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► **To cite this version:**

Rym Smaï. Causal completion of a globally hyperbolic conformally flat spacetime. *Journal of Mathematical Physics*, 2024, 65 (10), pp.102501. 10.1063/5.0204994 . hal-04331774

HAL Id: hal-04331774

<https://cnrs.hal.science/hal-04331774v1>

Submitted on 8 Dec 2023

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Causal completion of a globally hyperbolic conformally flat spacetime

Rym SMAÏ

Abstract. In [6], Geroch, Kronheimer and Penrose introduced a way to attach ideal points to a spacetime M , defining the *causal completion* of M . They established that this is a topological space which is Hausdorff when M is globally hyperbolic. In this paper, we prove that if, in addition, M is *simply-connected* and *conformally flat*, its causal completion is a topological manifold with boundary homeomorphic to $S \times [0, 1]$ where S is a Cauchy hypersurface of M . We also introduce three remarkable families of globally hyperbolic conformally flat spacetimes and provide a description of their causal completions.

1 Introduction

An important concept in Lorentzian geometry is *causality*. Indeed, the tangent vectors to a Lorentzian manifold split into three classes: those of negative, positive and null norm, known as *timelike*, *spacelike* and *lightlike* vectors respectively. The curves whose tangent vectors are timelike or lightlike are called *causal curves*. The causal structure describes which points of the manifold can be connected, or not, by a causal curve. Lorentzian manifolds under consideration are usually oriented and time-oriented, referred to as *spacetimes*. A time-orientation provides an orientation for every causal curve, making it either future-directed or past-directed.

Among causal properties, special importance is given to *global hyperbolicity*. This is a standard assumption for spacetimes considered as cosmological models in general relativity. According to a classical theorem by Geroch [7], a spacetime is globally hyperbolic (abbrev. GH) if it admits a topological hypersurface that is intersected exactly once by every inextendible causal curve, known as a *Cauchy hypersurface*. It turns out that Cauchy hypersurfaces of a GH spacetime are homeomorphic.

In [6], Geroch, Kronheimer and Penrose defined how to attach to a spacetime M ideal points. Their construction is entirely based on the causal structure of M and formalizes the concept of *endpoints at infinity* of inextendible causal curves. Those in the future form *the future causal boundary* while those in the past form *the past causal boundary* of M . The union of M , its future and its past causal boundary define *the causal completion* of M . The authors proved that the causal completion is a topological space in which M is a dense open subset. Moreover, it is Hausdorff when M is GH. It is worth noting that, in general, the topology of the causal completion can be quite intricate (see Section 4).

In this paper, we explore the causal completion of a GH *conformally flat* spacetime of dimension $n \geq 3$. These are locally homogeneous manifolds, no longer equipped with a single Lorentzian metric but with a *conformal class* of Lorentzian metrics. The model space of these geometric structures is the Lorentzian analogue of the conformal sphere,

the so-called *Einstein universe* $Ein_{1,n-1}$; its group of conformal transformations is the linear group $O(2,n)$. Notice that there is still a notion of causality in this setting since the sign of the norm of a tangent vector is invariant under conformal changes of metrics. In Section 4, we prove the following result:

Theorem 1. *Let M be a simply-connected conformally flat GH spacetime of dimension $n \geq 3$ without conjugate points. The causal completion of M is a topological manifold with boundary, homeomorphic to $S \times [0, 1]$ where S is any Cauchy hypersurface of M .*

The assumptions *simply-connected* and *without conjugate points* (see Section 2.3 for the definition) are natural in the sense that if one of them is not satisfied, we construct easily counterexamples. In fact, Theorem 1 still holds for the larger class of *developable* spacetimes (see Section 2.3).

The proof of Theorem 1 centrally relies on the notion of the *enveloping space* of a simply-connected GH conformally flat spacetime, introduced in a previous paper [16]. We established that M can be embedded in a conformally flat spacetime $E(M) = \mathcal{B} \times \mathbb{R}$, where \mathcal{B} is a conformally flat Riemannian manifold diffeomorphic to any Cauchy hypersurface of M . This embedding is realized as the domain bounded by the graphs of two real functions, f^+ and f^- , defined on an open subset of \mathcal{B} . The space $E(M)$ is referred to as *an enveloping space of M* . The proof of Theorem 1 consists in proving that the graphs of f^+ and f^- precisely correspond to the future and past causal boundary of M .

1.1 Application: equivalence between \mathcal{C}_0 -maximality and C-maximality

There is a natural partial order relation on globally hyperbolic conformal spacetimes. Given two GH conformal spacetimes M and N , we say that N is a Cauchy-extension of M if there exists a conformal embedding f from M to N sending every Cauchy hypersurface of M on a Cauchy hypersurface of N . The map f is called *a conformal Cauchy-embedding*. A GH conformal spacetime M is said *C-maximal* if every conformal Cauchy-embedding from M to a GH conformal spacetime N is surjective. Given a GH conformal spacetime M , the existence of a C-maximal Cauchy-extension of M is *a priori* not insured. Nevertheless, it is within the category of conformally flat spacetimes. More precisely, a GH *conformally flat* spacetime M is said *\mathcal{C}_0 -maximal* if every conformal Cauchy-embedding from M to a GH *conformally flat* spacetime is surjective. A result due to C. Rossi [13] states that any conformally flat spacetime admits a \mathcal{C}_0 -maximal Cauchy-extension and furthermore, this extension is unique up to conformal diffeomorphism. We provided a new proof of this result, which involves the notion of enveloping space as detailed in [16, Sec. 5].

A natural question arises: Is the \mathcal{C}_0 -maximal extension also C-maximal? *A priori*, there is no reason for this to be true. However, we establish that it is indeed the case.

Theorem 2. *Let M be a GH conformally flat spacetime. If M is \mathcal{C}_0 -maximal, it is C-maximal.*

The converse assertion is obviously true: any C-maximal conformally flat spacetime is \mathcal{C}_0 -maximal. Thus, Theorem 2 says that the notions of \mathcal{C}_0 -maximality and C-maximality

are equivalent for conformally flat spacetimes. Hence, a \mathcal{C}_0 -maximal spacetime will be simply referred to as *maximal* (abbrev. GHM) - while keeping in mind that it is for the ordering relation defined by *conformal* Cauchy-embeddings.

This result was initially proved by Rossi in her thesis (see [12, Chap. 7, Sec. 2]). The key idea of the proof is that the C-maximality of a conformal spacetime M can be characterized by the points of its causal boundary (see Section 5.1). When M is conformally flat, we prove that the description of the causal completion of M given by Theorem 1 immediately implies Theorem 2.

1.2 Future work: complete photons

In [17], Witten proposed the problem of classifying GH spacetimes of constant curvature, and more generally conformally flat spacetimes. GH spacetimes of constant curvature have been extensively studied by several authors as Mess, Scannell, Barbot, M erigot and Bonsante. A direction still little investigated is the study of GH *conformally flat* spacetimes. The first results on this topic are attributed to C. Rossi. A fundamental result of hers states that a GHM conformally flat spacetime whose universal cover admits conjugate points is a finite quotient of the universal cover of Einstein universe (see [13, Theorem 10]).

Rossi's result completely classifies GHM conformally flat spacetimes with conjugate points. In continuation, we investigate GHM conformally flat spacetimes *without conjugate points*. An intriguing scenario arises when the universal cover of the spacetime contains *complete* photons (see Definition 9). In Section 6, we outline three notable families of such spacetimes and provide descriptions of their causal boundaries.

A strategy for studying globally hyperbolic (GH) conformally flat spacetimes with complete photons involves exploring the points of the causal boundary that serve as endpoints for complete photons, called *complete ideal points*. In a forthcoming paper, the description of the causal boundary provided by Theorem 1 will allow us to conduct an in-depth study of the complete ideal points.

Overview of the paper

Section 4.1 provides an overview of basic notions in Lorentzian geometry, Einstein universe, and conformally flat Lorentzian structures. We also recall the notion of enveloping space introduced in [16, Sec. 4]. In Section 3, we define the causal completion of a GH spacetime and we prove that it is a topological space which is Hausdorff. Section 4 is dedicated to the proof of Theorem 1 while Section 5 focuses on the proof of Theorem 2. Last but not least, we introduce in Section 6 three remarkable families of globally hyperbolic maximal conformally flat spacetimes and we describe their causal completion.

Acknowledgement

I would like to express my gratitude to my PhD advisor for the insightful discussions that made this work possible, for his valuable remarks, and for his guidance. I am

also thankful to Charles for his interest in my work and for his thoughtful comments and suggestions. This work of the Interdisciplinary Thematic Institute IRMIA++, as part of the ITI 2021-2028 program of the University of Strasbourg, CNRS and Inserm, was supported by IdEx Unistra (ANR-10-IDEX-0002), and by SFRI-STRAT'US project (ANR-20-SFRI-0012) under the framework of the French Investments for the Future Program.

2 Einstein universe and conformally flat spacetimes

This section introduces the model of conformally flat Lorentzian structures, the so-called *Einstein universe*. Let us start with some basics in Lorentzian geometry.

2.1 Preliminaries on Lorentzian geometry

We give here a short exposition on the causality of spacetimes. We direct to [11, Chap. 14] for more details.

Spacetime. A Lorentzian manifold is a smooth manifold of dimension n equipped with a non-degenerate symmetric 2-tensor g of signature $(1, n - 1)$.

In a Lorentzian manifold (M, g) , we say that a non-zero tangent vector v is *timelike*, *lightlike*, *spacelike* if $g(v, v)$ is respectively negative, zero, positive. The set of timelike vectors is the union of two convex open cones. When it is possible to make a continuous choice of a connected component in each tangent space, the manifold M is said *time-orientable*. The timelike vectors in the chosen component are said *future-directed* while those in the other component are said *past-directed*. A *spacetime* is an oriented and time-oriented Lorentzian manifold.

Future, past. In a spacetime M , a differential curve is *timelike*, *lightlike*, *spacelike* if its tangent vectors are timelike, lightlike, spacelike. It is *causal* if its tangent vectors are either timelike or lightlike.

Given a point p in M , the *future* (resp. *chronological future*) of p , denoted $J^+(p)$ (resp. $I^+(p)$), is the set of endpoints of future-directed causal (resp. timelike) curves starting from p . More generally, the future (resp. chronological future) of a subset A of M , denoted $J^+(A)$ (resp. $I^+(A)$), is the union of $J^+(a)$ (resp. $I^+(a)$) where $a \in A$. An open subset U of M is a spacetime and the intrinsic causality relations of U imply the corresponding ones in M . We denote $J^+(A, U)$ (resp. $I^+(A, U)$) the future (resp. chronological future) in the manifold U of a set $A \subset U$. Then, $I^+(A, U) \subset I^+(A) \cap U$. Dual to the preceding definitions are corresponding *past* versions. In general, *past* definitions and proofs follows from future versions (and vice versa) by reversing time-orientation.

Achronal, acausal subsets. A subset A of a spacetime M is called *achronal* (resp. *acausal*) if no timelike (resp. causal) curve intersects A more than once.

A subset A of M is said to be *edgeless* if for every $p \in A$, there exists an open neighborhood U of p such that:

- $U \cap A$ is achronal in U ;
- every causal curve contained in U joining a point of $I^-(p, U)$ to a point of $I^+(p, U)$ intersects $U \cap A$.

Example 1. The subset $A = \{(0, x) \in \mathbb{R}^{1,1}; x \in [0, 1]\}$ of the 2-dimensional Minkowski spacetime is not edgeless since the second condition is not satisfied at the points $(0, 0)$ and $(0, 1)$ as it is shown in Figure 1.

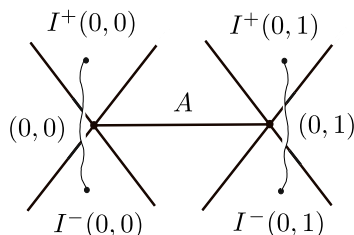


Figure 1: Achronal subset of $\mathbb{R}^{1,1}$ which is not edgeless.

Causal convexity. In Riemannian geometry, it is often useful to consider open neighborhoods which are *geodesically convex*. In Lorentzian geometry, there is, in addition, a *causal convexity* notion. A subset U of M is said *causally convex* if for every $p, q \in U$, any causal curve of M joining p to q is contained in U .

Global hyperbolicity. A spacetime M is said *strongly causal* if for every point $p \in M$ and every neighborhood U of p , there exists a neighborhood V of p contained in U , which is causally convex in M .

A spacetime M is said *globally hyperbolic* (abbrev. GH) if the two following conditions hold:

1. M is strongly causal.
2. all the intersections $J^-(p) \cap J^+(q)$, where $p, q \in M$, are compact.

By a classical theorem of Geroch [7], a spacetime M is GH if and only if it admits a topological hypersurface which is met by every inextendible causal curve exactly once. This result was improved by Bernal and Sanchez [2]: they proved that GH spacetimes admit *smooth* Cauchy hypersurfaces. It turns out that the Cauchy hypersurfaces of a GH spacetime are all homeomorphic (even diffeomorphic if they are smooth). Thus, if one of them is compact, they are all compact. In this case, the spacetime is said *Cauchy-compact*.

Conformal spacetimes. Two Lorentzian metrics g and g' on a manifold M are said *conformally equivalent* if there is a smooth function f from M to \mathbb{R} such that $g' = e^f g$. The conformal class of g is the set of Lorentzian metrics conformally equivalent to g .

Causality is a conformal notion. Indeed, given a Lorentzian manifold (M, g) , the type of a tangent vector to M depends only the conformal class of g . Then, it is relevant to consider manifolds equipped with a conformal class of Lorentzian metrics. We call *conformal spacetime* an oriented smooth manifold equipped with a conformal class of Lorentzian metrics and a time-orientation. Let us point out that, in general, geodesics are not preserved by conformal changes of metrics, except lightlike geodesics as non-parametrized curves.

2.2 Geometry of Einstein universe

Let $\mathbb{R}^{2,n}$ be the vector space \mathbb{R}^{n+2} of dimension $(n+2)$ equipped with the nondegenerate quadratic form $q_{2,n}$ of signature $(2, n)$ given by

$$q_{2,n}(u, v, x_1, \dots, x_n) = -u^2 - v^2 + x_1^2 + \dots + x_n^2$$

in the coordinate system (u, v, x_1, \dots, x_n) associated to the canonical basis of \mathbb{R}^{n+2} .

The Klein model. *Einstein universe* of dimension n , denoted by $\text{Ein}_{1,n-1}$, is the space of isotropic lines of $\mathbb{R}^{2,n}$ with respect to the quadratic form $q_{2,n}$, namely

$$\text{Ein}_{1,n-1} = \{[x] \in \mathbb{P}(\mathbb{R}^{2,n}) : q_{2,n}(x) = 0\}.$$

In practice, it is more convenient to work with the double cover of the Einstein universe, denoted by $\text{Ein}_{1,n-1}$:

$$\text{Ein}_{1,n-1} = \{[x] \in \mathbb{S}(\mathbb{R}^{2,n}) : q_{2,n}(x) = 0\}$$

where $\mathbb{S}(\mathbb{R}^{2,n})$ is the sphere of rays, namely the quotient of $\mathbb{R}^{2,n} \setminus \{0\}$ by positive homotheties.

Conformal structure. The choice of a *timelike* 2-plane of $\mathbb{R}^{2,n}$, i.e. a 2-plane on which the restriction of $q_{2,n}$ is negative definite, defines a *spatio-temporal decomposition* of Einstein universe:

Lemma 1. *Any timelike plane $P \subset \mathbb{R}^{2,n}$ defines a diffeomorphism between $\mathbb{S}^{n-1} \times \mathbb{S}^1$ and $\text{Ein}_{1,n-1}$.*

Proof. Consider the orthogonal splitting $\mathbb{R}^{2,n} = P^\perp \oplus P$ and call q_{P^\perp} and q_P the positive definite quadratic form induced by $\pm q_{2,n}$ on P^\perp and P respectively. The restriction of the canonical projection $\mathbb{R}^{2,n} \setminus \{0\}$ on $\mathbb{S}(\mathbb{R}^{2,n})$ to the set of points $(x, y) \in P^\perp \oplus P$ such that $q_{P^\perp}(x) = q_P(y) = 1$ defines a map from $\mathbb{S}^{n-1} \times \mathbb{S}^1$ to $\text{Ein}_{1,n-1}$. It is easy to check that this map is a diffeomorphism. \square

For every timelike plane $P \subset \mathbb{R}^{2,n}$, the quadratic form $q_{2,n}$ induces a Lorentzian metric g_P on $\mathbb{S}^{n-1} \times \mathbb{S}^1$ given by

$$g_P = d\sigma^2(P) - d\theta^2(P)$$

where $d\sigma^2(P)$ is the round metric on $\mathbb{S}^{n-1} \subset (P^\perp, q_{P^\perp})$ induced by q_{P^\perp} and $d\theta^2(P)$ is the round metric on $\mathbb{S}^1 \subset (P, q_P)$ induced by q_P .

An easy computation shows that if $P' \subset \mathbb{R}^{2,n}$ is another timelike plane, the Lorentzian metric $g_{P'}$ is conformally equivalent to g_P , i.e. $q_{P'}$ and $g'_{P'}$ are proportionnal by a positive smooth function on $\mathbb{S}^{n-1} \times \mathbb{S}^1$. As a result, Einstein universe is naturally equipped with a conformal class of Lorentzian metrics. Moreover, $Ein_{1,n-1} \simeq \mathbb{S}^{n-1} \times \mathbb{S}^1$ is oriented and time-oriented by the timelike vector field ∂_θ . Hence, Einstein universe is a conformal spacetime. It turns out that the causal structure of Einstein universe is trivial: any point is connected to any other point by a causal curve (see e.g. [12, Chap. 2, Cor. 2.10]). We will see that the causal structure of the universal cover of Einstein universe is more interesting.

Conformal group. The subgroup $O(2, n) \subset Gl_{n+2}(\mathbb{R})$ preserving $q_{2,n}$, acts conformally on $Ein_{1,n-1}$. When $n \geq 3$, the conformal group of $Ein_{1,n-1}$ is *exactly* $O(2, n)$. This is a consequence of the following result, which is an extension to Einstein universe, of a classical theorem of Liouville in Euclidean conformal geometry (see e.g. [5]):

Theorem 3. *Let $n \geq 3$. Any conformal transformation between two open subsets of $Ein_{1,n-1}$ is the restriction of an element of $O(2, n)$.*

Photons, lightcones and conformal spheres. Let us characterize some remarkable subsets of $Ein_{1,n-1}$.

1. A *photon* is the projectivization of a totally isotropic 2-plane of $\mathbb{R}^{2,n}$ (see e.g. [12, Chap. 2, Lemme 2.12]).
2. The *lightcone* of a point $x \in Ein_{1,n-1}$, denoted $\mathcal{C}(x)$, is the intersection of $Ein_{1,n-1}$ with the projectivization of the orthogonal of a representant $x \in \mathbb{R}^{2,n}$ of x in $\mathbb{R}^{2,n}$. Topologically, it is a double pinched torus.
3. A *conformal $(k-1)$ -sphere* is a connected component of the intersection of $Ein_{1,n-1}$ with the projectivization of a Lorentzian $(k+1)$ -plane of $\mathbb{R}^{2,n}$.

Affine charts. For every $x \in Ein_{1,n-1}$, let $M(x)$ denote the intersection of $Ein_{1,n-1}$ with the affine chart $A(x) := \{ \langle x, \cdot \rangle_{2,n} < 0 \}$ of $\mathbb{S}(\mathbb{R}^{2,n})$:

$$M(x) := \{ y \in Ein_{1,n-1} : \langle x, y \rangle_{2,n} < 0 \}.$$

Definition 1. *We call affine chart of $Ein_{1,n-1}$ any open subset of the form $M(x)$.*

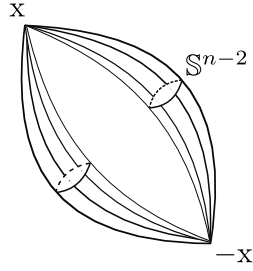


Figure 2: Lightcone of $x \in Ein_{1,2}$.

We prove that any affine chart $M(x)$ of $Ein_{1,n-1}$ is naturally a conformal Minkowski spacetime. First, we define an affine structure on $M(x)$. This depends on the choice of a representant $x \in \mathbb{R}^{2,n}$ of x but it turns out that the class of these affine structures up to homotheties is canonical in the sense that it depends only on x . Indeed, consider the map $f_x : M(x) \times M(x) \rightarrow x^\perp / \text{vect}(x)$ given by $f_x(y, z) = [y - z]$ where y and z are representant of y and z respectively such that $\langle y, x \rangle_{2,n} = \langle z, x \rangle_{2,n} = -1/2$.

Lemma 2. *The map f_x defines an affine structure on $M(x)$ of direction $x^\perp / \text{vect}(x)$. Moreover, if $x, x' \in \mathbb{R}^{2,n}$ are two distinct representant of x , there exists $\lambda \in \mathbb{R}^*$ such that $f_{x'} = \lambda f_x$. \square*

The orthogonal of x is degenerate. The kernel is the vector line in the direction of x . A supplement of $\text{vect}(x)$ in x^\perp is a subspace such that the restriction of $q_{2,n}$ is of signature $(1, n-1)$. Therefore, $q_{2,n}$ induces a quadratic form of signature $(1, n-1)$ on the quotient $x^\perp / \text{vect}(x)$. Hence, it results from Lemma 2 the following statement.

Proposition 1. *The affine chart $M(x)$ is a conformal Minkowski spacetime. \square*

The boundary of $M(x)$ in $Ein_{1,n-1}$ is the lightcone of x . Notice that the complement of the lightcone $\mathcal{C}(x)$ in $Ein_{1,n-1}$ is the disjoint union of the affine charts $M(x)$ and $M(-x)$.

Lemma 3. *The intersection of the affine chart $M(x)$ with a conformal $(k-1)$ -sphere of $Ein_{1,n-1}$ going through x is a spacelike $(k-1)$ -plane of $M(x)$.*

Proof. Let S be a conformal $(k-1)$ -sphere going through x . It is the intersection of $Ein_{1,n-1}$ with the projectivization of a Lorentzian $(k+1)$ -plane P of $\mathbb{R}^{2,n}$ containing $\text{vect}(x)$ where $x \in \mathbb{R}^{2,n}$ is a representant of x . It is easy to check that the restriction of f_x to $S \cap M(x)$ defines an affine structure with direction $x^{\perp P} / \text{vect}(x)$ where $x^{\perp P}$ denotes the orthogonal of x in P . Since P is Lorentzian, the restriction of $q_{2,n}$ to P induces a positive definite quadratic form on $x^{\perp P} / \text{vect}(x)$. The lemma follows. \square

Penrose boundary. Let $M(x)$ be an affine chart of $Ein_{1,n-1}$.

Definition 2. *The regular part of the lightcone of x is called the Penrose boundary of the affine chart $M(x)$ and is denoted $\mathcal{J}(x)$.*

The Penrose boundary of $M(x)$ is the union of two connected components $\mathcal{J}^+(x)$ and $\mathcal{J}^-(x)$ where

- $\mathcal{J}^+(x)$ fibers trivially over the sphere $S^+(x)$ of future lightlike directions at x ;
- $\mathcal{J}^-(x)$ fibers trivially over the sphere $S^-(x)$ of past lightlike directions at x ;

The fiber over a direction $[v] \in S^\pm(x)$ is the lightlike geodesic contained in $\mathcal{J}^\pm(x)$ tangent to v at x .

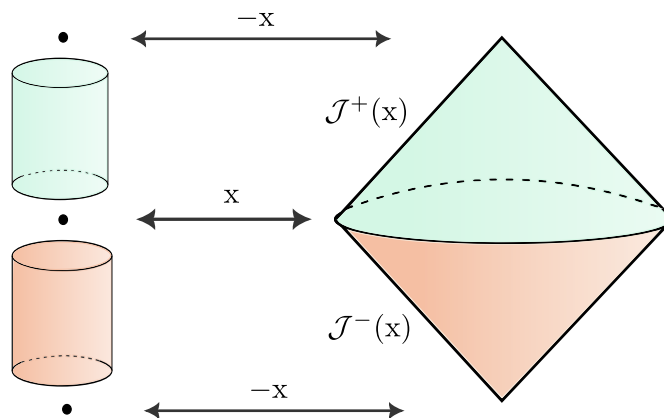


Figure 3: Penrose diagram of an affine chart $M(x)$ of $Ein_{1,2}$: the interior of the diamond represents $M(x)$; the upper and lower cones represent $\mathcal{J}^+(x)$ and $\mathcal{J}^-(x)$ respectively; the equatorial circle is identified to x and the vertices are identified to $-x$.

Penrose boundary and degenerate affine hyperplanes of an affine chart. Now, we prove that each connected component of $\mathcal{J}(x)$ is in bijection with the space of degenerate affine hyperplanes of $M(x)$. We write it for $\mathcal{J}^+(x)$ but of course it is similar for $\mathcal{J}^-(x)$.

Lemma 4. *The intersection of the affine chart $M(x)$ with the lightcone of a point $y \in \mathcal{J}^+(x)$ is a degenerate affine hyperplane of $M(x)$.*

Proof. Let $x, y \in \mathbb{R}^{2,n}$ be two representants of x and y respectively. On the one hand, $\langle y, x \rangle_{2,n} = 0$, i.e. $\text{vect}(y) \subset x^\perp$. On the other hand, since $y \notin \{x, -x\}$, the lightlike line $\text{vect}(y)$ is transverse to $\text{vect}(x)$. Therefore, the projection $[y]$ of y in the quotient $x^\perp / \text{vect}(x)$ is a non-trivial isotropic vector. It is easy to check that the map f_x defined above induces on the intersection of $M(x)$ with the lightcone of y an affine structure of direction the orthogonal of $[y]$ in $x^\perp / \text{vect}(x)$. The lemma follows. \square

Lemma 5. *Every lightlike geodesic of the Penrose boundary $\mathcal{J}^+(x)$ defines a unique lightlike direction of the affine chart $M(x)$ and vice versa.*

Proof. Let $x \in \mathbb{R}^{2,n}$ be a representant of x . Recall that the vector space associated to the affine chart $M(x)$ is $x^\perp/\text{vect}(x)$. A lightlike geodesic of $\mathcal{J}^+(x)$ is the intersection of $\text{Ein}_{1,n-1}$ with the projectivization of a totally isotropic 2-plane containing x . Therefore, a lightlike geodesic of $\mathcal{J}^+(x)$ is equivalent to the data of an isotropic vector in x^\perp transverse to $\text{vect}(x)$, in other words the data of an isotropic vector of $x^\perp/\text{vect}(x)$. \square

Lemma 6. *The intersection of the Penrose boundary $\mathcal{J}^+(x)$ with the lightcone of a point of the affine chart $M(x)$ is a section of the trivial fiber bundle $\mathcal{J}^+(x) \rightarrow S^+(x)$.*

Proof. Let $x_0 \in M(x)$. Let $x, x_0 \in \mathbb{R}^{2,n}$ be two representant of x and x_0 respectively. The intersection of $\mathcal{J}^+(x)$ with the lightcone of x_0 is a connected component of the intersection of $\text{Ein}_{1,n-1}$ with the projectivization of $x^\perp \cap x_0^\perp$. Notice that $x^\perp \cap x_0^\perp = \text{vect}(x, x_0)^\perp$. Since $\langle x, x_0 \rangle_{2,n} < 0$, the subspace $\text{vect}(x, x_0)$ is of type $(1, 1)$. Then, $\text{vect}(v, v_0)^\perp$ is of type $(1, n-1)$. It follows that the intersection of $\mathcal{J}^+(x)$ with the lightcone of x_0 is a conformal $(n-2)$ -sphere that meets every lightlike geodesic of $\mathcal{J}^+(x)$. The lemma follows. \square

Let f be the map which associates to every point $y \in \mathcal{J}^+(x)$ the intersection of the lightcone of y with the affine chart $M(x)$.

Proposition 2. *The map f is a natural bijection between $\mathcal{J}^+(x)$ and the space of degenerate affine hyperplans of the affine chart $M(x)$.*

Proof. We construct the inverse of L . Let P be a degenerate affine hyperplane of $M(x)$. It is directed by the orthogonal of a lightlike direction of $M(x)$. By Lemma 5, to this lightlike direction corresponds a unique lightlike geodesic φ of $\mathcal{J}^+(x)$. Let $x_0 \in P$. By Lemma 6, the intersection of the lightcone of x_0 with $\mathcal{J}^+(x)$ meets every lightlike geodesic of $\mathcal{J}^+(x)$ in a unique point, in particular it meets φ in a unique point p . We call g the map which sends P on p . It is easy to check that $g = f^{-1}$. \square

Universal Einstein universe. *The universal Einstein universe is the cyclic cover $\mathbb{S}^{n-1} \times \mathbb{R}$ equipped with the conformal class of $d\sigma^2 - dt^2$. Notice that in dimension $n \geq 3$, it is the universal cover of Einstein universe but this is not true in dimension $n = 2$. The timelike vector ∂_t defines a time-orientation on $\widetilde{\text{Ein}}_{1,n-1}$. Thus, $\widetilde{\text{Ein}}_{1,n-1}$ is a conformal spacetime.*

Conjugate points. Let $\pi : \widetilde{\text{Ein}}_{1,n-1} \rightarrow \text{Ein}_{1,n-1}$ and $\bar{\pi} : \widetilde{\text{Ein}}_{1,n-1} \rightarrow \text{Ein}_{1,n-1}$ be the cyclic covering maps. We denote by δ (resp. σ) : $\widetilde{\text{Ein}}_{1,n-1} \rightarrow \widetilde{\text{Ein}}_{1,n-1}$ the generator of the Galois group of π (resp. $\bar{\pi}$) defined by $\delta(x, t) = (x, t + 2\pi)$ (resp. $\sigma(x, t) = (-x, t + \pi)$).

Definition 3. *Two points p and q of $\widetilde{\text{Ein}}_{1,n-1}$ are said conjugate if $q = \sigma(p)$.*

Remark 1. *If $p, q \in \widetilde{\text{Ein}}_{1,n-1}$ are conjugate, then $\pi(p) = -\pi(q)$.*

Unlike Einstein universe, the universal Einstein universe has a rich causal structure. We give a brief description of its causal structure below. We direct to [12, Chap. 2] and [13, Sec. 2] for the proofs.

Causal curves. Causal curves of $\widetilde{Ein}_{1,n-1}$ are, up to reparametrization, the curves $(x(t), t)$ where $x : I \subset \mathbb{R} \rightarrow \mathbb{S}^{n-1}$ is a 1-Lipschitz curve on the sphere defined on an interval I of \mathbb{R} . In particular, lightlike geodesics are the causal curves for which $x : I \rightarrow \mathbb{S}^{n-1}$ is a geodesic of the sphere (see e.g. [13, Lemma 5]). Inextendible causal curves are those for which $I = \mathbb{R}$ in the previous parametrization. It is then easy to see that the inextendible lightlike geodesics of $\widetilde{Ein}_{1,n-1}$ going through a point (x_0, t_0) have common intersections at the points $\sigma^k(x_0, t_0)$, for $k \in \mathbb{Z}$; and are pairwise disjoint outside these points. This description of inextendible causal curves shows that any sphere $\mathbb{S}^{n-1} \times \{t\}$, where $t \in \mathbb{R}$, is a Cauchy hypersurface. Hence, $\widetilde{Ein}_{1,n-1}$ is globally hyperbolic.

Lightcone, future and past.

- The lightcone of a point (x_0, t_0) is the set of points (x, t) such that $d(x, x_0) = |t - t_0|$ where d is the distance on the sphere \mathbb{S}^{n-1} induced by the round metric.
- The chronological future of (x_0, t_0) : this is the set of points (x, t) of $\mathbb{S}^{n-1} \times \mathbb{R}$ such that $d(x, x_0) < t - t_0$.
- The chronological past of (x_0, t_0) : this is the set of points (x, t) of $\mathbb{S}^{n-1} \times \mathbb{R}$ such that $d(x, x_0) < t_0 - t$.

Achronal sets. Every achronal subset of $\widetilde{Ein}_{1,n-1}$ is the graph of a 1-Lipschitz real-valued function f defined on a subset of \mathbb{S}^{n-1} . Achronal embedded topological hypersurfaces of $\widetilde{Ein}_{1,n-1}$ are exactly the graphs of 1-Lipschitz real-valued functions defined on \mathbb{S}^{n-1} .

Although there is no achronal subsets in $Ein_{1,n-1}$, we can keep track of the notion of achronality in $Ein_{1,n-1}$.

Proposition 3. *Two distinct points x and y of $Ein_{1,n-1}$ can be lifted to points p and q of $\widetilde{Ein}_{1,n-1}$ which are not extremities of a causal curve if and only if the sign of $\langle x, y \rangle_{2,n}$ is negative.*

Affine charts of the universal Einstein universe. For every $p \in \widetilde{Ein}_{1,n-1}$, let $\text{Mink}_0(p)$ denote the set of points which are not causally related to p . Notice that $\text{Mink}_0(p)$ is the interior of the diamond $J(\sigma(p), \sigma^{-1}(p))$. It is easy to see that this last one does not contain conjugate points. Hence, the restriction of the covering map $\pi : \widetilde{Ein}_{1,n-1} \rightarrow Ein_{1,n-1}$ to $\text{Mink}_0(p)$ is injective. Moreover, using Proposition 3, one can prove that the image of this restriction is exactly the affine chart $M(x)$ where $x = \pi(p)$. This motivates the following definition.

Definition 4. *We call affine chart of $\widetilde{Ein}_{1,n-1}$ any open subset of the form $\text{Mink}_0(p)$.*

Besides $\text{Mink}_0(p)$, the point p defines two other affines charts (see Figure 4):

- the set of points non-causally related to $\sigma(p)$, contained in the chronological future of p , denoted $\text{Mink}_+(p)$;
- the set of points non-causally related to $\sigma^{-1}(p)$, contained in the chronological past of p , denoted $\text{Mink}_-(p)$.

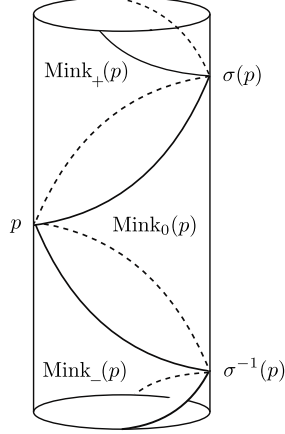


Figure 4: Affine charts defined by a point $p \in \widetilde{Ein}_{1,n-1}$ for $n = 2$.

The boundary of $\text{Mink}_0(p)$ is the union of $\partial I^+(p)$ and $\partial I^-(p)$. The regular parts of $\partial I^+(p)$ and $\partial I^-(p)$ are called *the future* and *the past Penrose boundary* of $\text{Mink}_0(p)$ and are denoted $\mathcal{J}^+(p)$ and $\mathcal{J}^-(p)$.

Remark 2. If $x = \pi(p)$, the restriction of π to $\mathcal{J}^\pm(p)$ is injective and its image is exactly $\mathcal{J}^\pm(x)$.

Given a point $q \in \mathcal{J}^+(p)$, the intersection of the past lightcone of q with $\text{Mink}_0(p)$ is a degenerate hyperplane $H(q)$. It follows that:

- the intersection of $I^-(q)$ with $\text{Mink}_0(p)$ is the chronological past of $H(q)$ in $\text{Mink}_0(p)$;
- the complement of $I^-(q)$ in $\text{Mink}_0(p)$ is the chronological future of $H(q)$ in $\text{Mink}_0(p)$.

Notice that $I^+(q)$ is disjoint from $\text{Mink}_0(p)$. Hence, the chronological future of $H(q)$ is exactly the set of points of $\text{Mink}_0(p)$ which are not causally related to q .

Regular domains. Penrose boundary is closely related to the notion of *regular domains*. We adopt here the definition of [4, Sec. 2, p. 7].

Definition 5. A future-regular domain is a non-empty convex open domain of Minkowski spacetime obtained as the intersection of strict future half-spaces bounded by a degenerate hyperplane.

We define similarly past-regular domains by reversing the time-orientation.

Let $\Lambda \subset \mathcal{J}^+(p)$. It defines naturally a convex domain in both affine charts $\text{Mink}_0(p)$ and $\text{Mink}_+(p)$:

- the set $\Omega^+(\Lambda)$ of points of $\text{Mink}_0(p)$ non-causally related to any point of Λ ;
- the set $\Omega^-(\Lambda)$ of points of $\text{Mink}_+(p)$ non-causally related to any point of Λ .

The set $\Omega^+(\Lambda)$ corresponds to the intersection of the strict future half-spaces of $\text{Mink}_0(p)$ bounded by a degenerate hyperplane of Λ and so is a *future* convex set. Similarly, $\Omega^-(\Lambda)$ is the intersection of the strict past half-spaces of $\text{Mink}_+(p)$ bounded by a degenerate hyperplane of Λ and so is a *past* convex set. Indeed, notice that $\mathcal{J}^+(p) = \mathcal{J}^-(\sigma(p))$ is the past Penrose boundary of the affine chart $\text{Mink}_0(\sigma(p)) = \text{Mink}_+(p)$. Hence, $\Omega^\pm(\Lambda)$ are regular domains if and only if they are non-empty and open. In [1, Cor. 4.11], the author shows that $\Omega^+(\Lambda)$ and $\Omega^-(\Lambda)$ are regular if and only if Λ is compact.

Remark 3. Set $x := \pi(p)$. The restriction of π to $\Omega^+(\Lambda)$ is injective (since it is contained in the affine chart $\text{Mink}_0(p)$) and, by Proposition 3, its image is the future convex domain of $M(x)$ defined as

$$\pi(\Omega^+(\Lambda)) = \{y \in M(x) : \langle y, y_0 \rangle_{2,n} < 0, \forall y_0 \in \pi(\Lambda)\}.$$

Similarly, the restriction of π to $\Omega^-(\Lambda)$ is injective and its image is the past convex domain of $M(-x)$ defined as

$$\pi(\Omega^-(\Lambda)) = \{y \in M(-x) : \langle y, y_0 \rangle_{2,n} < 0, \forall y_0 \in \pi(\Lambda)\}.$$

Example 2 (Misner domains). Regular domains defined by a conformal $(k-1)$ -sphere of the Penrose boundary are remarkable: they can be described as the chronological future/past of a spacelike $(n-k-1)$ -plane of Minkowski spacetime. These domains appeared naturally in the study of GHCM flat spacetimes (see [1, Sec. 3.2]) and are called Misner domains¹.

2.3 Conformally flat spacetimes

A spacetime is said *conformally flat* if it is locally conformal to Minkowski spacetime. In dimension $n \geq 3$, by Liouville theorem, a spacetime is conformally flat if and only if it is equipped with (G, X) -structure where $X = \widetilde{\text{Ein}}_{1,n-1}$ and $G = \text{Conf}(\widetilde{\text{Ein}}_{1,n-1})$ is its group of conformal transformation. Therefore, a conformally flat Lorentzian structure on a manifold M of dimension $n \geq 3$ is encoded by the data of a development pair (D, ρ) where $D : \widetilde{M} \rightarrow \widetilde{\text{Ein}}_{1,n-1}$ is a local diffeomorphism called *developing map* and $\rho : \pi_1(M) \rightarrow \text{Conf}(\widetilde{\text{Ein}}_{1,n-1})$ is the associated *holonomy morphism*². Let us make some remarks and introduce some vocabulary:

¹These spacetimes have been called after the mathematician Charles W. Misner since they can be seen as a generalization of the two-dimensional spacetime described by Misner in [10], namely the quotient by a boost of a half space of $\mathbb{R}^{1,1}$ bounded by a lightlike straight line.

²We direct the reader not familiar with (G, X) -structures to [8, Chapter 5].

- In general, a developing map is only a local diffeomorphism, neither injective nor surjective. When D is a global diffeomorphism, we say that the conformally flat Lorentzian structure on M is *complete*.
- A conformally flat spacetime M is said *developable* if any developing map descends to the quotient, giving a local diffeomorphism from M to $\widetilde{Ein}_{1,n-1}$.
- Two points p, q of a developable conformally flat spacetime M are said to be *conjugate* if their images under a developing map are conjugate in $\widetilde{Ein}_{1,n-1}$.

Enveloping space of a developable GH conformally flat spacetime. Let M be a developable GH conformally flat spacetime. In [16, Section 4.2.], we constructed a developable conformally flat spacetime $E(M)$ with the following properties:

- $E(M)$ fibers trivially over a conformally flat Riemannian manifold \mathcal{B} , diffeomorphic to a Cauchy hypersurface of M ;
- M embeds conformally in $E(M)$ as a causally convex open subset;
- all the conformally flat Cauchy-extensions of M embeds conformally in $E(M)$ as causally convex open subsets. In particular, the \mathcal{C}_0 -maximal extension of M is the Cauchy development of a Cauchy hypersurface of M in $E(M)$.

Such a spacetime $E(M)$ is called *an enveloping space of M* .

3 Causal completion of GH spacetimes

This section introduce the notion of *causal boundary of a spacetime*, due to Geroch-Kronheimer-Penrose [6], in the setting of GH spacetimes. Let M denote a GH spacetime.

3.1 IPs and IFs

Let $U \subset M$ be an open subset.

Definition 6. We say that U is a past set if $I^-(U) \subset U$.

The first obvious examples of past sets are chronological pasts of points, and more generally, chronological pasts of causal curves.

Definition 7. We say that U is an indecomposable past set (abbrev. IP) if U is a past set which can not be written as the union of two distinct past open subsets of M .

Lemma 7. Let U be a past open subset of M . Suppose that for every $p, q \in U$ the intersection of the chronological futures of p and q in U is non-empty. Then U is an IP.

Proof. Suppose $U = V \cup W$ where V and W are two distinct past open subsets of M . Let $v \in V \setminus W$ and $w \in W \setminus V$. Let $u \in I^+(v) \cap I^+(w) \cap U$. Then, $v, w \in I^-(u)$. Suppose $u \in V$. Since V is a past set, $I^-(u) \subset V$; hence $w \in V$. Contradiction. Similarly, if $u \in W$, we obtain $v \in I^-(u) \subset W$. Contradiction. \square

It follows immediately from Lemma 7 the following statement.

Corollary 1. *The chronological past of a causal curve of M is an IP.* □

In [6], the authors proved that conversely, any IP is the chronological past of some *timelike* curve of M (see [6, Theorem 2.1]). Their proof is based on the fact that the condition on U given in Lemma 7 is not only sufficient but also necessary.

The IFs are defined similarly for the reverse time-orientation. All the results stated above are true for the reverse time-orientation.

PIPs and TIPs. Let $P = I^-(\gamma)$ be an IP, where γ is a causal curve of M . We distinguish two cases:

1. The curve γ admits a future endpoint p in M . Then, P equals $I^-(p)$. In this case, P is called a *proper indecomposable past set* (abbrev. PIP).
2. The curve γ is inextendible in the future. In this case, P is called a *terminal indecomposable past set* (abbrev. TIP).

Similarly, the IFs split in two classes: the PIFs, namely the chronological futures of points, and the TIFs $I^+(\gamma)$ where γ is a causal curve inextendible in the past.

Definition 8. We call \hat{M} (resp. \check{M}) the set of IPs (resp. IFs).

Maximal TIPs. The inclusion defines a partial ordering relation on the set of TIPs. We say that a TIP is *maximal* if it is maximal for this ordering relation.

Proposition 4. *The set of TIPs admits at least a maximal element.*

Proof. We use Zorn's lemma which states that any partially ordered set containing upper bounds for every chain necessarily contains at least one maximal element.

Let C a chain of TIPs, i.e. a totally ordered subset of TIPs. Let P denote the union of the elements of C . We prove that P is a TIP. Clearly, P is a past-set as union of past-sets. Moreover, P is indecomposable. Indeed, if P is the union of two distinct past-sets Q and R then any TIP P_i of C is the union of the $P_i \cap Q$ and $P_i \cap R$. Hence, $P_i = P_i \cap Q$ or $P_i = P_i \cap R$; in other words $P_i \subset Q$ or $P_i \subset R$ for every TIP P_i of C . It follows that $P = Q$ or $P = R$. Lastly, suppose that P is a PIP, i.e. $P = I^-(p)$ with $p \in M$. Let P_i a TIP of C . There exists a future-inextendible causal curve γ_i such that $P_i = I^-(\gamma_i)$. Since $P_i \subset P$, we have $I^-(\gamma_i(t)) \subset I^-(p)$ for every $t \geq 0$. Hence, $\gamma_i(t) \in J^-(p)$ for every $t \geq 0$. Contradiction. □

Causal completion of M . Since M is globally hyperbolic, it is in particular *past- and future-distinguishing* (see [9, Remark 3.23]), that is the maps $p \in M \mapsto I^\pm(p)$ are injective. Therefore, the set of PIPs (PIFs) identifies with M . The TIPs (resp. TIFs) can be seen as the future (resp. past) endpoints at infinity of inextendible causal curves of M called *ideal points* in [6]. The future ideal points form *the future causal boundary*

of M while the past ideal points form the *past causal boundary* of M . The disjoint union of M and its future and past causal boundaries is called *the causal completion* of M , denoted by M^\sharp . In other words, M^\sharp is the quotient of the disjoint union of the set \hat{M} of IPs and the set \check{M} of IFs by the equivalence relation $I^+(p) \sim I^-(p)$.

Example 3 (Universal Einstein universe). *The description of inextendible causal curves of the universal Einstein universe shows that any TIP and any TIF is equal to the whole space. In other words, the future and the past causal boundary of the universal Einstein universe are both reduced to a single point.*

Example 4 (Minkowski spacetime). *In Minkowski spacetime, there are remarkable TIPs (resp. TIFs) which are the chronological futures (resp. pasts) of causal straight lines: the chronological future (resp. past) of a timelike straight line is the whole space and the chronological future (resp. past) of a lightlike straight line is equal to chronological future (resp. past) of the unique degenerate hyperplane containing this line (see e.g. [15, Chapter 1, Lemme 1.2.3]). It turns out that these are the only TIPs and TIFs of Minkowski spacetime (see e.g. [15, Annexe C.1, Exemple 2]).*

The causal boundary of Minkowski spacetime can be interpreted as the conformal boundary of Minkowski spacetime (see Section 6.2.1).

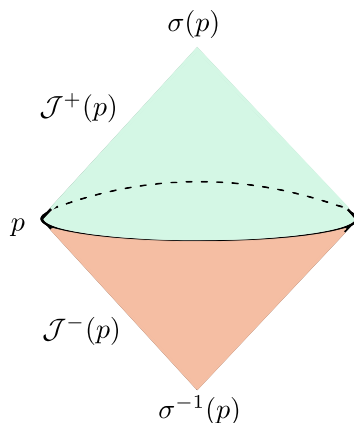


Figure 5: Causal boundary of an affine chart $\text{Mink}_0(p)$ in $\widetilde{Ein}_{1,2}$: the future causal boundary is the union of $\mathcal{J}^+(p)$ and $\sigma(p)$ and the past causal boundary is the union of $\mathcal{J}^-(p)$ and $\sigma^{-1}(p)$.

Remark 4. *Minkowski spacetime $\text{Mink}_0(p)$ admits one maximal TIP, $\sigma(p)$ and one maximal TIF, $\sigma^{-1}(p)$.*

3.2 Topology on the causal completion

We show here that the topology of M extends naturally to the causal completion of M . We start with defining a topology on the set \hat{M} of IPs. Since M is globally hyperbolic,

the topology of M coincide with Alexandrov topology, namely the topology generated by the open subsets of the form $I^+(p)$, $I^-(p)$, $M \setminus J^+(p)$, $M \setminus J^-(p)$ where p is a point of M . In other words, any open subset of M is a union of finite intersections of the previous subsets. We extend this base of topology on M to a base of topology on \hat{M} . We define two families of subsets of \hat{M} :

1. The first family is indexed by open subsets of M of the form $I^+(p)$ or $M \setminus J^-(p)$ where p is a point of M . For any open subset U of this form, we call \mathcal{O}_U the subset of \hat{M} consisting in IPs P such that $P \cap U \neq \emptyset$.
2. The second family is indexed by open subsets of M of the form $I^-(p)$ or $M \setminus J^+(p)$. For any open subset U of this form, we call \mathcal{O}'_U the subset of \hat{M} consisting in IPs P such that $\bar{P} \subset U$.

Remark 5. *The intersection of the set of PIPs with a subset \mathcal{O}_U or \mathcal{O}'_U is in bijection with U .*

Remark 6. *1. A subset of the form \mathcal{O}'_U where $U = I^-(p)$ does not contain any TIP. Indeed, suppose there exists a TIP P contained in \mathcal{O}'_U , i.e. such that $\bar{P} \subset I^-(p)$. The TIP P is the chronological past of some future causal curve γ inextendible in the future. Then, γ would be confined in $J^-(p)$. Contradiction.*

2. *A subset of the form \mathcal{O}_U where $U = M \setminus J^-(p)$ contains all the TIPs of M . Indeed, let P be a TIP. Suppose $P \notin \mathcal{O}_U$. Then, $P \cap (M \setminus J^-(p)) = \emptyset$, equivalently $P \subset J^-(p)$. Contradiction.*

Let τ be the topology on \hat{M} generated by the subsets of the form \mathcal{O}_U and \mathcal{O}'_U . It follows from Remark 5 the following statement.

Lemma 8. *The map from M to \hat{M} which sends a point $p \in M$ on the PIP $I^-(p)$ is a topological embedding. \square*

Proposition 5. *The topology τ is Hausdorff.*

Proof. Let U and V two distinct IPs. Then, $V \setminus U$ or $U \setminus V$ is non-empty. We suppose without loss of generality that $V \setminus U$ is non-empty. We prove that consequently, $V \setminus \bar{U}$ is non-empty. Suppose $V \setminus \bar{U} = \emptyset$, i.e. $V \subset \bar{U}$. Since U is a past set, the interior of \bar{U} is exactly U ³. Then, $V \subset U$, i.e. $V \setminus U = \emptyset$. Contradiction. Hence, $V \setminus \bar{U} \neq \emptyset$. Let $w \in V \setminus \bar{U}$. Set $W = I^+(w)$ and $W' = M \setminus J^+(w)$. Then, $\mathcal{O}'_{W'}$ and \mathcal{O}_W are two disjoint neighborhoods of U and V respectively. \square

Proposition 6. *The subspace of \hat{M} consisting in PIPs of M is an open subset of \hat{M} dense in \hat{M} .*

³This follows immediately from the fact that the boundary of a past set is a closed achronal topological hypersurface of M (see [11, Corollary 27, p. 415]).

Proof. By Lemma 8, the set of PIPs is open in \hat{M} . Let $P = I^-(\gamma)$ be an TIP where $\gamma : [a, b[\rightarrow M$ is a future-inextendible causal curve. Let $\{t_n\}$ be an increasing sequence of times in the interval $[a, b[$. We call P_n the PIP defined as the chronological past of $\gamma(t_n)$. We prove that $\{P_n\}$ converges to P . By Remark 6, it is sufficient to prove that any open neighborhood of P of the form \mathcal{O}_U where $U = I^+(p)$ with $p \in M$, or of the form \mathcal{O}'_U where $U = M \setminus J^+(p)$ with $p \in M$, contains all the P_n for n big enough.

Since $P_n \subset P$, it is clear that any open neighborhood \mathcal{O}'_U of P contains all the P_n . Let \mathcal{O}_U be a neighborhood of P where $U = I^+(p)$. Then, $p \in P$. Since $\{t_n\}$ is increasing, for n big enough, $p \in I^-(\gamma(t_n)) = P_n$. Hence, $P_n \cap I^+(p) \neq \emptyset$, i.e. $P_n \in \mathcal{O}_U$ for n big enough. We deduce that $\{P_n\}$ converges to P . \square

Remark 7. *The same construction with the reverse time orientation defines a topology on the set \check{M} of IFs. Propositions 5 and 6 still hold for \check{M} .*

Propositions 5 and 6 are the best we can say for a general globally hyperbolic space-time.

4 Causal completion of developable GH conformally flat spacetimes

We devote this section to the proof of our main result:

Theorem 4. *Let M be a developable conformally flat GH spacetime of dimension $n \geq 3$ without conjugate points. The causal completion of M is a topological manifold with boundary, homeomorphic to $S \times [0, 1]$ where S is any Cauchy hypersurface of M .*

Notice that the assumptions *developable* and *without conjugate points* are necessary. Indeed, if one of them is not satisfied, we can easily point out counter-examples:

- The universal Einstein universe $\widetilde{Ein}_{1,n-1}$ is developable but contains conjugate points. Each of its future and past causal boundary is reduced to a single point.
- Consider the quotient of a regular domain Ω of $\mathbb{R}^{1,2}$ (in the sense of [3]) by a discrete subgroup of isometries Γ of $\mathbb{R}^{1,2}$. In general, the action of Γ on the singular points of the boundary of Ω in $\mathbb{R}^{1,2}$ (see [3, Section 4]) is neither free nor properly discontinuous. The quotient of the causal completion of Ω by Γ is then far from being a topological manifold.

In Section 4.1, we prove Theorem 1 for causally convex open subsets of $\widetilde{Ein}_{1,n-1}$ before dealing with the general case in Section 4.2.

4.1 The case of causally convex open subsets of Einstein universe

Let Ω be a causally convex open subset of $\widetilde{Ein}_{1,n-1}$ without conjugate points. By [16, Proposition 3], there exist two 1-Lipschitz real-valued functions f^+ and f^- defined on an open subset U of \mathbb{S}^{n-1} whose extensions to ∂U coincide, such that

$$\Omega = \{(x, t) \in U \times \mathbb{R}; f^-(x) < t < f^+(x)\}.$$

Proposition 7. *The future (resp. past) causal boundary of Ω is homeomorphic to the graph of f^+ (resp. f^-).*

Proof. Let f be the map which associates to every point p in the graph of f^+ the TIP $I^-(p) \cap \Omega$ of Ω . We prove that f is bijective. Let p, q two points in the graph of f^+ such that $I^-(p) \cap \Omega = I^-(q) \cap \Omega$. Then, $I^-(I^-(p) \cap \Omega) = I^-(I^-(q) \cap \Omega)$. Hence, $I^-(p) = I^-(q)$. Since $\widetilde{Ein}_{1, n-1}$ is past-distinguishing, we deduce that $p = q$. Thus, f is injective.

Now, let P be a TIP of Ω . Then, there exists an inextensible timelike curve γ of Ω such that $P = I^-(\gamma, \Omega)$. Since Ω is causally convex, γ is the intersection of an inextensible timelike curve $\tilde{\gamma}$ with Ω . By [16, Lemma 6], $\tilde{\gamma}$ intersects the graph of f^+ exactly once in a point p . Thus, $P = I^-(p) \cap \Omega$. The map f is then surjective.

We deduce that f is bijective. It is easy to check that it is a homeomorphism. \square

Corollary 2. *The causal completion of Ω is a topological manifold with boundary homeomorphic to $S \times [0, 1]$ where S is a Cauchy hypersurface of Ω .* \square

4.2 The general case

Let M be a developable GH conformally flat spacetime without conjugate points. Let $E(M)$ be an enveloping space of M : this is a conformally flat developable spacetime which fibers trivially over a conformally flat Riemannian manifold \mathcal{B} diffeomorphic to a Cauchy hypersurface of M , in which M embeds conformally as a causally convex open subset Ω (see [16, Theorem 1]). By [16, Prop. 10], Ω is the domain bounded by the graphs of two real-valued functions f^+ and f^- defined on an open subset of \mathcal{B} :

$$\Omega = \{(x, t) \in U \times \mathbb{R}; f^-(x) < t < f^+(x)\}.$$

Moreover, the graphs of f^+ and f^- satisfy the property of being achronal in $E(M)$.

Proposition 8. *The future (resp. past) causal boundary of Ω is homeomorphic to the graph of f^+ (resp. f^-).*

Proof. Consider the map f which associates to every point p in the graph of f^+ the TIP $I^-(p) \cap \Omega$. We show that f is bijective. Since $E(M)$ is developable, it is strongly causal (see [16, Lemma 8]) and thus distinguishing (see [9, Remark 3.23]). Therefore, the same arguments used in the proof of Proposition 7 show that f is injective.

Let P be a TIP of Ω . Then, there exists an inextensible timelike curve γ of Ω such that $P = I^-(\gamma)$. Let $p_0 \in \gamma$. Since Ω is GH, the restriction of \hat{D} to $I^+(p_0, \Omega)$ is injective and its image is causally convex in $\widetilde{Ein}_{1, n-1}$ (see [12, Prop. 2.7 and Cor. 2.8, p.151]). Clearly, $\gamma \cap I^+(p_0, \Omega)$ is a future-inextensible timelike curve of $I^+(p_0, \Omega)$, denoted γ_0 . Therefore, by Proposition 7, γ_0 admits a future endpoint p in the intersection of the graph of f^+ with $I^+(p_0, \Omega)$. Then, $P = I^-(p) \cap \Omega$. Thus, f is surjective. It is easy to check that f and its inverse are continuous. \square

Theorem 4 follows immediately from Proposition 8.

5 \mathcal{C}_0 -maximality and C-maximality

We devote this section to the proof of Theorem 2: we prove that any \mathcal{C}_0 -maximal spacetime is C-maximal. In Section 5.1, we recall a criterion of C-maximality involving the causal boundary. Then, in Section 5.2, we use this criterion to prove Theorem 2.

5.1 Criterion of C-maximality

Let f be a conformal Cauchy-embedding between two globally hyperbolic conformal spacetimes M and N .

Proposition 9 ([13, Lemma 10]). *The boundary of $f(M)$ in N is the union of two disjoint closed achronal edgeless sets (eventually empty) $\partial^+ f(M)$ and $\partial^- f(M)$ such that*

$$I^-(\partial^+ f(M)) \cap I^+(\partial^- f(M)) \subset f(M).$$

There is subtle relation between the boundary of $f(M)$ in N and the causal boundary of M :

Proposition 10. *Suppose $\partial^+ f(M)$ is non-empty. Then, for every point p of $\partial^+ f(M)$, the pre-image under f of the intersection of the chronological past of p with $f(M)$ is a TIP of M .*

There is a similar statement for $\partial^- f(M)$ with the reverse time-orientation.

Proof. Set $P = f^{-1}(I^-(p) \cap f(M))$.

- *P is a past set, i.e. $P = I^-(P)$:* Since P is open, $P \subset I^-(P)$. Conversely, let $x \in I^-(P)$. There exists $y \in P$ such that $x \in I^-(y)$. Then, since f is conformal, $f(x) \in I^-(f(y))$. Moreover, since $y \in P$, we have $f(y) \in I^-(p)$. By transitivity, we get $f(x) \in I^-(p)$. Hence, $x \in P$. Thus, $I^-(P) \subset P$.
- *P is indecomposable:* Suppose P is the union of two distinct past sets Q and R . Then, $f(P) = f(Q) \cup f(R)$, i.e. $I^-(p) \cap f(M) = f(Q) \cup f(R)$. Hence, $I^-(I^-(p) \cap f(M)) = I^-(f(Q) \cup f(R)) = I^-(f(Q)) \cup I^-(f(R))$. It is easy to see that $I^-(I^-(p) \cap f(M)) = I^-(p)$. Thus, $I^-(p) = I^-(f(Q)) \cup I^-(f(R))$. Since $I^-(p)$ is indecomposable, we deduce that $I^-(p) = I^-(f(Q))$ or $I^-(p) = I^-(f(R))$. Without loss of generality, we suppose that $I^-(p) = I^-(f(Q))$. Thus, $I^-(p) \cap f(M) = I^-(f(Q)) \cap f(M)$. Since $f(M)$ is causally convex, the intersection $I^-(f(Q)) \cap f(M)$ is exactly the chronological past of $f(Q)$ in $f(M)$. But $f(Q)$ is a past set in $f(M)$. Thus, $I^-(f(Q)) \cap f(M) = f(Q)$. Hence, $I^-(p) \cap f(M) = f(Q)$. It follows that $P = Q$. In other words, P is indecomposable.
- *P is terminal:* Suppose there exists $q \in M$ such that $P = I^-(q)$. Then, we have $I^-(p) \cap f(M) = f(I^-(q))$. Since $f(M)$ is causally convex in N (see [13, Lemma 8]), $f(I^-(q)) = I^-(f(q)) \cap f(M)$. Hence, $I^-(p) \cap f(M) = I^-(f(q)) \cap f(M)$. Then, $I^-(I^-(p) \cap f(M)) = I^-(f(q) \cap f(M))$. Hence, $I^-(p) = I^-(f(q))$. Since M is globally hyperbolic, it follows that $p = f(q)$. Contradiction.

□

Corollary 3. *The map which associates to every point p of $\partial^+ f(M)$ the TIP of M defined by $f^{-1}(I^-(p) \cap f(M))$ is injective.*

Proof. Let $p, q \in \partial^+ f(M)$ such that $f^{-1}(I^-(p) \cap f(M)) = f^{-1}(I^-(q) \cap f(M))$. Then, $I^-(p) \cap f(M) = I^-(q) \cap f(M)$. Thus, $I^-(I^-(p) \cap f(M)) = I^-(I^-(q) \cap f(M))$, i.e. $I^-(p) = I^-(q)$. Hence, $p = q$. □

A consequence of the description above is the following criterion of C-maximality.

Proposition 11 (Criterion of maximality). *Let M be a globally hyperbolic conformal spacetime. Suppose that M admits a non-compact Cauchy hypersurface S . Then, if the intersection of S with any TIP and any TIF of M is non-compact, M is C-maximal.*

Proof. Let f be a conformal Cauchy embedding from M to a globally hyperbolic spacetime N . Let S be a Cauchy hypersurface of M . Suppose that f is not surjective. Then $\partial^+ f(M)$ (or $\partial^- f(M)$) is non-empty. Let $p \in \partial^+ f(M)$. By Proposition 10, $f^{-1}(I^-(p) \cap f(M))$ is a TIP of M . The intersection of this TIP with S is not compact. But, its image, equal to $I^-(p) \cap f(S)$ is compact. Contradiction. The proposition follows. □

5.2 \mathcal{C}_0 -maximality implies C-maximality

Let M be a globally hyperbolic conformally flat spacetime. We suppose that M is \mathcal{C}_0 -maximal. In what follows, we prove that M is C-maximal.

By [16, Corollary 6], the universal cover of M , denoted \tilde{M} , is \mathcal{C}_0 -maximal. If \tilde{M} is Cauchy-compact, it is conformally equivalent to the universal cover of Einstein universe (see [13, Theorem 9]).

Proposition 12. *The spacetime $\widetilde{Ein}_{1,n-1}$ is C-maximal.*

Proof. Let f be a conformal Cauchy-embedding from $\widetilde{Ein}_{1,n-1}$ in a globally hyperbolic spacetime N . Suppose that f is not surjective. Then, $\partial^+ f(M)$ (or $\partial^- f(M)$) is non-empty. By Corollary 3, there is an injection from $\partial^+ f(M)$ to the future causal boundary of $\widetilde{Ein}_{1,n-1}$. This last one is reduced to a single point. Thus, $\partial^+ f(M)$ is equal to this point. But, $\partial^+ f(M)$ is edgeless (see Proposition 9). Contradiction. Hence, $\widetilde{Ein}_{1,n-1}$ is C-maximal. □

Now, suppose that \tilde{M} admits a non-compact Cauchy hypersurface. Let $E(\tilde{M})$ be an enveloping space of \tilde{M} ; it fibers trivially over a conformally flat Riemannian manifold \mathcal{B} . The spacetime \tilde{M} embeds conformally in $E(\tilde{M})$ as a causally convex open subset Ω . Let f^+, f^- two 1-Lipschitz real-valued functions defined on an open subset U of \mathcal{B} such that

$$\Omega = \{(x, t) \in U \times \mathbb{R}; f^-(x) < t < f^+(x)\}.$$

The \mathcal{C}_0 -maximality of Ω is characterized by the following property of the graphs of f^+ and f^- (see [16, Prop. 18]):

Fact 1. *For every point p in the graph of f^+ , there exists a past-directed lightlike geodesic starting from p , entirely contained in the graph of f^+ and with no past endpoint in the graph of f^+ .*

Proposition 13. *The causally convex open subset Ω of $E(\tilde{M})$ is C-maximal.*

Proof. We use the criterion given by Proposition 11. Let S be a Cauchy hypersurface of Ω . Let P be a TIP of Ω . By Proposition 8, there exists a unique point p in the graph of f^+ such that $P = I^-(p) \cap \Omega$. By Fact 1, there exists a past-directed lightlike geodesic starting from p , entirely contained in the graph of f^+ . Therefore, $P \cap S$ is not compact. Similarly, the intersection of any TIF of Ω with S is not compact. Thus, M is C-maximal. \square

Proof of Theorem 2. Let f be a conformal Cauchy-embedding from M to a globally hyperbolic spacetime N . It is easy to see that the lift \tilde{f} of f is a Cauchy-embedding from \tilde{M} to \tilde{N} . It follows from Proposition 13 that \tilde{f} is surjective. Thus, f is surjective. In other words, M is C-maximal. \square

From now on, by Theorem 2, we simply say that a conformally flat spacetime is *maximal* (while keeping in mind that it is for the ordering relation defined by *conformal* Cauchy-embeddings).

6 Complete photons

In [13], Rossi proved that any developable GHM conformally flat spacetime containing conjugate points is conformally equivalent to $\widetilde{Ein}_{1,n-1}$. As a consequence, any GHM conformally flat spacetime whose universal cover admits conjugate points is a finite quotient of $\widetilde{Ein}_{1,n-1}$.

In this section, we introduce three remarkable families of GHM conformally flat spacetimes obtained as the quotient of a causally convex open subset $\Omega \subsetneq \widetilde{Ein}_{1,n-1}$ by a discrete group of conformal transformations of $\widetilde{Ein}_{1,n-1}$. Rossi's result insures that the domains Ω do not contain conjugate points. Nevertheless, we prove that they satisfy the remarkable property of containing *complete* photons:

Definition 9. *Let M be a developable conformally flat spacetime. A photon of M is said **complete** if it develops on a segment of photon of $\widetilde{Ein}_{1,n-1}$ connecting two conjugate points.*

The GHCM conformally flat spacetimes that we describe typically arise from the data of some appropriate P_1 -Anosov representation of a Gromov hyperbolic group into $O(2, n)$.

6.1 Anosov representations and GHCM conformally flat spacetimes

We recall here the notion of P_1 -Anosov representation in $O_0(2, n)$ and how they produce examples of GHCM conformally flat spacetimes.

Let Γ be a Gromov hyperbolic group and let ρ be a representation of Γ in $O_0(2, n)$. The notion of P_1 -Anosov representation involves two dynamics. On the one hand, the north-south dynamic on the Gromov boundary of Γ , denoted $\partial_\infty\Gamma$; and on the other hand, the dynamic on $E\text{in}_{1, n-1}$ under the action of some sequences of $O(2, n)$, namely P_1 -divergent sequences. A sequence $\{g_i\}$ of $O(2, n)$ is said P_1 -divergent if it is divergent - i.e. leaves every compact set - and if there exist a subsequence $\{g_{j_i}\}$, an attracting point p_+ and a repulsing point p_- in $E\text{in}_{1, n-1}$ such that $\{g_{j_i}\}$ converges uniformly to p_+ on every compact set disjoint from the lightcone of p_- .

Definition 10. *The representation $\rho : \Gamma \rightarrow O_0(2, n)$ is said P_1 -Anosov if*

1. *all sequences of $\rho(\Gamma)$ are P_1 -divergent;*
2. *there exists a continuous ρ -equivariant map $\xi : \partial_\infty\Gamma \rightarrow E\text{in}_{1, n-1}$ which is*
 - (a) *transverse meaning that for every pair (η, η') of distinct points of $\partial_\infty\Gamma$, the points $\xi(\eta)$ and $\xi(\eta')$ are not connected by a lightlike geodesic;*
 - (b) *dynamics-preserving meaning that if $\eta \in \partial_\infty\Gamma$ is an attracting point of $\gamma \in \Gamma$, then $\xi(\eta) \in E\text{in}_{1, n-1}$ is an attracting point of $\rho(\gamma)$.*

Any P_1 -Anosov representation preserves a specific closed subset of $E\text{in}_{1, n-1}$ called *limit set*:

- *The limit set of ρ , denoted Λ_ρ , is the set of all attracting points of the elements of $\rho(\Gamma)$. It coincides with the image of the boundary map $\xi : \partial_\infty\Gamma \rightarrow E\text{in}_{1, n-1}$.*
- *We say that the representation ρ is *negative* if the limit set Λ_ρ lifts to an *acausal* subset of $\widetilde{E\text{in}}_{1, n-1}$.*

In a previous paper, we proved the following result.

Theorem 5 ([14, Theorem 5.1]). *Any P_1 -Anosov representation ρ of a Gromov hyperbolic group Γ in $O_0(2, n)$ with negative limit set which is not a topological $(n-1)$ -sphere, is the holonomy of a GHCM conformally flat spacetime M_ρ of dimension n .*

The spacetime M_ρ is constructed as follow. We define *the invisible domain* $\Omega(\Lambda_\rho)$ of the limit set Λ_ρ :

$$\Omega(\Lambda_\rho) := \{\xi \in E\text{in}_{1, n-1} : \langle \xi, \xi_0 \rangle_{2, n} < 0 \ \forall \xi_0 \in \Lambda_\rho\}.$$

In the terminology of [16, Def. 9], $\Omega(\Lambda_\rho)$ is *the dual* of Λ_ρ . It corresponds to the set of points of $E\text{in}_{1, n-1}$ which are not causally related to any point of the limit set in the following sense:

Let $\tilde{\Lambda}_\rho$ be a lift of Λ_ρ in $\widetilde{Ein}_{1,n-1}$. We call $\Omega(\tilde{\Lambda}_\rho)$ the set of points of $\widetilde{Ein}_{1,n-1}$ which are not causally related to any point of $\tilde{\Lambda}_\rho$:

$$\Omega(\tilde{\Lambda}_\rho) = \widetilde{Ein}_{1,n-1} \setminus (J^+(\Lambda_\rho) \cup J^-(\Lambda_\rho)).$$

The restriction of the projection $\pi : \widetilde{Ein}_{1,n-1} \rightarrow Ein_{1,n-1}$ to $\Omega(\tilde{\Lambda}_\rho)$ is injective and its image is exactly $\Omega(\Lambda_\rho)$ (see [14, Lemmas 5.11 & 5.12]).

We prove that the action of $\rho(\Gamma)$ on $\Omega(\Lambda)$ is free and properly discontinuous (see [14, Sec. 5.2]). The spacetime M_ρ is then the quotient of $\Omega(\Lambda_\rho)$ by $\rho(\Gamma)$ (see [14, Sec. 5.3]).

6.2 GHCM conformally flat spacetimes with complete photons

Throughout this section, Λ denotes a closed negative subset of $Ein_{1,n-1}$ - which typically arises as the limit set of some P_1 -Anosov representation in $O(2, n)$. We describe the invisible domain $\Omega(\Lambda)$ in the following cases:

1. Λ is contained in a Penrose boundary $\mathcal{J}^+(x)$ where $x \in Ein_{1,n-1}$, defining *black-white holes*;
2. Λ is a conformal sphere of dimension $0 \leq k \leq n - 3$, defining *Misner domains of Einstein universe*;
3. Λ is strictly contained in a conformal sphere of dimension $0 \leq k \leq n - 3$, defining *Misner extensions*.

6.2.1 Black-white holes

Let $x \in Ein_{1,n-1}$. Suppose that Λ is contained in $\mathcal{J}^+(x)$ and that it is not a topological $(n - 2)$ -sphere (i.e. it is not a section of $\mathcal{J}^+(x)$).

Horizons. The intersection of $\Omega(\Lambda)$ with $\mathcal{J}^+(x)$ is the union of the lightlike geodesics with extremities x and $-x$ disjoint from Λ . Every connected component of this union is called a **horizon**. The assumption that Λ is not a topological $(n - 2)$ -sphere insures the existence of horizons. By definition, each horizon is foliated by complete photons.

Black hole, white hole. $\mathcal{J}^+(x)$ is the future Penrose boundary of the affine chart $M(x)$ and the past Penrose boundary of $M(-x)$. Hence,

1. the intersection of $\Omega(\Lambda)$ with $M(x)$ is a future-regular domain of $M(x)$, denoted $\Omega^+(\Lambda)$;
2. the intersection of $\Omega(\Lambda)$ with $M(-x)$ is a past-regular domain of $M(-x)$, denoted $\Omega^-(\Lambda)$.

The domain $\Omega^-(\Lambda)$ satisfies the property that no photon of $\Omega(\Lambda)$ going through a point of $\Omega^-(\Lambda)$ can escape from $\Omega^-(\Lambda)$ in the future, it can only escape through one of the horizons in the past. The domain $\Omega^+(\Lambda)$ satisfies a similar property for the reverse time-orientation. For this reason, $\Omega^-(\Lambda)$ is called a *black hole* and $\Omega^+(\Lambda)$ a *white-hole*.

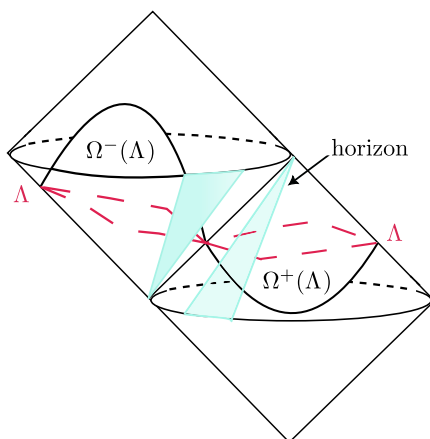


Figure 6: Black-white hole domain of $Ein_{1,2}$.

Definition 11. *The invisible domain $\Omega(\Lambda)$ is called a black-white hole domain of $Ein_{1,n-1}$. The points x and $-x$ are the endpoints of the horizons of $\Omega(\Lambda)$.*

Causal boundary. The future causal boundary of $\Omega(\Lambda)$ is the boundary of $\Omega^-(\Lambda)$ in $M(-x)$ while its past causal boundary is the boundary of $\Omega^+(\Lambda)$ in $M(x)$.

Black-white holes conformally flat spacetimes. Let Γ be a discrete group of $O_0(2, n)$ preserving $\Omega(\Lambda)$ such that

1. the action of Γ on $\Omega(\Lambda)$ is free and properly discontinuous;
2. Γ fixes x .

The point 1. insures that the quotient $M := \Omega(\Lambda) \backslash \Gamma$ is a conformally flat spacetime. The point 2. insures that Γ preserves the decomposition of $\Omega(\Lambda)$ as the disjoint union of the black hole $\Omega^-(\Lambda)$, the horizons H_i , and the white hole $\Omega^+(\Lambda)$. As a consequence, the conformally flat spacetime M can be written as the disjoint union of the black hole $\mathcal{B} := \Omega^-(\Lambda) \backslash \Gamma$, the horizons $\mathcal{H}_i := H_i \backslash \Gamma$, and the white hole $\mathcal{W} := \Omega^+(\Lambda) \backslash \Gamma$. This motivates the following definition.

Definition 12. *We call black-white hole any GHCM conformally flat spacetime obtained as the quotient of a black-white hole domain $\Omega(\Lambda)$ of $Ein_{1,n-1}$ by a discrete subgroup of $O(2, n)$ fixing the endpoints of the horizons of $\Omega(\Lambda)$.*

Example 5. *The data of a negative P_1 -Anosov representation ρ of a Gromov hyperbolic group Γ in $O_0(2, n)$ fixing a point $x \in Ein_{1,n-1}$ such that the limit set Λ_ρ*

1. *is contained in $\mathcal{J}^+(x)$,*
2. *is not a topological $(n - 2)$ -sphere;*

defines a black-white hole $\Omega(\Lambda_\rho) \backslash \rho(\Gamma)$ (see [14, Theorem 5.1]).

6.2.2 Conformally flat Misner spacetimes

Suppose that Λ is a $(\ell - 1)$ -conformal sphere where $\ell \in \mathbb{N}^*$ such that $\ell < n - 1$. It is defined by the data of a Lorentzian subspace $\mathbb{R}^{1,\ell} \subset \mathbb{R}^{2,n}$. Consider the orthogonal splitting $\mathbb{R}^{2,n} = \mathbb{R}^{1,\ell} \oplus^\perp \mathbb{R}^{1,k}$ where $k \in \mathbb{N}^*$ such that $k + \ell = n$. We denote $q_{1,\ell}$ and $q_{1,k}$ the restrictions of the quadratic form $q_{2,n}$ to $\mathbb{R}^{1,\ell}$ and $\mathbb{R}^{1,k}$ respectively; notice that $q_{2,n} = q_{1,\ell} + q_{1,k}$. The sphere Λ is a connected component of the projectivization of the quadric

$$\{(x, 0) : \in \mathbb{R}^{1,\ell} \oplus \mathbb{R}^{1,k} : q_{1,\ell}(x) = 0\}.$$

Moreover, the projectivization of the quadric

$$\{(0, y) : \in \mathbb{R}^{1,\ell} \oplus \mathbb{R}^{1,k} : q_{1,k}(y) = 0\}$$

is the disjoint union of two antipodal conformal $(k-1)$ -spheres \mathbb{S}_+^{k-1} and \mathbb{S}_-^{k-1} of $Ein_{1,n-1}$. Notice that Λ is a connected component of the intersection of the lightcones of the points of \mathbb{S}_+^{k-1} .

Homogeneous model of the invisible domain. Let \mathcal{C}_ℓ be the connected component of the causal cone of $\mathbb{R}^{1,\ell}$ defining Λ . In other words, Λ is the projectivization of the boundary of \mathcal{C}_ℓ in $\mathbb{R}^{1,\ell}$. Let

$$\mathbb{H}^\ell = \{x \in \mathcal{C}_\ell : q_{1,\ell}(x) = -1\}$$

be the hyperbolic space of dimension ℓ and let

$$dS_{1,k-1} = \{y \in \mathbb{R}^{1,k} : q_{1,k}(y) = 1\}$$

be the de Sitter space of dimension k .

Proposition 14. *The invisible domain $\Omega(\Lambda)$ is conformally equivalent to the homogeneous space $\mathbb{H}^\ell \times dS_{1,k-1}$.*

Proof. We denote by $\langle \cdot, \cdot \rangle_{1,\ell}$ the bilinear form on $\mathbb{R}^{1,\ell}$ associated to the quadratic form $q_{1,\ell}$. Let $[x : y] \in Ein_{1,n-1}$. We have

$$\begin{aligned} [x : y] \in \Omega(\Lambda) &\Leftrightarrow \langle (x; y), (z; 0) \rangle_{2,n} < 0, \quad \forall [z : 0] \in \Lambda \\ &\Leftrightarrow \langle x, z \rangle_{1,\ell} < 0, \quad \forall z \in \partial\mathcal{C}_\ell \end{aligned}$$

where $\partial\mathcal{C}_\ell$ denotes the boundary of \mathcal{C}_ℓ in $\mathbb{R}^{1,\ell}$. This last condition can be rephrase by saying that x is in the intersection of the strict future half-spaces of $\mathbb{R}^{1,\ell}$ bounded by the degenerate hyperplanes $\langle \cdot, z \rangle_{1,\ell} = 0$ where $z \in \partial\mathcal{C}_\ell$. It is easy to see that this intersection is exactly the cone \mathcal{C}_ℓ . Hence, $[x : y] \in \Omega(\Lambda)$ if and only if $x \in \mathcal{C}_\ell$. Up to rescaling, one can suppose that $q_{1,\ell}(x) = -1$, i.e. $x \in \mathbb{H}^\ell$. Since $0 = q_{2,n}(x; y) = q_{1,\ell}(x) + q_{1,k}(y)$, we deduce that $q_{1,k}(y) = 1$, i.e. $y \in dS_{1,k-1}$. The proposition follows. \square

Black-white hole decomposition. Fix a point $x \in \mathbb{S}_-^{k-1}$. Then, the sphere Λ is contained in $\mathcal{J}^+(x)$. Hence, the description presented in Section 6.2.1 still holds here:

The invisible domain $\Omega(\Lambda)$ is the disjoint union of a black hole B , horizons H_i foliated by complete photons and a white hole W . The fact that Λ is a conformal $(\ell - 1)$ -sphere implies that B and W are past and future Misner domains (see Lemma 3). Moreover, there is a single horizon when $\ell < n - 1$ and there are exactly two horizons when $\ell = n - 2$ (see Figure 7).

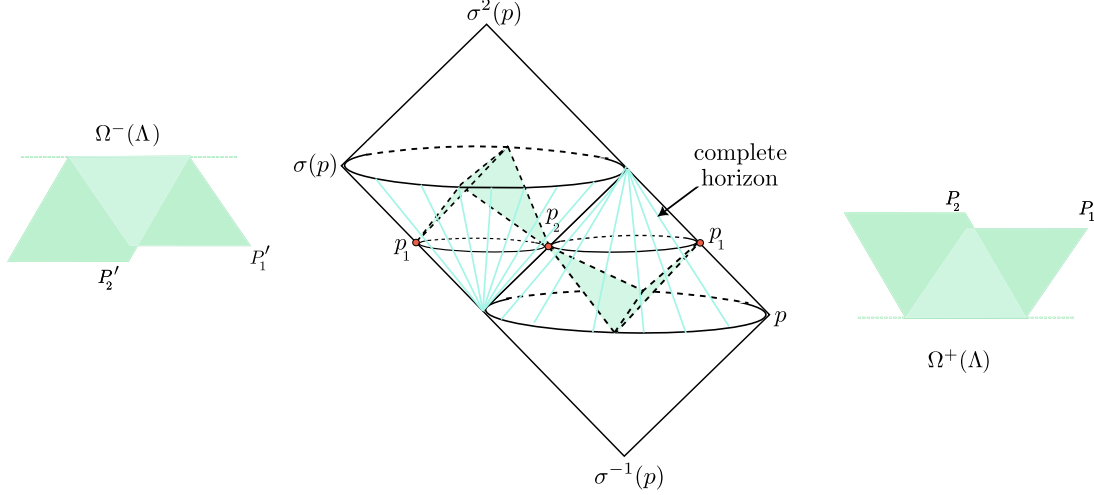


Figure 7: Black-white hole decomposition of a Misner domain of $Ein_{1,2}$. The sphere Λ is of dimension 0: this is the union of two non-causally related points p_1 and p_2 of $Ein_{1,2}$.

In short, every point of \mathbb{S}_+^{k-1} defines a decomposition of the invisible domain as the union of two Misner domains separated by one or two horizons. Since \mathbb{S}_-^{k-1} is antipodal to \mathbb{S}_+^{k-1} , this statement also holds for the points of \mathbb{S}_+^{k-1} . This motivates the following definition.

Definition 13. We call Misner domain of Einstein universe any open subset of the form $\Omega(\Lambda)$ where Λ is a conformal $(\ell - 1)$ -sphere with $0 < \ell < n - 1$.

Causal boundary. The future/past causal boundary of $\Omega(\Lambda)$ is the union of \mathbb{S}_\pm^{k-1} and the union of the lightlike geodesics joining a point of \mathbb{S}_\pm^{k-1} to a point of Λ . The set of maximal TIPs is exactly the sphere \mathbb{S}_+^{k-1} and the set of maximal TIFs is exactly the sphere \mathbb{S}_+^{k-1} (see Figure 7).

Conformally flat Misner spacetimes. Let Γ be a subgroup preserving the conformal sphere Λ . *A priori*, Γ does not fix Λ point by point, so the black-white decomposition of $\Omega(\Lambda)$ described above is not preserved by Γ . However, since Γ preserves Λ , it preserves \mathbb{S}_\pm^k . Hence, every element $\gamma \in \Gamma$ sends a black-white hole decomposition of $\Omega(\Lambda)$ on a conformally equivalent black-white hole decomposition.

Definition 14. We call conformally flat Misner spacetime any GHCM conformally flat spacetime obtained as the quotient of a Misner domain $\Omega(\Lambda)$ of Einstein universe by a discrete subgroup Γ of $O(2, n)$ preserving Λ .

Example 6. Let $O_0(1, \ell) \times O_0(1, k)$ the subgroup of $O_0(2, n)$ preserving the splitting $\mathbb{R}^{2, n} = \mathbb{R}^{1, \ell} \oplus \mathbb{R}^{1, k}$. Consider a negative P_1 -Anosov representation ρ of a Gromov hyperbolic group Γ in $O_0(2, n)$ defined by a pair (ρ_ℓ, ρ_k) where

1. $\rho_\ell : \Gamma \rightarrow O_0(1, \ell)$ is a cocompact representation whose limit set is the conformal $(\ell - 1)$ -sphere Λ ;
2. $\rho_k : \Gamma \rightarrow O_0(1, k)$ is a relatively compact representation, i.e. the image $\rho_k(\Gamma)$ is contained in a compact of $O_0(1, k)$.

The limit set Λ_ρ of the representation $\rho = (\rho_\ell, \rho_k)$ is the conformal $(\ell - 1)$ -sphere Λ . Hence, the quotient $\Omega(\Lambda_\rho) \backslash \rho(\Gamma)$ is a conformally flat Misner spacetime.

Extensions of conformally flat Misner spacetimes. We ask the following question:

Is there a discrete subgroup Γ of $O_0(2, n)$ preserving the sphere Λ and a causally convex open subset Ω' of $\text{Ein}_{1, n-1}$ containing strictly $\Omega(\Lambda)$?

The answer to this question is yes! Anosov representations give examples of such subgroups Γ . Indeed, let ρ be a negative P_1 -Anosov representation of a Gromov hyperbolic group Γ in $O_0(2, n)$ defined by a pair (ρ_ℓ, ρ_k) where

1. $\rho_\ell : \Gamma \rightarrow O_0(1, \ell)$ is a convex cocompact representation such that the limit set Λ_{ρ_ℓ} is *strictly* contained in the conformal $(\ell - 1)$ -sphere Λ ;
2. $\rho_k : \Gamma \rightarrow O_0(1, k)$ is a relatively compact representation, i.e. the image $\rho_k(\Gamma)$ is contained in a compact of $O_0(1, k)$.

The limit set Λ_ρ of the representation ρ is exactly Λ_{ρ_ℓ} . Since $\Lambda_\rho \subsetneq \Lambda$, the invisible domain $\Omega(\Lambda_\rho)$ contains strictly $\Omega(\Lambda)$. Moreover, since $\rho(\Gamma)$ preserves Λ , it preserves $\Omega(\Lambda)$. Since $\Omega(\Lambda) \subset \Omega(\Lambda_\rho)$, the action of $\rho(\Gamma)$ on $\Omega(\Lambda)$ is free and properly discontinuous. The quotient spacetime $M_\rho = \Omega(\Lambda_\rho) \backslash \rho(\Gamma)$ is then a GHCM conformally flat extension of the conformally flat Misner spacetime $\Omega(\Lambda) \backslash \rho(\Gamma)$. Extensions of conformally flat Misner spacetimes form a larger class of GHCM conformally flat spacetimes with complete photons.

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