3$ such that $\delta_1(E_{n_0}) \geq 1 - \varepsilon_m$. Call $F := \mu E_{n_0}$, with $\mu = (m/\lambda(E_{n_0}))^{1/N} > 0$, hence $\lambda(F) = m$, $\delta_1(F) = \delta_1(E_{n_0}) \geq 1 - \varepsilon_m$ and $F \in \mathcal{E}$ by Proposition 3.3. The bounded non-negligible Borel sets E_0 , E_1 and F are equimeasurable and satisfy

$$\delta_1(E_2) < \delta_1(E_1) < \delta_1(F),$$

while

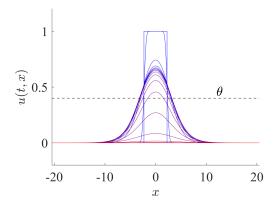
$$E_2 \in \mathcal{E}, E_1 \in \mathcal{I}, \text{ and } F \in \mathcal{E}.$$

Theorem 3.6, sustained by numerical simulations (see Figure 1), is proved in Subsection 4.3. We provide a proof based on some homogenization results, holding in any dimension $N \geq 1$. We also give another proof, holding in dimensions $N \geq 2$, based on completely different geometric arguments and the construction of suitable initial sets as intersections of balls with cubes. The first proof, based on homogenization techniques, appears to highlight an intrinsically positive effect of fragmentation on invasibility, as shown in Fig. 1. The second proof, which is more geometric, is instead based on the lack of the δ_1 index to capture certain types of fragmentation, rather than on an intrinsic effect of fragmentation. Therefore, we believe that the arguments of the first proof, hence the non-monotonicity of the large-time dynamics of the solutions of (1.1) and (1.7) with respect to the fragmentation of E in the class of equimeasurable sets, should remain valid for a large class of fragmentation indices (unlike the arguments of the second proof, which should no longer hold for indices that capture the distance between the connected components of E more finely, see Fig. 2 in Subsection 4.3). Other indices could be based, for example, on the Wasserstein distance between the measures $\mathbb{1}_E$ and $\mathbb{1}_{B_{R_E}(x)}$.

4 Proofs of the main results

4.1 Comparisons between the indices δ_1 and δ_H : proof of Proposition 2.1

Proof of Proposition 2.1. As already underlined in Section 2 after the statement of Proposition 2.1, only parts (i) and (iii) remain to be proved. For the proof of (i), one can restrict



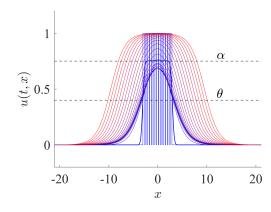


Figure 1: Evolution of the numerical solution u(t,x) of (1.1), with N=1, $f(s)=s(1-s)(s-\theta)$, and $\theta=0.4$, starting with an initial condition $u_0(x)=\mathbbm{1}_E(x)$ with (left) $E=E_1:=(-L/2,L/2)$ and (right) $E=E_2:=\bigcup_{x\in\mathbb{Z}\cap[-k,k]}\left[\frac{x}{z}+\left(-\frac{\alpha}{2z},\frac{\alpha}{2z}\right)\right]$. The curves correspond to the solution for t=0, t=0.05 and for t ranging from 4 to 80 with a step size of 4. The gradient color goes from blue for the earliest times to red for the latest times. In the left panel, $\lambda(E_1)=L=4.55$ and $\delta_1(E_1)=0$; in the right panel $\alpha=3/4$, z=2.16, k=6, $\lambda(E_2)=4.52$ and $\delta_1(E_2)=0.23$. We observe here that, though z is not that large and $\lambda(E_2)$ is even smaller than $\lambda(E_1)$, fragmentation (right panel) improves invasion success. We also note that $u(t,x)\approx \alpha\mathbbm{1}_{B_R}$ with $R=k/z+\alpha/(2z)$ when $t\ll 1$. The Matlab code used for the computations is available at http://doi.org/10.17605/0SF.IO/ZM479.

without loss of generality to the class of bounded Borel sets E such that $\rho_E = 1$, since the indices δ_1 and δ_H are invariant by dilation or contraction of the sets. Consider any such set E. Assume first that

$$\delta_H(E) < \frac{1}{4}.$$

From (2.7) and $0 < R_E \le \rho_E = 1$, there is then a point $y_E \in \mathbb{R}^N$ such that

$$d_H(E, B_{R_E}(y_E)) = \delta_H(E) (\rho_E + R_E) \le 2\delta_H(E) < \frac{1}{2},$$

hence $E \subset B_{R_E+1/2}(y_E)$ up to a negligible set. It follows that $R_E+1/2 \ge \rho_E=1$, that is, $R_E \ge 1/2$. On the other hand,

$$\delta_{1}(E) \leq \frac{\lambda(E \setminus B_{R_{E}}(y_{E}))}{\lambda(E)} \leq \frac{\omega_{N}(R_{E} + d_{H}(E, B_{R_{E}}(y_{E})))^{N} - \omega_{N}R_{E}^{N}}{\omega_{N}R_{E}^{N}}$$
$$\leq \frac{(R_{E} + 2\delta_{H}(E))^{N} - R_{E}^{N}}{R_{E}^{N}}.$$

Since $1/2 \le R_E \le R_E + 2\delta_H(E) \le 3/2$, one infers from the mean value theorem that $\delta_1(E) \le 2^N N(3/2)^{N-1} \times (2\delta_H(E)) = 4N3^{N-1}\delta_H(E)$.

If now $\delta_H(E) \ge 1/4$, then $\delta_1(E) < 1 \le 4\delta_H(E)$. Finally, $\delta_1(E) \le 4N3^{N-1}\delta_H(E)$ for every bounded non-negligible Borel set E, that is, (2.10) holds with $\gamma := 4N3^{N-1} > 0$.

For the proof of (iii), one can restrict without loss of generality to the class of bounded Borel sets E such that $R_E = \sqrt{N}/2$. Consider any such set E. Assume first that

$$0 < 1 - \delta_1(E) < \frac{1}{6^N}.$$

Then

$$\frac{1}{2} \times \left(\frac{1}{(1 - \delta_1(E))^{1/N}} - 1 \right) - 1 > \frac{1}{4} \times \frac{1}{(1 - \delta_1(E))^{1/N}} > \frac{3}{2},$$

and there is an integer $m_E \geq 2$ such that

$$\frac{1}{4} \times \frac{1}{(1 - \delta_1(E))^{1/N}} < m_E < \frac{1}{2} \times \left(\frac{1}{(1 - \delta_1(E))^{1/N}} - 1\right). \tag{4.1}$$

Now, since any cube with sides of unit length is included into a ball of radius $R_E = \sqrt{N}/2$, one has in particular

$$\max_{k \in \mathbb{Z}^N} \frac{\lambda(E \cap (k + (0, 1)^N))}{\lambda(E)} \le \max_{x \in \mathbb{R}^N} \frac{\lambda(E \cap B_{R_E}(x))}{\lambda(E)} = 1 - \delta_1(E),$$

where the last equality follows from (2.3). On the other hand, since $\lambda(E) > 0$ and the cubes $k + (0,1)^N$ cover \mathbb{R}^N up to a negligible set as k describes \mathbb{Z}^N , there is $k_E \in \mathbb{Z}^N$ such that

$$0 < \max_{k \in \mathbb{Z}^N} \lambda(E \cap (k + (0, 1)^N)) = \lambda(E \cap (k_E + (0, 1)^N)) \le (1 - \delta_1(E)) \lambda(E).$$

Thus,

$$\lambda \Big(E \cap \bigcup_{k \in \mathbb{Z}^N, ||k||_{\infty} \le m_E} (k_E + k + (0, 1)^N) \Big) \le (2m_E + 1)^N \times (1 - \delta_1(E)) \lambda(E).$$

Since $\lambda(E) > 0$ and $(2m_E + 1)^N (1 - \delta_1(E)) < 1$ from the right inequality in (4.1), it follows that

$$\lambda \Big(E \cap \Big(\mathbb{R}^N \setminus \bigcup_{k \in \mathbb{Z}^N, ||k||_{\infty} \le m_E} (k_E + k + (0, 1)^N) \Big) \Big) > 0,$$

hence there is $k'_E \in \mathbb{Z}^N$ with $||k'_E||_{\infty} > m_E$ (that is, $||k'_E||_{\infty} \geq m_E + 1$) such that $\lambda(E \cap (k_E + k'_E + (0,1)^N)) > 0$. Remembering that $\lambda(E \cap (k_E + (0,1)^N)) > 0$, one gets that diam $(E) \geq m_E$, hence

$$\rho_E \ge \frac{\text{diam}(E)}{2} \ge \frac{m_E}{2} \ge \frac{1}{8(1 - \delta_1(E))^{1/N}},$$

by using (2.2) and the left inequality in (4.1). Together with (2.9) and the normalization $R_E = \sqrt{N}/2$, one infers that

$$1 - \delta_H(E) \le \frac{2R_E}{\rho_E + R_E} \le \frac{2R_E}{\rho_E} \le 8\sqrt{N} (1 - \delta_1(E))^{1/N}.$$

If now $1 - \delta_1(E) \ge 1/6^N$, then $1 - \delta_H(E) \le 1 \le 6(1 - \delta_1(E))^{1/N}$. Finally, $1 - \delta_H(E) \le 8\sqrt{N}(1 - \delta_1(E))^{1/N}$ for every bounded non-negligible Borel set E, that is, (2.11) holds with $\eta := 8\sqrt{N}$. The proof of Proposition 2.1 is thereby complete.

4.2 Extinction vs. invasion for dilated and highly or weakly fragmented sets

This subsection is devoted to the proofs of Propositions 3.1-3.4 on the large-time dynamics of solutions of (1.1) for close, dilated, highly fragmented, or weakly fragmented indicator sets E in (1.7).

Proof of Proposition 3.1. Let us first consider E in \mathcal{I} and let us show that $F \in \mathcal{I}$ for any bounded Borel set F such that $d_1(E,F)$ is small enough. Let u_E and u_F denote the solutions of (1.1) with initial conditions $\mathbb{1}_E$ and $\mathbb{1}_F$ respectively. Pick any $\alpha \in (\theta,1)$ with $\theta \in (0,1)$ as in (1.2), remember the definition of $R_{\alpha} > 0$ in (1.6), and pick any $R \in (R_{\alpha}, +\infty)$. As $u_E(t, \cdot) \to 1$ as $t \to +\infty$ locally uniformly in \mathbb{R}^N , there is T > 0 such that $\min_{\overline{B_R}} u_E(T, \cdot) > \alpha$. Since

$$||u_E(T,\cdot) - u_F(T,\cdot)||_{L^{\infty}(\mathbb{R}^N)} \le \frac{e^{M'T}}{(4\pi T)^{N/2}} \times ||\mathbb{1}_E - \mathbb{1}_F||_{L_1(\mathbb{R}^N)}$$
(4.2)

from the maximum principle, with $M' = \max_{[0,1]} |f'|$ as in (3.2), there is then r > 0 such that $\min_{\overline{B_R}} u_F(T, \cdot) \ge \alpha$ if $d_1(E, F) = \|\mathbb{1}_E - \mathbb{1}_F\|_{L_1(\mathbb{R}^N)} < r$. Therefore, for every $F \in \mathbb{B}_r(E)$, there holds $1 \ge u_F(T, \cdot) \ge \alpha \mathbb{1}_{B_R}$ in \mathbb{R}^N , hence $u_F(t, \cdot) \to 1$ as $t \to +\infty$ locally uniformly in \mathbb{R}^N , from the comparison principle and (1.6) again. In other words, $\mathbb{B}_r(E) \subset \mathcal{I}$.

Let us now assume that $E \in \mathcal{E}$. Therefore, there is T > 0 such that $0 \le u_E(T, \cdot) \le \theta/3$ in \mathbb{R}^N . From (4.2), there is then s > 0 such that $u_F(T, \cdot) \le \theta/2$ in \mathbb{R}^N for every bounded Borel set F satisfying $d_1(E, F) < s$, and then the nonnegative function u_F converges to 0 as $t \to +\infty$ uniformly in \mathbb{R}^N , from (1.2) and the comparison principle. In other words, $\mathbb{B}_s(E) \subset \mathcal{E}$.

Proof of Proposition 3.2. As already underlined in Section 3 after the statement of Proposition 3.2, for the proof of (3.3), only the case of dilated sets μE with large μ and $\dot{E} = \emptyset$ remains to be dealt with. Let E be such a set. Since $\lambda(E) > 0$ and since

$$\frac{1}{\lambda(B_r(x))} \int_{B_r(x)} \mathbb{1}_E(y) \, dy \to \mathbb{1}_E(x) \text{ as } r \stackrel{>}{\to} 0 \text{ for almost every } x \in \mathbb{R}^N$$

by Lebesgue's differentiation theorem, there is $x_0 \in \mathbb{R}^N$ such that $\lambda(E \cap B_r(x_0)) \sim \lambda(B_r(x_0))$ as $r \stackrel{>}{\to} 0$. Since (1.1) is invariant by translation, one can assume without loss of generality that $x_0 = 0$. Pick any $R > R_1$, with R_1 defined in (1.6) with $\alpha = 1$ (hence, $B_R \in \mathcal{I}$). Since

$$\lambda(B_R) \ge \lambda(\mu E \cap B_R) = \mu^N \lambda(E \cap B_{R/\mu}) \underset{\mu \to +\infty}{\sim} \mu^N \lambda(B_{R/\mu}) = \lambda(B_R),$$

it follows that $d_1(B_R, \mu E \cap B_R) \to 0$ as $\mu + \infty$. Therefore, $\mu E \cap B_R \in \mathcal{I}$ for all $\mu > 0$ large enough, from Proposition 3.1. Finally, $\mu E \in \mathcal{I}$ for all $\mu > 0$ large enough, from the comparison principle.

Let us now turn to the proof of the openness of the sets $\{\mu > 0 : \mu E \in \mathcal{E}\}$ and $\{\mu > 0 : \mu E \in \mathcal{I}\}$. In the case when the bounded Borel set $E \subset \mathbb{R}^N$ has a negligible boundary for the N-dimensional Lebesgue measure λ , that is, if $\lambda(\partial E) = 0$, then, for

any $\mu_0 > 0$, $\lambda(\partial(\mu_0 E)) = 0$ and

$$d_{1}(\mu E, \mu_{0}E) = \int_{\mathbb{R}^{N}} |\mathbb{1}_{\mu E}(x) - \mathbb{1}_{\mu_{0}E}(x)| dx$$

$$= \int_{\widehat{\mu_{0}E}} |\mathbb{1}_{\mu E}(x) - \mathbb{1}_{\mu_{0}E}(x)| dx + \int_{\mathbb{R}^{N} \setminus \overline{\mu_{0}E}} |\mathbb{1}_{\mu E}(x) - \mathbb{1}_{\mu_{0}E}(x)| dx \xrightarrow{\mu \to \mu_{0}} 0$$

from Lebesgue's dominated convergence theorem. Therefore, in this case, the openness of the sets $\{\mu > 0 : \mu E \in \mathcal{E}\}\$ and $\{\mu > 0 : \mu E \in \mathcal{I}\}\$ follows from Proposition 3.1.

Consider now the case of a general bounded Borel set $E \subset \mathbb{R}^N$. Assume first that $\mu_0 > 0$ is such that

$$\mu_0 E \in \mathcal{I}$$
.

For $\mu > 0$, let u_{μ} denote here the solution of (1.1) with initial condition $\mathbb{1}_{\mu E}$. As in the proof of Proposition 3.1, pick any $\alpha \in (\theta, 1)$ and any $R > R' > R_{\alpha}$, with $R_{\alpha} > 0$ as in (1.6). By hypothesis, there holds $u_{\mu_0}(t, \cdot) \to 1$ as $t \to +\infty$ locally uniformly in \mathbb{R}^N . There is then T > 0 such that $\min_{\overline{B_R}} u_{\mu_0}(T, \cdot) > \alpha$. On the other hand, for every $\mu > 0$, the function

$$(t,x) \mapsto v_{\mu}(t,x) := u_{\mu} \left(\frac{\mu^2}{\mu_0^2} t, \frac{\mu}{\mu_0} x \right)$$

ranges in [0,1] and satisfies

$$\frac{\partial v_{\mu}}{\partial t} = \Delta v_{\mu} + \frac{\mu^2}{\mu_0^2} f(v_{\mu}), \quad t > 0, \ x \in \mathbb{R}^N$$

with initial condition $v_{\mu}(0,\cdot) = \mathbb{1}_{\mu_0 E} = u_{\mu_0}(0,\cdot)$. Therefore, the function $w_{\mu} := v_{\mu} - u_{\mu_0}$ vanishes at time t = 0 and satisfies

$$\frac{\partial w_{\mu}}{\partial t} \le \Delta w_{\mu} + f(v_{\mu}) - f(u_{\mu_0}) + M \left| \frac{\mu^2}{\mu_0^2} - 1 \right| \le \Delta w_{\mu} + M' |w_{\mu}| + M \left| \frac{\mu^2}{\mu_0^2} - 1 \right|$$

in $(0, +\infty) \times \mathbb{R}^N$, with $M := \max_{[0,1]} |f|$ and $M' = \max_{[0,1]} |f'|$. It then follows from the maximum principle that

$$w_{\mu}(t,x) \leq \frac{M}{M'} \left| \frac{\mu^2}{\mu_0^2} - 1 \right| \times \left(e^{M't} - 1 \right) \text{ for all } t \geq 0 \text{ and } x \in \mathbb{R}^N.$$

By arguing similarly with $u_{\mu_0} - v_{\mu}$, one gets the same bound from above for $|v_{\mu} - u_{\mu_0}|$. In particular, at time t = T, by rewriting v_{μ} in terms of u_{μ} and changing x into $(\mu_0/\mu)x$, one infers that

$$\left| u_{\mu} \left(\frac{\mu^2}{\mu_0^2} T, x \right) - u_{\mu_0} \left(T, \frac{\mu_0}{\mu} x \right) \right| \le \frac{M}{M'} \left| \frac{\mu^2}{\mu_0^2} - 1 \right| \times \left(e^{M'T} - 1 \right) \text{ for all } x \in \mathbb{R}^N.$$
 (4.3)

As a consequence, remembering that $\min_{\overline{B_R}} u_{\mu_0}(T,\cdot) > \alpha$ and that R > R', there is $\varepsilon \in (0, \mu_0)$ such that, if $|\mu - \mu_0| \le \varepsilon$, then

$$\min_{\overline{B_{R'}}} u_{\mu} \left(\frac{\mu^2}{\mu_0^2} T, \cdot \right) \ge \alpha,$$

hence $1 \ge u_{\mu}((\mu^2/\mu_0^2)T, \cdot) \ge \alpha \mathbb{1}_{B_{R'}}$ in \mathbb{R}^N . Since $R' > R_{\alpha}$, one concludes from (1.6) and the maximum principle that $u_{\mu}(t, \cdot) \to 1$ as $t \to +\infty$ locally uniformly in \mathbb{R}^N for

each $\mu \in (\mu_0 - \varepsilon, \mu_0 + \varepsilon)$, that is, $\mu E \in \mathcal{I}$ for any such μ . From the arbitrariness of μ_0 such that $\mu_0 E \in \mathcal{I}$, the openness of $\{\mu > 0 : \mu E \in \mathcal{I}\}$ has been shown.

Let now $\mu_0 > 0$ be such that

$$\mu_0 E \in \mathcal{E}$$
,

and let T > 0 be such that $0 \le u_{\mu_0}(T, \cdot) \le \theta/3$ in \mathbb{R}^N . The inequality (4.3), which holds independently of the hypothesis $\mu_0 E \in \mathcal{E}$, provides the existence of $\varepsilon \in (0, \mu_0)$ such that, if $|\mu - \mu_0| \le \varepsilon$, then $u_{\mu}((\mu^2/\mu_0^2)T, \cdot) \le \theta/2$ in \mathbb{R}^N , hence $u_{\mu}(t, \cdot) \to 0$ as $t \to +\infty$ uniformly in \mathbb{R}^N , from (1.2) and the comparison principle. Finally, $\mu E \in \mathcal{E}$ for all $\mu \in (\mu_0 - \varepsilon, \mu_0 + \varepsilon)$, and the set $\{\mu > 0 : \mu E \in \mathcal{E}\}$ is open.

Lastly, property (3.4) immediately follows from (3.3) and the openness of the sets $\{\mu > 0 : \mu E \in \mathcal{E}\}\$ and $\{\mu > 0 : \mu E \in \mathcal{I}\}\$, together with the fact that $\mu E \in \mathcal{E} \cup \mathcal{T} \cup \mathcal{I}$ for every $\mu > 0$. The proof of Proposition 3.2 is thereby complete.

Proof of Proposition 3.3. Let m > 0 be given, define

$$s := \frac{1}{\sqrt{N}} \times \left(\frac{m}{\omega_N}\right)^{1/N} > 0,$$

and observe that any cube of measure $(2s)^N$ is included into a ball of measure m. Consider now any bounded Borel set E with $\lambda(E) = m$, and let u be the solution of (1.1) with initial condition (1.7). Formula (2.3) and the previous observations then imply that $\lambda(Q \cap E) \leq m(1 - \delta_1(E))$ for every cube Q of measure $(2s)^N$. Therefore, remembering the definition $M' = \max_{[0,1]} |f'|$, it follows from the comparison principle that, for every $x \in \mathbb{R}^N$,

$$\begin{array}{lcl} 0 & \leq & u(1,x) & \leq & \frac{e^{M'}}{(4\pi)^{N/2}} \times \int_{E} e^{-|x-y|^{2}/4} dy \\ \\ & = & \frac{e^{M'}}{(4\pi)^{N/2}} \times \sum_{k \in \mathbb{Z}^{N}} \int_{E \cap (x+2ks+(-s,s)^{N})} e^{-|x-y|^{2}/4} dy \\ \\ & \leq & \frac{e^{M'}}{(4\pi)^{N/2}} \times m \times (1-\delta_{1}(E)) \times \sum_{k \in \mathbb{Z}^{N}} e^{-\max(2\|k\|_{\infty}-1;0)^{2}s^{2}/4}, \end{array}$$

where the above two series converge. Since the right-hand side of the last inequality does not depend on x, there is then $\varepsilon_m > 0$ such that $0 \le u(1, \cdot) \le \theta/2$ in \mathbb{R}^N as soon as $\delta_1(E) \ge 1 - \varepsilon_m$, hence $u(t, \cdot) \to 0$ as $t \to +\infty$ uniformly in \mathbb{R}^N . In other words, every bounded Borel set E such that $\lambda(E) = m$ and $\delta_1(E) \ge 1 - \varepsilon_m$ belongs to \mathcal{E} .

Proof of Proposition 3.4. Fix any $m > \lambda(B_{R_1}) = \omega_N R_1^N$. Define $\alpha = (\theta+1)/2 \in (\theta,1)$ and $R = R_{\alpha} + 1$, with $\theta \in (0,1)$ and R_{α} given in (1.2) and (1.6). Let v denote the solution of (1.1) with initial condition $\mathbbm{1}_{B_{(m/\omega_N)^{1/N}}}$. Since $(m/\omega_N)^{1/N} > R_1$ by assumption, it follows from (1.6) that $v(t,\cdot) \to 1$ as $t \to +\infty$ locally uniformly in \mathbb{R}^N , hence there is T > 0 such that $\min_{\overline{B_R}} v(T,\cdot) > \alpha$. Consider now any bounded Borel set E such that $\lambda(E) = m$ and, from (2.3), let $z_E \in \mathbb{R}^N$ be such that

$$\|\mathbb{1}_E - \mathbb{1}_{B_{(m/\omega_N)^{1/N}}(z_E)}\|_{L_1(\mathbb{R}^N)} = 2\lambda(E \setminus B_{(m/\omega_N)^{1/N}}(z_E)) = 2\lambda(E)\delta_1(E) = 2m\delta_1(E).$$

As in (4.2), the solution u_E of (1.1) with initial condition $\mathbb{1}_E$ satisfies

$$||u_E(T,\cdot) - v(T,\cdot - z_E)||_{L^{\infty}(\mathbb{R}^N)} \le \frac{e^{M'T}}{(4\pi T)^{N/2}} \times ||\mathbb{1}_E - \mathbb{1}_{B_{(m/\omega_N)^{1/N}}(z_E)}||_{L_1(\mathbb{R}^N)},$$

hence $||u_E(T, \cdot + z_E) - v(T, \cdot)||_{L^{\infty}(\mathbb{R}^N)} \leq 2e^{M'T}(4\pi T)^{-N/2}m \times \delta_1(E)$. Observing that α and T only depend on f and m (and on the dimension N), and remembering that $\min_{\overline{B_R}} v(T, \cdot) > \alpha$, there is then $\eta_m > 0$ such that $\min_{\overline{B_R}} u_E(T, \cdot + z_E) \geq \alpha$ as soon as $\delta_1(E) \leq \eta_m$. For any such E, one then has $u_E(T, \cdot + z_E) \geq \alpha \mathbb{1}_{B_R}$, hence $u_E(t, \cdot) \to 1$ as $t \to +\infty$ locally uniformly in \mathbb{R}^N , by (1.6) and the definition $R = R_{\alpha} + 1$. Therefore, any bounded Borel set E such that $\lambda(E) = m$ and $\delta_1(E) \leq \eta_m$ belongs to \mathcal{I} .

4.3 Proof of Theorem 3.6

We provide two different proofs of Theorem 3.6. The first one holds in any dimension $N \geq 1$ and is based on homogenization techniques, and the constructed sets E_1 and E_2 have similarities with the sets F_n defined in (2.6). The alternate proof holds in dimensions $N \geq 2$, and the constructed sets involve cubes and their intersections with balls.

Proof of Theorem 3.6 in any dimension $N \ge 1$. Consider any $\alpha \in (\theta, 1)$, with $\theta \in (0, 1)$ as in (1.2). Fix then some real numbers R and R', close enough to R_{α} , such that

$$0 < R' < R_{\alpha} < R, \quad \alpha^{1/N} R < R', \quad \text{and} \quad 0 < \alpha R^N - \alpha R'^N < \min\left(\frac{R_1^N}{2^N}, \alpha^2 R^N\right), \quad (4.4)$$

with $R_{\alpha} > 0$ and $R_1 > 0$ as in (1.6). Owing to the definition of R_{α} , the solution \widetilde{v} of (1.1) with initial condition $\alpha \mathbbm{1}_{B_{R'}}$ converges to 0 as $t \to +\infty$ uniformly in \mathbb{R}^N , thus there is T > 0 such that $\widetilde{v}(T,\cdot) \leq \theta/3$ in \mathbb{R}^N . For $\beta \in (\alpha,1]$, let \widetilde{v}^{β} be the solution of (1.1) with initial condition $\beta \mathbbm{1}_{B_{R'}}$. Since $0 \leq \widetilde{v}^{\beta}(T,\cdot) - \widetilde{v}(T,\cdot) \leq e^{M'T}(4\pi T)^{-N/2} \|\beta \mathbbm{1}_{B_{R'}} - \alpha \mathbbm{1}_{B_{R'}}\|_{L^1(\mathbb{R}^N)}$ from the maximum principle, with $M' = \max_{[0,1]} |f'|$ as in (3.2), one can choose $\beta \in (\alpha,1)$, close enough to α , so that $0 \leq \widetilde{v}^{\beta}(T,\cdot) \leq \theta/2$ in \mathbb{R}^N , hence

$$\widetilde{v}^{\beta}(t,\cdot) \to 0 \text{ as } t \to +\infty \text{ uniformly in } \mathbb{R}^N$$
 (4.5)

by (1.2) and the maximum principle. Even if it means decreasing β in $(\alpha, 1)$, one can also assume without loss of generality that

$$0 < \alpha R^N - \beta R'^N < \min\left(\frac{R_1^N}{2^N}, \alpha^2 R^N\right). \tag{4.6}$$

Consider the following bounded Borel sets

$$F_n := \bigcup_{z \in (\mathbb{Z}^{N/n}) \cap B_R} \left[z + \left(-\frac{\alpha^{1/N}}{2n}, \frac{\alpha^{1/N}}{2n} \right)^N \right]$$

and

$$G_n := \bigcup_{z \in (\mathbb{Z}^N/n) \cap B_{R'}} \left[z + \left(-\frac{\beta^{1/N}}{2n}, \frac{\beta^{1/N}}{2n} \right)^N \right]$$

for $n \geq 1$. The sets F_n and G_n are such that (see (2.6) and below)

$$\lim_{n \to +\infty} \lambda(F_n) = \alpha \,\omega_N R^N > \beta \,\omega_N {R'}^N = \lim_{n \to +\infty} \lambda(G_n) \tag{4.7}$$

and

$$\lim_{n \to +\infty} \delta_1(F_n) = 1 - \alpha > 1 - \beta = \lim_{n \to +\infty} \delta_1(G_n) > 0.$$
(4.8)

Together with (4.6), there is then $n_1 \in \mathbb{N}$ such that

$$0 < \lambda(G_n) < \lambda(F_n) \text{ and } \lambda(F_n) - \lambda(G_n) < \min\left(\omega_N \frac{R_1^N}{2^N}, \alpha^2 \omega_N R^N\right)$$
 (4.9)

for all $n > n_1$.

Let then u_n, v_n, u and $v = \tilde{v}^{\beta}$ be the solutions of (1.1) with respective initial conditions $\mathbb{1}_{F_n}$, $\mathbb{1}_{G_n}$, $\alpha \, \mathbb{1}_{B_R}$ and $\beta \, \mathbb{1}_{B_{R'}}$. Let similarly $U_n, \, V_n, \, U$ and V be the solutions of the heat equation $\frac{\partial z}{\partial t} = \Delta z$ with respective initial conditions $\mathbb{1}_{F_n}, \, \mathbb{1}_{G_n}, \, \alpha \, \mathbb{1}_{B_R}$ and $\beta \, \mathbb{1}_{B_{R'}}$. We point out that these 8 functions are nonnegative, and that $u_n, \, v_n, \, u$ and v range in [0,1]. Since $\|\alpha \mathbb{1}_{B_R} - (\alpha - \varepsilon) \mathbb{1}_{B_{R-\varepsilon}}\|_{L^1(\mathbb{R}^N)} \to 0$ as $\varepsilon \to 0$, and since $u(t,\cdot) \to 1$ as $t \to +\infty$ locally uniformly in \mathbb{R}^N by (1.6) (since $R > R_{\alpha}$), it follows as in the proof of Proposition 3.1 that there is $\varepsilon_1 \in (0,R)$ such that the solution u^{ε_1} of (1.1) with initial condition

$$u_0^{\varepsilon_1} := (\alpha - \varepsilon_1) \mathbb{1}_{B_{R-\varepsilon_1}}$$

converges to 1 as $t \to +\infty$ locally uniformly in \mathbb{R}^N . Notice also that $U(t,\cdot) \to \alpha$ as $t \to 0$ locally uniformly in B_R . As a consequence, there is $t_1 > 0$ small enough such that

$$U(t_1, \cdot) \ge \left(\alpha - \frac{\varepsilon_1}{3}\right) \mathbb{1}_{B_{R-\varepsilon_1}} \text{ in } \mathbb{R}^N.$$

Even if it means decreasing $t_1 > 0$, one can assume without loss of generality that $(\alpha - \varepsilon_1/2)e^{-M't_1} \ge \alpha - \varepsilon_1$. Now, owing to the definition of F_n , it follows by homogenization that

$$0 \leq U_n(t_1, x) = \frac{1}{(4\pi t_1)^{N/2}} \sum_{z \in (\mathbb{Z}^N/n) \cap B_R} \int_{z + (-\alpha^{1/N}/(2n), \alpha^{1/N}/(2n))^N} e^{-|x-y|^2/(4t_1)} dy$$

$$\xrightarrow[n \to +\infty]{} \frac{1}{(4\pi t_1)^{N/2}} \int_{B_R} \alpha e^{-|x-y|^2/(4t_1)} dy = U(t_1, x)$$

uniformly with respect to $x \in \mathbb{R}^N$. Thus, there is $n_2 \geq n_1$ such that

$$U_n(t_1, \cdot) \ge \left(\alpha - \frac{\varepsilon_1}{2}\right) \mathbb{1}_{B_{R-\varepsilon_1}}$$
 in \mathbb{R}^N for all $n \ge n_2$.

Since $1 \geq u_n(t_1,\cdot) \geq e^{-M't_1}U_n(t_1,\cdot)$ in \mathbb{R}^N from the maximum principle, one infers that $1 \geq u_n(t_1,\cdot) \geq (\alpha - \varepsilon_1)\mathbb{1}_{B_{R-\varepsilon_1}} = u_0^{\varepsilon_1}$ in \mathbb{R}^N for all $n \geq n_2$, hence $u_n(t,\cdot) \to 1$ as $t \to +\infty$ locally uniformly in \mathbb{R}^N . In other words,

$$F_n \in \mathcal{I} \text{ for all } n \ge n_2.$$
 (4.10)

Consider now the solution $v = \tilde{v}^{\beta}$ of (1.1) with initial condition $\beta \mathbbm{1}_{B_{R'}}$. By (4.5), there is T' > 0 such that $0 \le v(T', \cdot) \le \theta/3$ in \mathbb{R}^N . For $\varepsilon \in (0, 1 - \beta)$, let v^{ε} be the solution of (1.1) with initial condition

$$v_0^\varepsilon := (\beta + \varepsilon) \, \mathbb{1}_{B_{R' + \varepsilon}} + \varepsilon \, \mathbb{1}_{\mathbb{R}^N \backslash B_{R' + \varepsilon}}.$$

Since

$$\begin{split} 0 \! \leq \! v^{\varepsilon}(T',x) - \! v(T',x) \! \leq \! \frac{e^{M'T'}}{(4\pi T')^{N/2}} \Big[\int_{B_{R'}} \! \varepsilon \, e^{-|x-y|^2/(4T')} dy + \! \int_{B_{R'+\varepsilon} \backslash B_{R'}} \! (\beta + \varepsilon) \, e^{-|x-y|^2/(4T')} dy \\ + \int_{\mathbb{R}^N \backslash B_{R'+\varepsilon}} \! \varepsilon \, e^{-|x-y|^2/(4T')} dy \Big] \end{split}$$

for all $x \in \mathbb{R}^N$, one infers that $v^{\varepsilon}(T',\cdot) \to v(T',\cdot)$ uniformly in \mathbb{R}^N as $\varepsilon \to 0$. Thus, there is $\varepsilon_2 \in (0,1-\beta)$ such that $0 \le v^{\varepsilon_2}(T',\cdot) \le \theta/2$ in \mathbb{R}^N , hence $v^{\varepsilon_2}(t,\cdot) \to 0$ as $t \to +\infty$ uniformly in \mathbb{R}^N . On the other hand, $\sup_{B_{R'+\varepsilon_2}} v(t,\cdot) \to \beta$ and $\sup_{\mathbb{R}^N \setminus B_{R'+\varepsilon_2}} v(t,\cdot) \to 0$ as $t \to 0$, while $V(t,\cdot) \le e^{M't}v(t,\cdot)$ in \mathbb{R}^N for all $t \ge 0$, from the maximum principle. As a consequence, there is $t_2 > 0$ (independent of n) small enough such that

$$V(t_2,\cdot) \le \left(\beta + \frac{\varepsilon_2}{3}\right) \mathbb{1}_{B_{R'+\varepsilon_2}} + \frac{\varepsilon_2}{3} \mathbb{1}_{\mathbb{R}^N \setminus B_{R'+\varepsilon_2}}.$$

Even if it means decreasing $t_2 > 0$, one can also assume without loss of generality that $(\beta + \varepsilon_2/2)e^{M't_2} \leq \beta + \varepsilon_2$ and $(\varepsilon_2/2)e^{M't_2} \leq \varepsilon_2$. Now, owing to the definition of G_n , there holds that $V_n(t_2, \cdot) \to V(t_2, \cdot)$ uniformly in \mathbb{R}^N as $n \to +\infty$ by homogenization, as for $U_n(t_1, \cdot)$ and $U(t_1, \cdot)$ above. Therefore, there is $n_3 \geq n_2$ such that

$$V_n(t_2,\cdot) \le \left(\beta + \frac{\varepsilon_2}{2}\right) \mathbb{1}_{B_{R'+\varepsilon_2}} + \frac{\varepsilon_2}{2} \mathbb{1}_{\mathbb{R}^N \setminus B_{R'+\varepsilon_2}} \text{ in } \mathbb{R}^N \text{ for all } n \ge n_3.$$

As $0 \leq v_n(t_2,\cdot) \leq e^{M't_2}V_n(t_2,\cdot)$ in \mathbb{R}^N from the maximum principle, one infers that $0 \leq v_n(t_2,\cdot) \leq (\beta + \varepsilon_2)\mathbbm{1}_{B_{R'+\varepsilon_2}} + \varepsilon_2\mathbbm{1}_{\mathbb{R}^N\setminus B_{R'+\varepsilon_2}} = v_0^{\varepsilon_2}$ in \mathbb{R}^N for all $n \geq n_3$, hence

$$0 \le v_n(t_2 + t, \cdot) \le v^{\varepsilon_2}(t, \cdot)$$
 in \mathbb{R}^N for all $t \ge 0$ and $n \ge n_3$ (4.11)

by the maximum principle. In particular, $v_n(t,\cdot) \to 0$ as $t \to +\infty$ uniformly in \mathbb{R}^N , that is, $G_n \in \mathcal{E}$, for all $n \geq n_3$.

Remember (4.9) and define

$$0 < \rho_n := \left(\frac{\lambda(F_n) - \lambda(G_n)}{\omega_N}\right)^{1/N} < \frac{R_1}{2} \tag{4.12}$$

for $n \ge n_1$. Consider finally a sequence of points $(x_n)_{n \ge n_1}$ in \mathbb{R}^N such that $B_{\rho_n}(x_n) \cap G_n = \emptyset$ for every $n \ge n_1$, and $|x_n| \to +\infty$ as $n \to +\infty$. The bounded Borel sets

$$H_n := G_n \cup B_{\rho_n}(x_n)$$

satisfy

$$\lambda(H_n) = \lambda(F_n) > 0 \text{ for all } n \ge n_1. \tag{4.13}$$

Since $\rho_n < R_1/2$, all the balls $B_{\rho_n}(x_n)$ belong to \mathcal{E} , by definition of R_1 .

For the Cauchy problem (1.1) with the initial condition $\mathbb{1}_{H_n}$, the two components $\mathbb{1}_{G_n}$ and $\mathbb{1}_{B_{\rho_n}(x_n)}$ act as almost independent initial conditions for n large enough and then H_n will belong to \mathcal{E} for every n large enough (as in the example of the sets D_a given in the first paragraph of Section 2). More precisely, to show this fact, fix first $\rho > 0$ such that $G_n \subset B_\rho$ and $\rho_n \leq \rho$ for all $n \geq n_1$. Denote w, w_n , and z_n , the solutions of (1.1) with initial conditions $\mathbb{1}_{B_{R_1/2}}$, $\mathbb{1}_{B_{\rho_n}(x_n)}$, and $\mathbb{1}_{H_n}$, respectively. From (4.11)-(4.12) and $\lim_{t\to+\infty} \|v^{\varepsilon_2}(t,\cdot)\|_{L^\infty(\mathbb{R}^N)} = 0$, together with the definition of R_1 , there is $\tau > 0$ such that

$$0 \le v_n(\tau, \cdot) \le \frac{\theta}{3}$$
 and $0 \le w_n(\tau, \cdot) \le w(\tau, \cdot - x_n) \le \frac{\theta}{3}$ in \mathbb{R}^N , for all $n \ge n_3$. (4.14)

Furthermore, since $H_n = G_n \cup B_{\rho_n}(x_n)$, the maximum principle yields

$$0 \le v_n(\tau, x) \le z_n(\tau, x) \le v_n(\tau, x) + \frac{e^{M'\tau}}{(4\pi\tau)^{N/2}} \int_{B_{\rho_n}(x_n)} e^{-|x-y|^2/(4\tau)} dy$$
$$\le v_n(\tau, x) + \frac{e^{M'\tau}}{(4\pi\tau)^{N/2}} \int_{B_{\rho}} e^{-|x-x_n-y|^2/(4\tau)} dy$$

and

$$0 \le w_n(\tau, x) \le z_n(\tau, x) \le w_n(\tau, x) + \frac{e^{M'\tau}}{(4\pi\tau)^{N/2}} \int_{G_n} e^{-|x-y|^2/(4\tau)} dy$$
$$\le v_n(\tau, x) + \frac{e^{M'\tau}}{(4\pi\tau)^{N/2}} \int_{B_n} e^{-|x-y|^2/(4\tau)} dy$$

for all $x \in \mathbb{R}^N$ and $n \ge n_1$. Let $\sigma > 0$ be such that

$$\frac{e^{M'\tau}}{(4\pi\tau)^{N/2}}\int_{B_{\varrho}}e^{-|\xi-y|^2/(4\tau)}dy \leq \frac{\theta}{6} \text{ for all } |\xi| \geq \sigma.$$

Together with (4.14), it follows that $0 \le z_n(\tau, x) \le \theta/2$ for all $n \ge n_3$ and for all $x \in \mathbb{R}^n$ such that either $|x - x_n| \ge \sigma$ or $|x| \ge \sigma$. Since $|x_n| \to +\infty$ as $n \to +\infty$, there is $n_4 \ge n_3$ such that $0 \le z_n(\tau, x) \le \theta/2$ for all $n \ge n_4$ and $x \in \mathbb{R}^N$, and then $z_n(t, \cdot) \to 0$ as $t \to +\infty$ uniformly in \mathbb{R}^N . In other words,

$$H_n \in \mathcal{E} \quad \text{for all } n \ge n_4.$$
 (4.15)

Let us finally estimate $\delta_1(H_n)$ for large n. Since $0 < \lambda(H_n) = \lambda(F_n) \to \alpha \omega_N R^N$ as $n \to +\infty$ by (4.7), one has

$$0 < R_{H_n} \to \alpha^{1/N} R < R' \quad \text{as } n \to +\infty, \tag{4.16}$$

where the inequality $\alpha^{1/N}R < R'$ holds by (4.4). From this, $G_n \subset B_\rho$, $|x_n| \to +\infty$ and $\rho_n < R_1/2$, there is then $n_5 \ge n_4$ such that, for every $n \ge n_5$ and $y \in \mathbb{R}^N$, one has either $G_n \cap B_{R_{H_n}}(y) = \emptyset$ or $B_{\rho_n}(x_n) \cap B_{R_{H_n}}(y) = \emptyset$. Consequently,

$$\max_{y \in \mathbb{R}^N} \lambda \big(H_n \cap B_{R_{H_n}}(y) \big) = \max \Big(\max_{y \in \mathbb{R}^N} \lambda \big(G_n \cap B_{R_{H_n}}(y) \big), \max_{y \in \mathbb{R}^N} \lambda \big(B_{\rho_n}(x_n) \cap B_{R_{H_n}}(y) \big) \Big)$$

for all $n \geq n_5$. But $\lambda(B_{\rho_n}(x_n)) = \lambda(F_n) - \lambda(G_n) < \alpha^2 \omega_N R^N < \alpha \beta \omega_N R^N$ for all $n \geq n_1$ by (4.9), while $\lambda(G_n \cap B_{R_{H_n}}) \to \beta \omega_N \alpha R^N$ as $n \to +\infty$ by (4.16) and the definition of G_n . Therefore, there is $n_6 \geq n_5$ such that

$$\max_{y \in \mathbb{R}^N} \lambda (H_n \cap B_{R_{H_n}}(y)) = \max_{y \in \mathbb{R}^N} \lambda (G_n \cap B_{R_{H_n}}(y))$$

for all $n \geq n_6$, and then

$$\lim_{n \to +\infty} \max_{y \in \mathbb{R}^N} \lambda (H_n \cap B_{R_{H_n}}(y)) = \lim_{n \to +\infty} \lambda (G_n \cap B_{R_{H_n}}) = \alpha \beta \omega_N R^N.$$

Therefore,

$$\delta_1(H_n) \to 1 - \frac{\alpha \beta \omega_N R^N}{\alpha \omega_N R^N} = 1 - \beta \text{ as } n \to +\infty.$$

From (4.8), there is then $n_7 \ge n_6$ such that

$$\delta_1(F_n) > \delta_1(H_n) > 0$$
 for all $n \ge n_7$.

Together with (4.10), (4.13) and (4.15), one concludes that, for each $n \geq n_7$, the sets $E_1 := F_n$ and $E_2 := H_n$ satisfy the desired properties of Theorem 3.6, completing its proof.

Alternate proof of Theorem 3.6 in dimensions $N \geq 2$. We assume here that $N \geq 2$. For a > 0, let

$$Q_a := \left(-\frac{a}{2}, \frac{a}{2}\right)^N$$

be the cube of \mathbb{R}^N centered at the origin and with sides of length a. For the solutions of (1.1) with the initial condition $\mathbb{1}_{Q_a}$, it follows from [33] that there is a unique threshold $a^* > 0$ between extinction and invasion, that is,

$$Q_a \in \mathcal{E} \text{ if } 0 < a < a^*, \quad Q_{a^*} \in \mathcal{T}, \quad \text{and} \quad Q_a \in \mathcal{I} \text{ if } a > a^*.$$
 (4.17)

For a > 0 and r > 0, define

$$C_{a,r} = Q_a \cap B_r$$
.

For each fixed a > 0, the family $C_{a,r}$ is continuously increasing in $L^1(\mathbb{R}^N)$ with respect to $r \in (0, a\sqrt{N}/2]$, while, for each fixed r > 0, the family $C_{a,r}$ is continuously increasing in $L^1(\mathbb{R}^N)$ with respect to $a \in (0, 2r]$. Consider

$$a = a^* + \varepsilon$$
,

with $\varepsilon > 0$. For $r \geq a\sqrt{N}/2 = (a^* + \varepsilon)\sqrt{N}/2$, one has $C_{a,r} = Q_a$, hence by (4.17) the solution of (1.1) with initial condition $\mathbb{1}_{C_{a,r}}$ converges to 1 as $t \to +\infty$ locally uniformly in \mathbb{R}^N . On the other hand, for r > 0 small enough, one has $C_{a,r} = B_r$ and extinction occurs, that is, $C_{a,r} \in \mathcal{E}$. Thus, by [33] again, there is a threshold value

$$r^*(\varepsilon) \in \left(0, \frac{a\sqrt{N}}{2}\right) = \left(0, \frac{(a^* + \varepsilon)\sqrt{N}}{2}\right)$$

such that

$$C_{a,r} \in \mathcal{E} \text{ if } 0 < r < r^*(\varepsilon), \quad C_{a,r^*(\varepsilon)} \in \mathcal{T}, \quad \text{and} \quad C_{a,r} \in \mathcal{I} \text{ if } r > r^*(\varepsilon).$$

When $\varepsilon \to 0$, we observe that $r^*(\varepsilon) \to a^*\sqrt{N}/2$. Indeed, on the one hand, $r^*(\varepsilon) < (a^* + \varepsilon)\sqrt{N}/2$ for each $\varepsilon > 0$, hence $\limsup_{\varepsilon \to 0, \varepsilon > 0} r^*(\varepsilon) \le a^*\sqrt{N}/2$. On the other hand, for each $r \in (0, a^*\sqrt{N}/2)$, one has $C_{a^*,r} \subset Q_{a^*}$ and $\lambda(Q_{a^*} \setminus C_{a^*,r}) > 0$, hence $C_{a^*,r} \in \mathcal{E}$ by [33] and then $C_{a^*+\varepsilon,r} \in \mathcal{E}$ for all $\varepsilon > 0$ small enough by Proposition 3.1. Finally, for each $r \in (0, a^*\sqrt{N}/2)$, one has $r < r^*(\varepsilon)$ for all $\varepsilon > 0$ small enough, hence $\lim\inf_{\varepsilon \to 0, \varepsilon > 0} r^*(\varepsilon) \ge a^*\sqrt{N}/2$, and therefore

$$\lim_{\varepsilon \to 0, \ \varepsilon > 0} r^*(\varepsilon) = \frac{a^* \sqrt{N}}{2}.$$
 (4.18)

When $\varepsilon > 0$ is large enough, so that $Q_{a^*+\varepsilon} \supset B_{R_1}$ with $R_1 > 0$ as in (1.6), then $r^*(\varepsilon) = R_1$. As the family $C_{a^*+\varepsilon,r}$ is nondecreasing with respect to $\varepsilon > 0$, the function $\varepsilon \mapsto r^*(\varepsilon)$ is nonincreasing in $(0, +\infty)$. This map is also continuous in $(0, +\infty)$, with similar arguments as above and by using Proposition 3.1 again. Notice also that $\omega_N = \lambda(B_1) > (2/\sqrt{N})^N$ (because $B_1 \supset Q_{2/\sqrt{N}}$ and $\lambda(B_1 \setminus Q_{2/\sqrt{N}}) > 0$, the assumption $N \ge 2$ is used here!). One can then choose $\sigma \in (1/\omega_N^{1/N}, \sqrt{N}/2)$ sufficiently close to $\sqrt{N}/2$ so that

$$0 < \lambda(Q_{a^*}) - \lambda(C_{a^*,\sigma a^*}) < \frac{\omega_N a^{*N}}{2^N}.$$

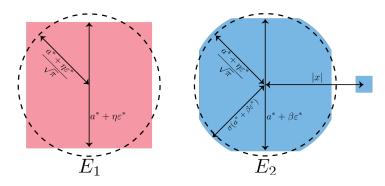


Figure 2: The two sets E_1 (in red) and E_2 (in blue) have the same Lebesgue measure $\lambda(E_1) = \lambda(E_2)$. Invasion occurs for (1.1) with initial condition $\mathbbm{1}_{E_1}$ but not with $\mathbbm{1}_{E_2}$. To understand why $0 < \delta_1(E_2) < \delta_1(E_1)$ for all |x| large enough, observe that the value of $\max_{B \in \mathcal{B}, \lambda(B) = \lambda(E_1)} \lambda(E_1 \cap B)$ corresponds in dimension N = 2 to the measure $\lambda(E_1 \cap \widetilde{B})$ of the intersection between E_1 and the disk \widetilde{B} inside the dashed circle. For |x| large enough, the value of $\max_{B \in \mathcal{B}, \lambda(B) = \lambda(E_2)} \lambda(E_2 \cap B)$ simply corresponds to the measure of $\lambda(E_2 \cap \widetilde{B})$, and is higher than $\lambda(E_1 \cap \widetilde{B}) = \max_{B \in \mathcal{B}, \lambda(B) = \lambda(E_1)} \lambda(E_1 \cap B)$, hence $\delta_1(E_2) < \delta_1(E_1)$. The details are provided below.

As $0 < \sigma < \sqrt{N}/2$, it follows from (4.18) and the above continuity and monotonicity properties of the map $\varepsilon \mapsto r^*(\varepsilon)$, that there is a unique $\varepsilon^* > 0$ such that

$$r^*(\varepsilon^*) = \sigma(a^* + \varepsilon^*).$$

Fix now $\beta \in (0,1)$ small enough so that

$$0 < \lambda(Q_{a^* + \beta \varepsilon^*}) - \lambda(C_{a^* + \beta \varepsilon^*, \sigma(a^* + \beta \varepsilon^*)}) < \frac{\omega_N a^{*N}}{2^N}, \tag{4.19}$$

define

$$0 < \frac{a^* + \beta \varepsilon^*}{\omega_N^{1/N}} < r := \sigma \left(a^* + \beta \varepsilon^* \right) < \frac{(a^* + \beta \varepsilon^*)\sqrt{N}}{2} \tag{4.20}$$

and choose $\eta \in (0, \beta)$ close enough to β so that

$$0 < \lambda(Q_{a^* + \eta \varepsilon^*}) - \lambda(C_{a^* + \beta \varepsilon^*, r}) < \frac{\omega_N a^{*N}}{2^N}, \tag{4.21}$$

and

$$\frac{a^* + \beta \varepsilon^*}{2} < \frac{a^* + \eta \varepsilon^*}{\omega_N^{1/N}} \tag{4.22}$$

(such a choice is possible because of (4.19) and because $\omega_N < 2^N$, due to the assumption $N \ge 2$). Define

$$E_1 := Q_{a^* + \eta \varepsilon^*}$$
 and $E_2 := C_{a^* + \beta \varepsilon^*, r} \cup Q^x$,

where Q^x denotes the cube of measure $\lambda(Q^x) := \lambda(Q_{a^*+\eta\varepsilon^*}) - \lambda(C_{a^*+\beta\varepsilon^*,r}) > 0$ and centered at x with |x| large enough so that $C_{a^*+\beta\varepsilon^*,r} \cap Q^x = \emptyset$. Let us check that the bounded Borel sets E_1 and E_2 fulfill the desired conclusions of Theorem 3.6, provided |x| is large enough (see Figure 2).

First of all, by definition of a^* in (4.17), invasion occurs for the solution of (1.1) with the initial condition $\mathbb{1}_{E_1}$, that is, $E_1 \in \mathcal{I}$. By definition of $r^*(\varepsilon^*)$, extinction occurs for the solution of (1.1) with the initial condition $\mathbb{1}_{C_{a^*+\beta\varepsilon^*,r}}$, as

$$r = \sigma\left(a^* + \beta \varepsilon^*\right) < \sigma\left(a^* + \varepsilon^*\right) = r^*(\varepsilon^*) \le r^*(\beta \varepsilon^*),$$

that is, $C_{a^*+\beta\varepsilon^*,r} \in \mathcal{E}$. Additionally,

$$0 < \lambda(Q^x) = \lambda(Q_{a^* + \eta \varepsilon^*}) - \lambda(C_{a^* + \beta \varepsilon^*, r}) < \frac{\omega_N a^{*N}}{2^N} < a^{*N}$$

$$(4.23)$$

by (4.21) and $\omega_N < 2^N$. As a consequence, by definition of a^* in (4.17) and the invariance of (1.1) by translation, extinction occurs for the solutions of (1.1) with the initial conditions $\mathbb{1}_{Q^x}$ (that is, $Q^x \in \mathcal{E}$) for all |x| large enough. As for H_n in (4.15) in the first proof of Theorem 3.6, for the Cauchy problem (1.1) with the initial condition $\mathbb{1}_{E_2}$, extinction occurs (that is, $E_2 \in \mathcal{E}$) for all |x| large enough.

Now, observe that $\lambda(E_1) = \lambda(E_2) > 0$ for all |x| large enough, and

$$\delta_1(E_1) = 1 - \frac{\lambda(C_{a^* + \eta \varepsilon^*, R_E})}{\lambda(E_1)},$$

with

$$\frac{a^* + \beta \varepsilon^*}{2} < R_E := \frac{a^* + \eta \varepsilon^*}{\omega_N^{1/N}} < \frac{a^* + \beta \varepsilon^*}{\omega_N^{1/N}}, \tag{4.24}$$

where the first inequality above holds because of (4.22). Lastly, for |x| large enough such that $C_{a^*+\beta\varepsilon^*,r} \cap Q^x = \emptyset$, one has

$$\lambda(Q^x) < \frac{\omega_N a^{*N}}{2^N} < \frac{\omega_N (a^* + \beta \varepsilon^*)^N}{2^N}$$
(4.25)

by (4.23). On the other hand, $R_E < r$ by (4.20) and (4.24), and then

$$\lambda(C_{a^*+\beta\varepsilon^*,r} \cap B_{R_E}) = \lambda(C_{a^*+\beta\varepsilon^*,R_E}) > \frac{\omega_N(a^*+\beta\varepsilon^*)^N}{2^N}$$
(4.26)

since $R_E > (a^* + \beta \varepsilon^*)/2$ by (4.24). For every |x| large enough, and for every $y \in \mathbb{R}^N$, one has either $C_{a^* + \beta \varepsilon^*, r} \cap B_{R_E}(y) = \emptyset$ or $Q^x \cap B_{R_E}(y) = \emptyset$. Hence

$$\max_{y \in \mathbb{R}^N} \lambda \left(E_2 \cap B_{R_E}(y) \right) = \max \left(\max_{y \in \mathbb{R}^N} \lambda \left(C_{a^* + \beta \varepsilon^*, r} \cap B_{R_E}(y) \right), \max_{y \in \mathbb{R}^N} \lambda \left(Q^x \cap B_{R_E}(y) \right) \right)$$

for all |x| large enough. Since $\lambda(Q^x) < \omega_N(a^* + \beta \varepsilon^*)^N/2^N < \lambda(C_{a^* + \beta \varepsilon^*, r} \cap B_{R_E})$ for all |x| large enough by (4.25)-(4.26), one infers that

$$\max_{y \in \mathbb{R}^N} \lambda \big(E_2 \cap B_{R_E}(y) \big) = \max_{y \in \mathbb{R}^N} \lambda \big(C_{a^* + \beta \varepsilon^*, r} \cap B_{R_E}(y) \big)$$

and then $\max_{y \in \mathbb{R}^N} \lambda(E_2 \cap B_{R_E}(y)) = \lambda(C_{a^* + \beta \varepsilon^*, r} \cap B_{R_E}) = \lambda(C_{a^* + \beta \varepsilon^*, R_E})$. Therefore,

$$\delta_1(E_2) = 1 - \frac{\lambda(C_{a^* + \beta \varepsilon^*, R_E})}{\lambda(E_1)} > 0$$

for all |x| large enough (notice that $\delta_1(E_2) > 0$ because E_2 is not a ball up to a negligible set). Finally, since $a^* + \eta \varepsilon^* < a^* + \beta \varepsilon^* < 2R_E$ by (4.24), it follows that $\lambda(C_{a^* + \eta \varepsilon^*, R_E}) < \lambda(C_{a^* + \beta \varepsilon^*, R_E})$, hence

$$0 < \delta_1(E_2) < \delta_1(E_1),$$

for all |x| large enough. The alternate proof of Theorem 3.6 in dimension $N \geq 2$ is thereby complete.

References

- [1] M. Alfaro, A. Ducrot, and G. Faye. Quantitative estimates of the threshold phenomena for propagation in reaction diffusion equations, SIAM J. Appl. Dyn. Syst. 19 (2020), 1291–1311.
- [2] L. Almeida, A. Hadon, C. Kermorvant, A. Léculier, Y. Privat, M. Strugarek, N. Vauchelet, and J. P. Zubelli. Optimal release of mosquitoes to control dengue transmission, ESAIM Proc. Surveys 67 (2020), 16-29.
- [3] D. G. Aronson and H. F. Weinberger. Multidimensional nonlinear diffusion arising in population genetics, *Adv. Math.* **30** (1978), 33–76.
- [4] H. Berestycki, F. Hamel, and N. Nadirashvili. The speed of propagation for KPP type problems. II: General domains, J. Amer. Math. Soc. 23 (2010), 1–34.
- [5] M. Bramson. Convergence of solutions of the Kolmogorov equation to travelling waves, *Memoirs Amer. Math. Soc.* 44, 1983.
- [6] Y. Du and H. Matano. Convergence and sharp thresholds for propagation in nonlinear diffusion problems, J. Europ. Math. Soc. 12 (2010), 279–312.
- [7] Y. Du and H. Matano. Radial terrace solutions and propagation profile of multistable reaction-diffusion equations over \mathbb{R}^N , https://arxiv.org/pdf/1711.00952.
- [8] Y. Du and P. Poláčik. Locally uniform convergence to an equilibrium for nonlinear parabolic equations on \mathbb{R}^N , Indiana Univ. Math. J. **64** (2015), 787–824.
- [9] A. Ducrot. On the large time behaviour of the multi-dimensional Fisher-KPP equation with compactly supported initial data, *Nonlinearity* **28** (2015), 1043–1076.
- [10] P. C. Fife and J. B. McLeod. The approach of solutions of non-linear diffusion equations to traveling front solutions, *Arch. Ration. Mech. Anal.* **65** (1977), 335–361.
- [11] R.A. Fisher. The advance of advantageous genes, Ann. Eugenics 7 (1937), 335–369.
- [12] J. Garnier, L. Roques, and F. Hamel. Success rate of a biological invasion in terms of the spatial distribution of the founding population, *Bull. Math. Biol.* **74** (2012), 453–473.
- [13] J. Gärtner. Location of wave fronts for the multi-dimensional KPP equation and Brownian first exit densities, *Math. Nachr.* **105** (1982), 317–351.
- [14] F. Hamel, J. Nolen, J.-M. Roquejoffre, and L. Ryzhik. A short proof of the logarithmic Bramson correction in Fisher-KPP equations, *Netw. Heterog. Media* 8 (2013), 275–289.
- [15] F. Hamel and L. Rossi. Spreading speeds and spreading sets of reaction-diffusion equations, https://arxiv.org/abs/2105.08344.
- [16] F. Hamel and L. Rossi. Asymptotic one-dimensional symmetry for the Fisher-KPP equation, https://arxiv.org/abs/2207.05147.
- [17] F. Hamel and L. Rossi. Spreading, flattening and logarithmic lag for reaction-diffusion equations in \mathbb{R}^N .
- [18] C. K. R. T. Jones. Spherically symmetric solutions of a reaction-diffusion equation, *J. Diff. Equations* **49** (1983), 142–169.
- [19] J. I. Kanel'. Stabilization of the solutions of the equations of combustion theory with finite initial functions, *Mat. Sb.* (N.S.) **65** (1964), 398–413.

- [20] A. N. Kolmogorov, I. G. Petrovsky, and N. S. Piskunov. Étude de l'équation de la diffusion avec croissance de la quantité de matière et son application à un problème biologique, *Bull. Univ. État Moscou, Sér. Intern. A* 1 (1937), 1–26.
- [21] E. Krätzel. Lattice points, Springer Science & Business Media, 33, Springer, 1989.
- [22] K.-S. Lau. On the nonlinear diffusion equation of Kolmogorov, Petrovsky, and Piscounov, *J. Diff. Equations* **59** (1985), 44–70.
- [23] M. A. Lewis and P. Kareiva. Allee dynamics and the spread of invading organisms, *Theo. Pop. Biol.* 43 (1993), 141–158.
- [24] H. Matano and M. Nara, Large time behavior of disturbed planar fronts in the Allen-Cahn equation, J. Diff. Equations 251 (2011), 3522–3557.
- [25] H. Matano, M. Nara, and M. Taniguchi. Stability of planar waves in the Allen-Cahn equation, *Comm. Part. Diff. Equations* **34** (2009), 976–1002.
- [26] H. Matano and P. Poláčik. Dynamics of nonnegative solutions of one-dimensional reaction-diffusion equations with localized initial data. Part I: A general quasiconvergence theorem and its consequences, Comm. Part. Diff. Equations 41 (2016), 785–811.
- [27] H. Matano and P. Poláčik. Dynamics of nonnegative solutions of one-dimensional reaction-diffusion equations with localized initial data. Part II: Generic nonlinearities, *Comm. Part. Diff. Equations* 45 (2020), 483–524.
- [28] I. Mazari, G. Nadin, and A. I. Toledo Marrero. Optimisation of the total population size with respect to the initial condition for semilinear parabolic equations: two-scale expansions and symmetrisations, *Nonlinearity* **34** (2021), 7510–7539.
- [29] C. B. Muratov and X. Zhong. Threshold phenomena for symmetric decreasing solutions of reaction-diffusion equations, *Nonlin. Diff. Equations Appl.* **20** (2013), 1519–1552.
- [30] C. B. Muratov and X. Zhong. Threshold phenomena for symmetric-decreasing radial solutions of reaction-diffusion equations, *Disc. Cont. Dyn. Syst. A* **37** (2017), 915–944.
- [31] G. Nadin and A. I. Toledo Marrero. On the maximization problem for solutions of reaction-diffusion equations with respect to their initial data, *Math. Model. Nat. Phenom.* **15** (2020), 71.
- [32] J. Nolen, J.-M. Roquejoffre, and L. Ryzhik. Convergence to a single wave in the Fisher-KPP equation, Chinese Ann. Math. Ser. B (special issue in honour of H. Brezis) 38 (2017), 629–646.
- [33] P. Poláčik. Threshold solutions and sharp transitions for nonautonomous parabolic equations on \mathbb{R}^N , Arch. Ration. Mech. Anal. 199 (2011), 69–97.
- [34] P. Poláčik. Convergence and quasiconvergence properties of solutions of parabolic equations on the real line: an overview, In: *Patterns of Dynamics*, Springer Proc. Math. Stat. 205, Springer, 2017, 172–183.
- [35] P. Poláčik. Planar propagating terraces and the asymptotic one-dimensional symmetry of solutions of semilinear parabolic equations, SIAM J. Math. Anal. 49 (2017), 3716–3740.
- [36] P. Poláčik. Propagating terraces and the dynamics of front-like solutions of reaction-diffusion equations on \mathbb{R} , Mem. Amer. Math. Soc. 264, 2020.
- [37] J.-M. Roquejoffre, L. Rossi, and V. Roussier-Michon. Sharp large time behaviour in N-dimensional Fisher-KPP equations, *Disc. Cont. Dyn. Syst. A* **39** (2019), 7265–7290.
- [38] J.-M. Roquejoffre and V. Roussier-Michon. Nontrivial large-time behaviour in bistable reaction-diffusion equations, *Ann. Mat. Pura Appl.* **188** (2009), 207–233.

- [39] J.-M. Roquejoffre and V. Roussier-Michon. Nontrivial dynamics beyond the logarithmic shift in two-dimensional Fisher-KPP equations, *Nonlinearity* **31** (2018), 3284–3307.
- [40] J.-M. Roquejoffre and V. Roussier-Michon. Sharp large time behaviour in N-dimensional reaction-diffusion equations of bistable type, J. Diff. Equations 339 (2022), 134–151.
- [41] L. Rossi. Symmetrization and anti-symmetrization in reaction-diffusion equations, *Proc. Amer. Math. Soc.* **145** (2017), 2527–2537.
- [42] L. Rossi. Stability analysis for semilinear parabolic problems in general unbounded domains, *J. Funct.* Anal. **279** (2020), 108657.
- [43] V. Roussier. Stability of radially symmetric travelling waves in reaction-diffusion equations, Ann. Inst. H. Poincaré, Anal. Non Linéaire 21 (2004), 341–379.
- [44] K. Uchiyama. The behavior of solutions of some semilinear diffusion equation for large time, *J. Math. Kyoto Univ.* **18** (1978), 453–508.
- [45] K. Uchiyama. Asymptotic behavior of solutions of reaction-diffusion equations with varying drift coefficients, Arch. Ration. Mech. Anal. 90 (1985), 291–311.
- [46] A. Zlatoš. Sharp transition between extinction and propagation of reaction, *J. Amer. Math. Soc.* **19** (2006), 251–263.